

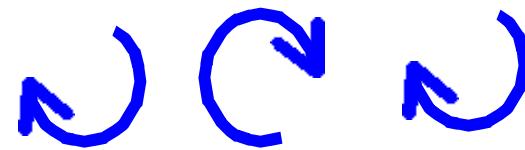
Steady-State Accretion Disc

- **thickness** $\frac{H}{R} \approx \frac{c_s}{V_q}$ $V_q = \left(\frac{GM}{R} \right)^{1/2}$
- **surface density**
 - (kg/m²) $\Sigma \equiv \int_{-\infty}^{\infty} r dz = \sqrt{2p} r_0 H$
- **accretion rate** $\dot{M} = 2p R (-V_R) \Sigma$
 - (kg / s) $= 3p n \Sigma \left(1 - \left(\frac{R_*}{R} \right)^{1/2} \right)$
- **inflow velocity** $V_R \approx -\frac{3n}{2R}$
- **viscosity**
 - (m² / s) $n \equiv \alpha c_s H$
 - -- alpha model (hides uncertain physics)

Anomalous Viscosity

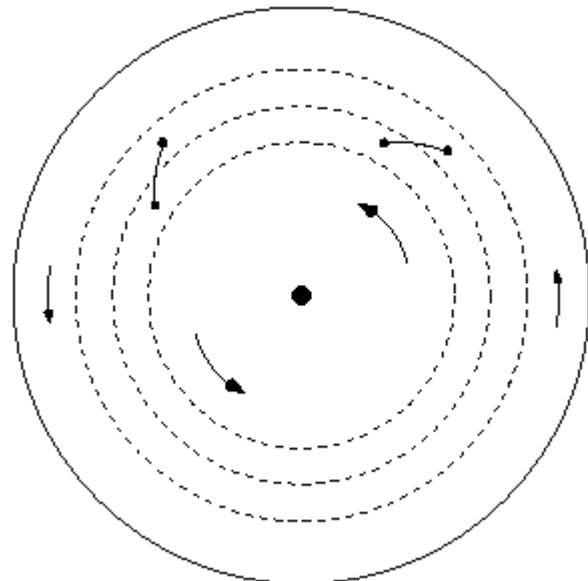
- Viscosity arises from turbulent eddies
- largest eddie size $\sim H$
- largest eddie velocity \sim sound speed

$$n = \alpha c_s H$$

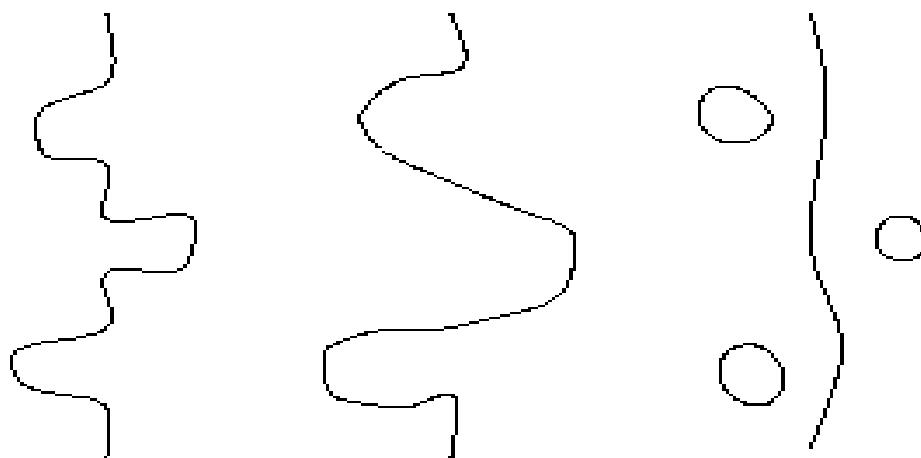
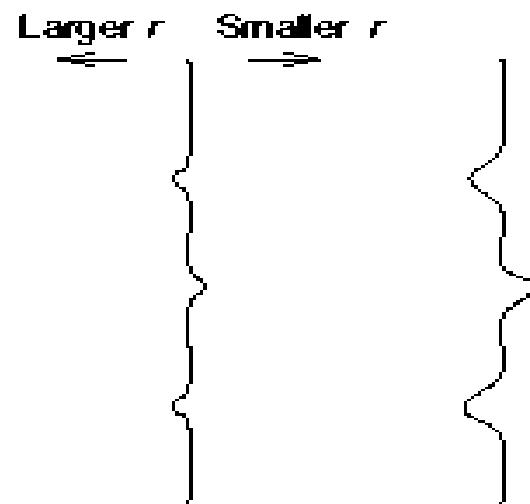


- alpha = dimensionless scale factor
- $0 < \alpha < 1$

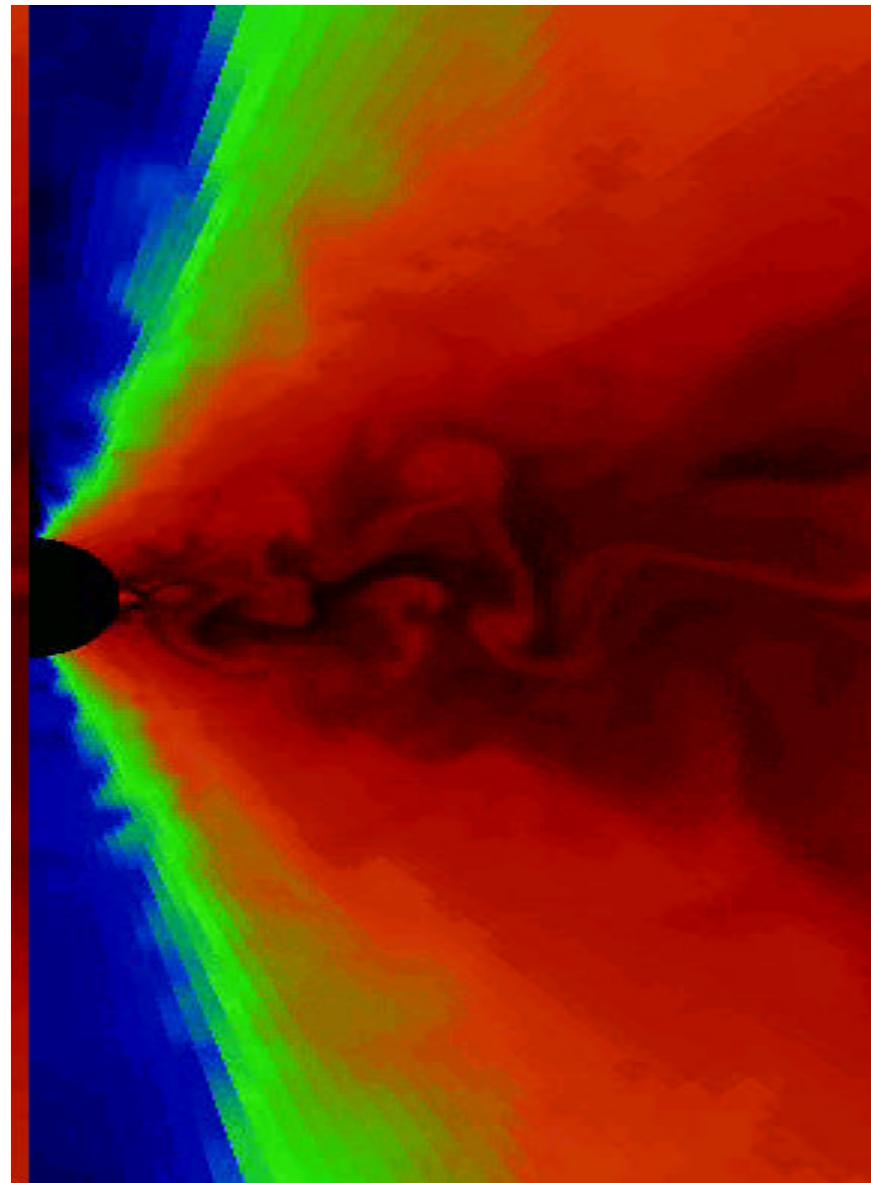
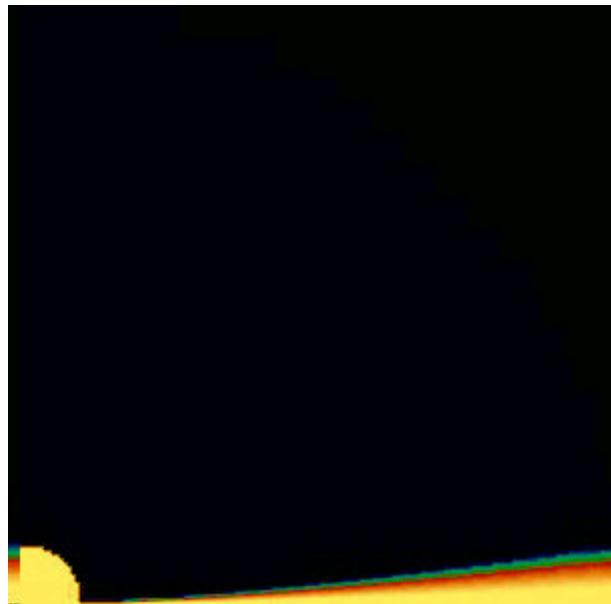
Magneto-Rotational Instability



**Magnetic fields link
different annuli
generate MHD turbulence**



Magneto-Rotational Instability



Timescales

- **Dynamical timescale -- Kepler orbit time**

$$t_{dyn} \sim \frac{R}{V_q} \sim \Omega^{-1}$$

- **Vertical hydrostatic equilibrium timescale**

$$t_{hse} \sim \frac{H}{c_s} \sim \Omega^{-1}$$

- **Radial inflow time**

$$t_{visc} \sim \frac{R}{-V_R} \sim \frac{R^2}{n} \sim \frac{R^2}{a c_s H} \sim a^{-1} \left(\frac{R}{H} \right)^2 \Omega^{-1}$$

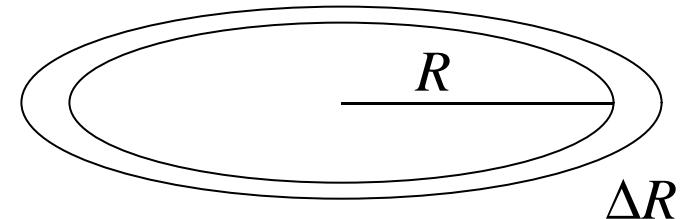
Steady-State Disc $T(R)$

- Energy (kinetic + potential) per mass

$$e = \frac{V_q^2}{2} + \Phi_{grav} = \frac{GM}{2R} - \frac{GM}{R} = -\frac{1}{2} \frac{GM}{R} \quad \frac{de}{dR} = \frac{1}{2} \frac{GM}{R^2}$$

- energy radiated = energy released

$$s T^4 \times 2\pi R \Delta R \times 2 = \dot{M} \frac{de}{dR} \Delta R$$



- Temperature profile

$$T^4 = \frac{GM\dot{M}}{8\pi s R^3} \quad T \propto (M \dot{M})^{1/4} R^{-3/4}$$

- Disc Luminosity

$$L_{disk} = 2 \int_{R_*}^{\infty} 2\pi R dR s T^4 = \frac{GM\dot{M}}{2} \int_{R_*}^{\infty} \frac{dR}{R^2} = \frac{GM\dot{M}}{2R_*}$$

Note: 1/2 of energy remains as kinetic energy

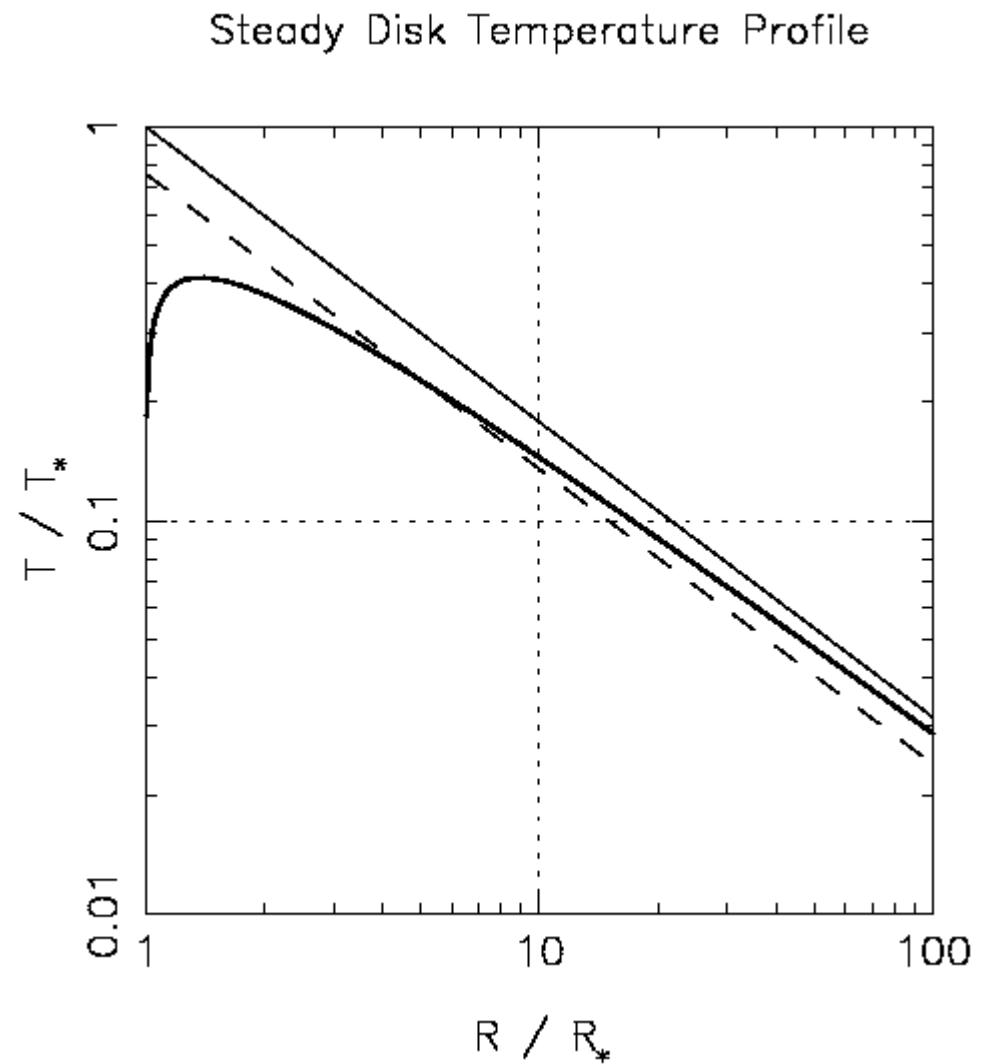
Temperature Profile

- Include work done by viscous torques

$$T^4 = \frac{3 G M \dot{M}}{8 \rho s R^3} \left[1 - \left(\frac{R_*}{R} \right)^{1/2} \right]$$
$$= T_*^4 \left(\frac{R}{R_*} \right)^{-3} \left[1 - \left(\frac{R_*}{R} \right)^{1/2} \right]$$

$$T_* = \left(\frac{3 G M \dot{M}}{8 \rho s R_*^3} \right)^{1/4}$$

Inner disc luminosity lower.
Outer disk luminosity 3x higher.
max T occurs outside min R



Maximum Temperature

- Including work done by viscous torques

$$\left(\frac{T}{T_*}\right)^4 = \left(\frac{R}{R_*}\right)^{-3} \left[1 - \left(\frac{R_*}{R}\right)^{1/2} \right] \quad T_* = \left(\frac{3 G M \dot{M}}{8 \rho s R_*^3} \right)^{1/4}$$

- Maximum temperature

$$\begin{aligned} \frac{d}{dR} [T^4] &= T_*^4 \frac{d}{dR} \left[\frac{1}{R^3} - \frac{R_*^{1/2}}{R^{5/2}} \right] \\ &= T_*^4 \left[-\frac{3}{R^4} + \frac{5}{2} \frac{R_*^{1/2}}{R^{7/2}} \right] \rightarrow \boxed{\frac{R}{R_*} = \left(\frac{6}{5}\right)^2 = 1.44} \end{aligned}$$

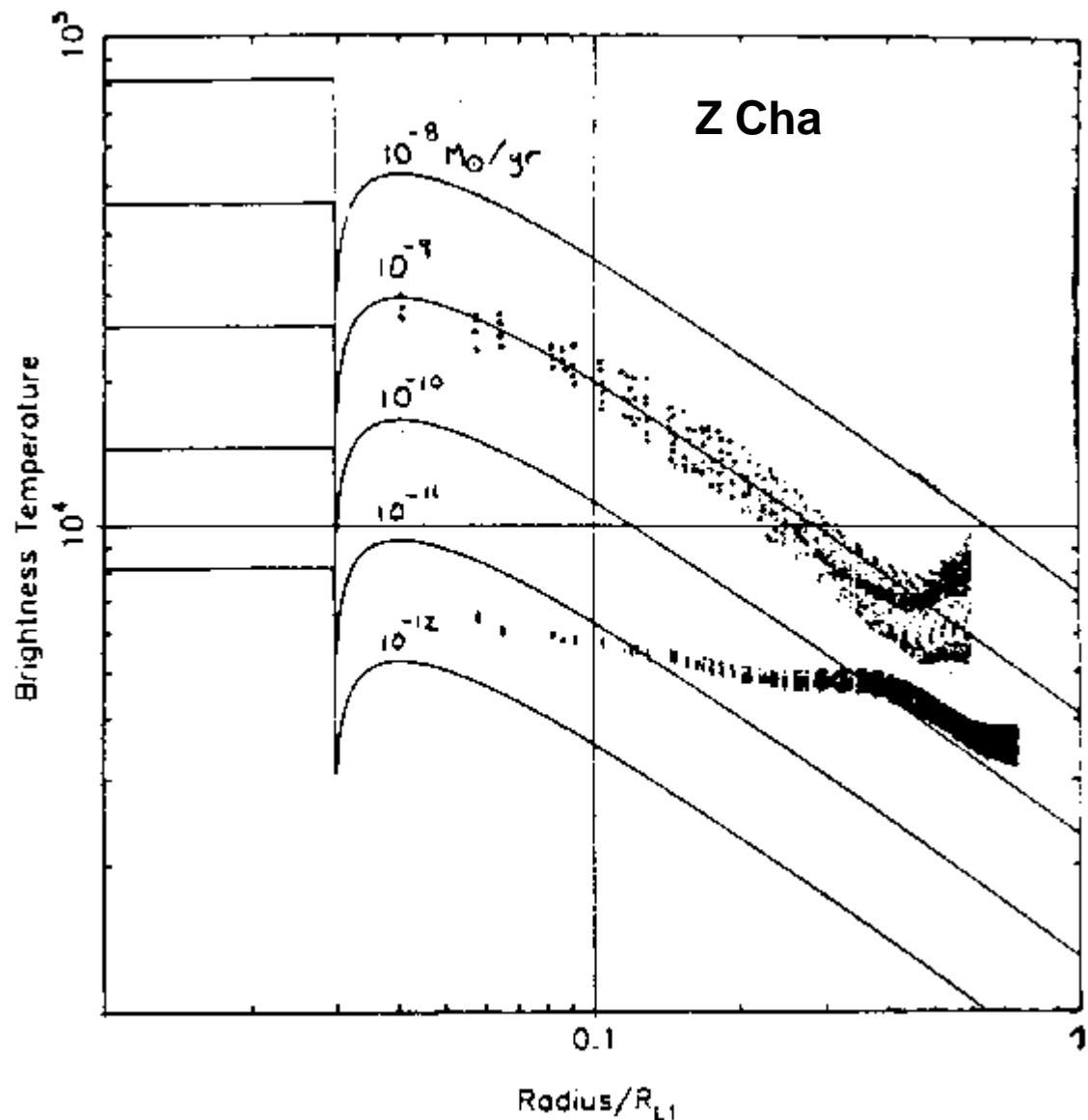
$$\boxed{\frac{T_{\max}}{T_*} = \left(\frac{5}{6}\right)^{3/2} \left(1 - \frac{5}{6}\right)^{1/4} = \frac{5^{3/2}}{6^{7/4}} \approx 0.486}$$

Observed Temperature Profile

Eclipse lightcurve shapes used to map temperature profile of disk.

Outburst disk matches steady disk theory

Quiescent disk does not (optically thin)



Disc Luminosity

- Including work done by viscous torques

$$T^4 = \frac{3 G M \dot{M}}{8 \rho s R^3} \left[1 - \left(\frac{R_*}{R} \right)^{1/2} \right]$$

- Disc luminosity $L_{disk} = 2 \int_{R_*}^{\infty} 2 \rho R dR s T^4$

$$= \frac{G M \dot{M}}{2} \int_{R_*}^{\infty} dR \left(\frac{1}{R^2} - \frac{R_*^{1/2}}{R^{5/2}} \right)$$

$$= \frac{G M \dot{M}}{2} \left[-\frac{1}{R} + \frac{3}{2} \frac{R_*^{1/2}}{R^{3/2}} \right]_{R_*}^{\infty}$$

$$= \frac{G M \dot{M}}{2 R_*}$$

Same as before.

Energy is conserved but
re-distributed in radius by
viscous torques.

Blackbody Disc Spectrum

$$F_n = \int B_n(T(R)) \frac{2\pi R dR \cos i}{D^2}$$

$$B_n(T) = \frac{2 h n^3}{c^2 (e^x - 1)} = \frac{2 k^3 T^3}{c^2 h^2} \frac{x^3}{e^x - 1} \quad x \equiv \frac{h n}{k T}$$

$$\frac{T}{T_*} = \left(\frac{R}{R_*} \right)^{-3/4} = \left(\frac{x}{x_*} \right)^{-1} \rightarrow B_n(I, T) = \frac{2 k^3 T_*^3}{c^2 h^2} \frac{x_*^3}{e^x - 1}$$

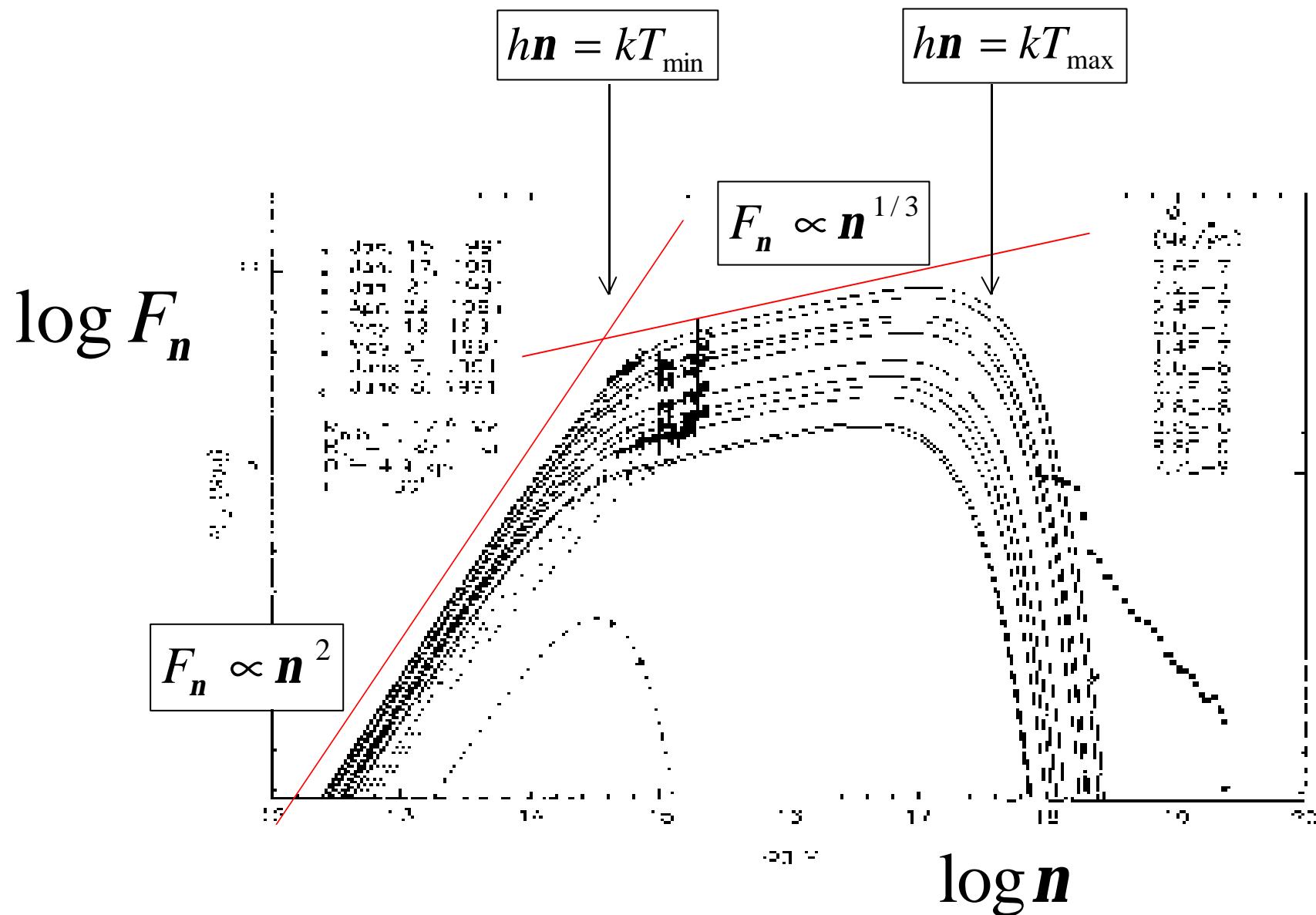
$$\frac{R}{R_*} = \left(\frac{x}{x_*} \right)^{4/3} \quad \frac{dR}{R_*} = \frac{4}{3} \left(\frac{x}{x_*} \right)^{1/3} \left(\frac{dx}{x_*} \right)$$

Blackbody Disc Spectrum

$$\begin{aligned} F_n &= \int_{R_*}^{\infty} B_n(T(R)) \frac{2\mathbf{p} R dR \cos i}{D^2} \\ &= \int_{x_*}^{\infty} \frac{2(k T_*)^3}{(c h)^2} \frac{x_*^3}{e^x - 1} \frac{2\mathbf{p} R_*^2 \cos i}{D^2} \frac{4}{3} \left(\frac{x}{x_*} \right)^{5/3} \left(\frac{dx}{x_*} \right) \\ &= \frac{16\mathbf{p}}{3} \frac{R_*^2 \cos i}{D^2} \frac{(k T_*)^3}{(h c)^2} \left(\frac{h \mathbf{n}}{k T_*} \right)^{1/3} \int_{x_*}^{\infty} \frac{x^{5/3} dx}{e^x - 1} \end{aligned}$$

$$F_n \propto \mathbf{n}^{1/3}$$

Blackbody Disc Spectrum



Surface Density Evolution

$$\Delta m = 2\mathbf{p} \cdot \mathbf{R} \Delta R \Sigma$$

$$\Delta \ell = 2\mathbf{p} \cdot \mathbf{R} \Delta R \Sigma R^2 \Omega$$

mass conservation :

$$\frac{\partial}{\partial t} (2\mathbf{p} \cdot \mathbf{R} \Delta R \Sigma) = V_R 2\mathbf{p} \cdot \mathbf{R} \Sigma \Big|_{R+\Delta R}^R$$

$$R \dot{\Sigma} + \frac{\partial}{\partial R} (R \Sigma V_R) = 0$$

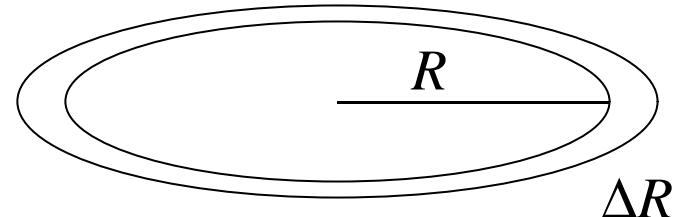
angular momentum :

$$R \frac{\partial}{\partial t} (\Sigma R^2 \Omega) + \frac{\partial}{\partial R} (R \Sigma V_R R^2 \Omega) = \frac{1}{2\mathbf{p}} \frac{\partial G}{\partial R}$$

viscous torque : $G(R,t) = 2\mathbf{p} \cdot \mathbf{R} \mathbf{n} \Sigma R^2 \frac{\partial \Omega}{\partial R}$

surface density evolution (diffusion) :

$$\dot{\Sigma} = \frac{3}{R} \frac{\partial}{\partial R} \left(R^{1/2} \frac{\partial}{\partial R} [\mathbf{n} \cdot \Sigma R^{1/2}] \right)$$



simple diffusion :

$$\dot{\Sigma} = \mathbf{n} \frac{\partial^2 \Sigma}{\partial R^2}$$

Ring spreads to form a disc

Pringle, Lynden-Bell 1974
MNRAS 168 603

Surface density
evolution

diffusion in radius

