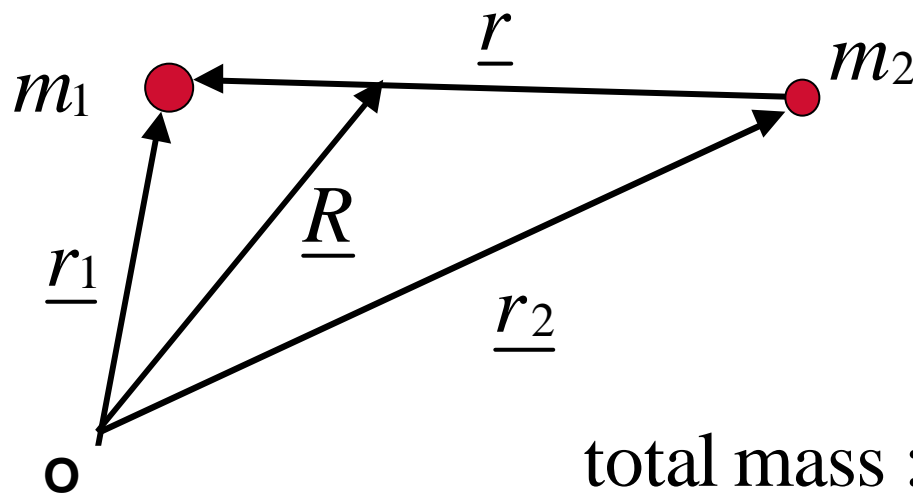


The Two-Body Problem

- **Newtonian gravity**
- **2 point masses**
 - good approx for stars



$$\underline{r} = \underline{r}_1 - \underline{r}_2$$

$$\hat{r} = \underline{r} / r$$

total mass : $M = m_1 + m_2$

centre of mass : $\underline{R} = \frac{m_1 \underline{r}_1 + m_2 \underline{r}_2}{M}$

– **O: origin**

Two-Body motion

- **Newtons laws of motion + law of gravitation**

Equations of motion :

$$\underline{F}_1 = m_1 \underline{\ddot{r}}_1 = \frac{-G m_1 m_2}{r^2} \hat{r} \qquad \underline{F}_2 = m_2 \underline{\ddot{r}}_2 = -\underline{F}_1$$

add the equations $m_1 \underline{\ddot{r}}_1 + m_2 \underline{\ddot{r}}_2 = 0$

integrate dt $m_1 \underline{\dot{r}}_1 + m_2 \underline{\dot{r}}_2 = \underline{A}$

again $m_1 \underline{r}_1 + m_2 \underline{r}_2 = \underline{A} t + \underline{B}$

definition of \underline{R} $M \underline{R} = \underline{A} t + \underline{B}$

- **\therefore Centre of mass moves with constant velocity**
 - (unless acted on by an external force)

Relative Motion

- Important for eclipses
- Subtract two equations of motion

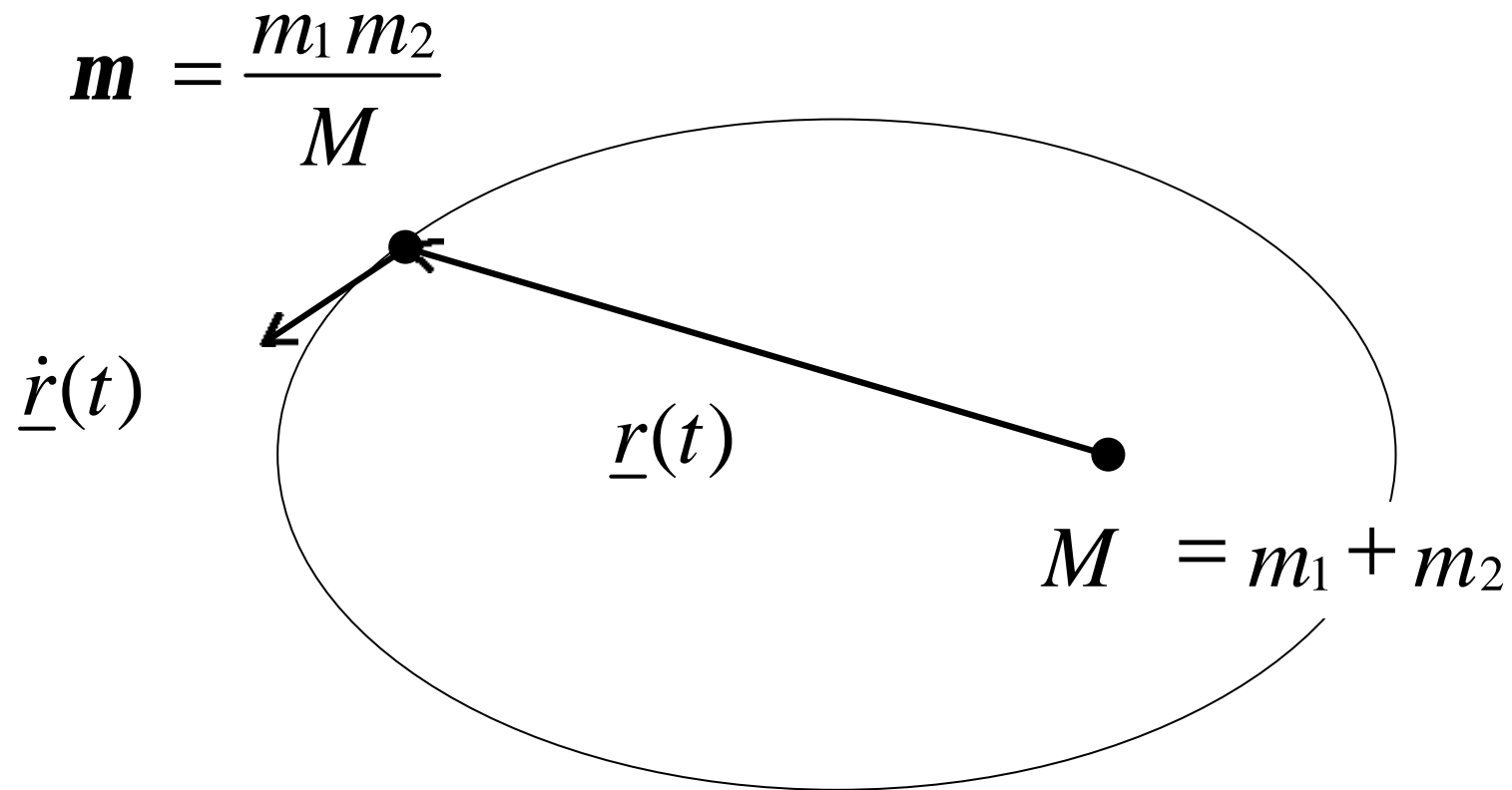
$$\underline{\ddot{r}} = \underline{\ddot{r}}_1 - \underline{\ddot{r}}_2 = \frac{-G M}{r^2} \underline{\hat{r}}$$

- Multiply by $\underline{m} = m_1 m_2 / M =$ “reduced mass”

$$\underline{m} \underline{\ddot{r}} = \frac{-G M \underline{m}}{r^2} \underline{\hat{r}} = \frac{-G m_1 m_2}{r^2} \underline{\hat{r}}$$

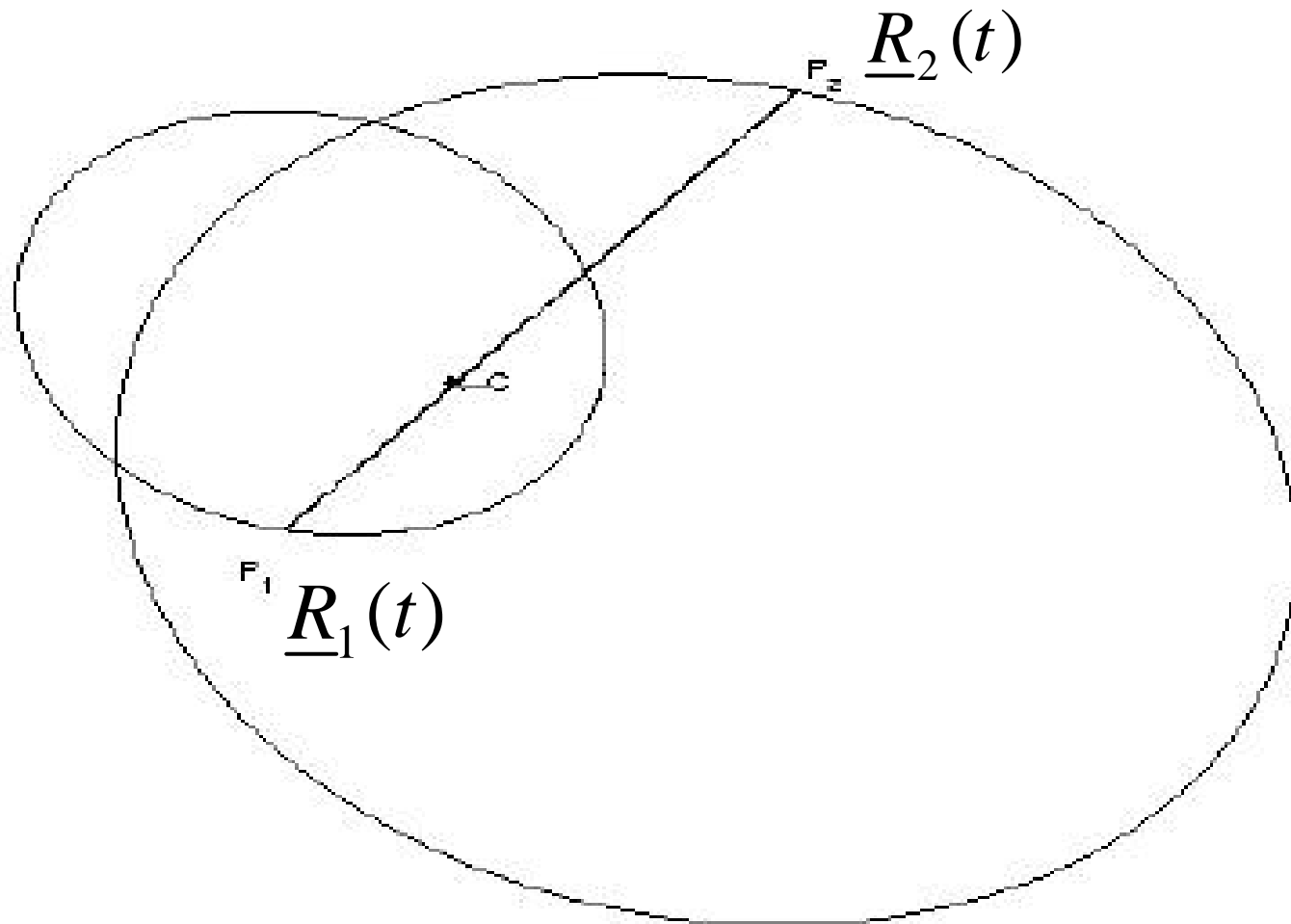
- Relative orbit is as if:
 - orbiter has reduced mass $\underline{m} =$ reduced mass
 - stationary central mass is $M =$ total mass.

The Relative Orbit



Important for Eclipses

2 Barycentric Orbits



Important for Radial Velocities

Barycentric Orbits

- **Important for radial velocity curves**

$$m_1 \underline{R}_1 + m_2 \underline{R}_2 = 0 \quad \text{centre of mass frame}$$

$$\underline{r} = \underline{R}_1 - \underline{R}_2 = \frac{m_1 + m_2}{m_2} \underline{R}_1 = \frac{M}{m_2} \underline{R}_1 = -\frac{M}{m_1} \underline{R}_2$$

equations of motion :

$$\ddot{\underline{R}}_1 = -\frac{G m_2}{r^3} \underline{r} \quad \ddot{\underline{R}}_2 = -\frac{G m_1}{r^3} (-\underline{r})$$

- Eliminate r^3 using $r = f(R_1)$ and $r = f(R_2)$:

$$\ddot{\underline{R}}_1 = -\frac{G m_2^3}{M^2} \frac{\underline{R}_1}{R_1^3} \quad \ddot{\underline{R}}_2 = -\frac{G m_1^3}{M^2} \frac{\underline{R}_2}{R_2^3}$$

- the acceleration of each star relative to the centre of mass

Relative vs Barycentric orbits

- (a, e, P, v) relative orbit
- $(a, e, P, v)_{1,2}$ barycentric orbits
- m_1 and m_2 on straight line thru C

$$P_1 = P_2 = P \quad e_1 = e_2 = e$$

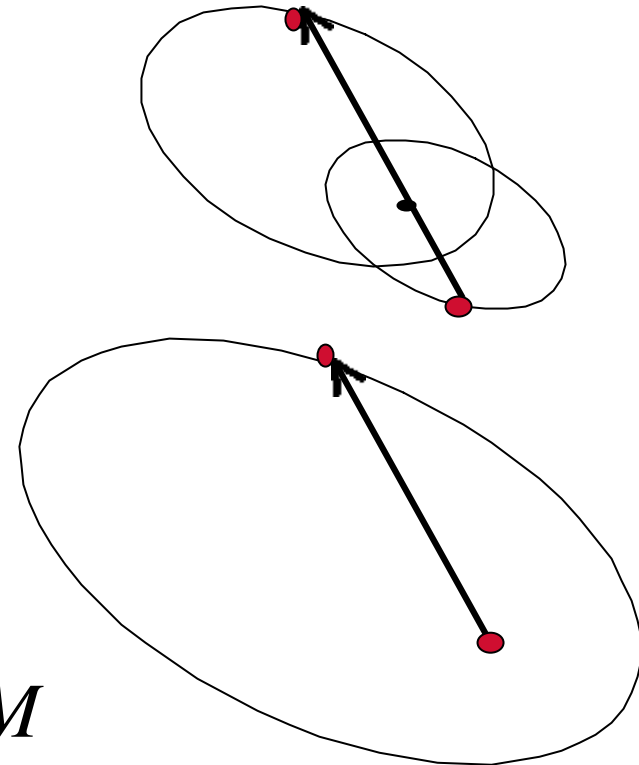
$$a_1 = a m_2 / M \quad a_2 = a m_1 / M$$

$$a = a_1 + a_2 \quad M = m_1 + m_2$$

$$a_1 : a_2 : a = V_1 : V_2 : V = m_2 : m_1 : M$$

Kepler :

$$\frac{4p^2}{GP^2} = \frac{M}{a^3} = \frac{m_2^3 / M^2}{a_1^3} = \frac{m_1^3 / M^2}{a_2^3}$$



Orbital Speed

orbital speed : $V^2 = \dot{r}^2 + r^2 \dot{\mathbf{q}}^2$

conic section : $r = \frac{\ell}{1 + e \cos \mathbf{q}}$

$$\frac{d}{dt} \left[1 + e \cos \mathbf{q} = \frac{\ell}{r} \right] \Rightarrow -e (\sin \mathbf{q}) \dot{\mathbf{q}} = -\frac{\ell}{r^2} \dot{r}$$

specific angular momentum : $r^2 \dot{\mathbf{q}} = L$

$$\therefore \dot{r} = \frac{r^2 \dot{\mathbf{q}}}{\ell} e \sin \mathbf{q} = \frac{L}{\ell} e \sin \mathbf{q}$$

$$r \dot{\mathbf{q}} = \frac{L}{r} = \frac{L}{\ell} (1 + e \cos \mathbf{q})$$

Orbital Speed

$$\begin{aligned}V^2 &= \left(\frac{L}{\ell}\right)^2 \left[e^2 \sin^2 \mathbf{q} + (1 + e \cos \mathbf{q})^2 \right] \\&= \left(\frac{L}{\ell}\right)^2 \left[e^2 + 1 + 2e \cos \mathbf{J} \right] \\&= \left(\frac{L}{\ell}\right)^2 \left[2(e \cos \mathbf{q} + 1) + e^2 - 1 \right] \\V^2 &= \frac{L^2}{\ell} \left[\frac{2}{r} - \frac{1 - e^2}{\ell} \right]\end{aligned}$$

$$\begin{aligned}\ell &= \frac{L^2}{G M} \\ \frac{L}{\ell} &= \frac{G M}{L}\end{aligned}$$

Orbital Speed

$$V^2 = \frac{L^2}{\ell} \left[\frac{2}{r} - \frac{1-e^2}{\ell} \right]$$

ellipse : $\ell = a(1 - e^2) \quad L^2 = GM \ell$

$$V^2 = GM \left[\frac{2}{r} - \frac{1}{a} \right]$$

<i>ellipse</i>	$e < 1$	$\ell = a(1 - e^2)$	$V^2 = GM \left[\frac{2}{r} - \frac{1}{a} \right]$
<i>circle</i>	$e = 0$	$\ell = a$	$V^2 = \frac{GM}{a}$
<i>parabola</i>	$e = 1$		$V^2 = \frac{2GM}{r}$
<i>hyperbola</i>	$e > 1$	$\ell = a(1 - e^2)$	$V^2 = GM \left[\frac{2}{r} + \frac{1}{a} \right]$

Energy of Orbit

Kinetic energy:

$$KE = T = \frac{1}{2} m_1 V_1^2 + \frac{1}{2} m_2 V_2^2 = \frac{m_1 m_2}{2 M} V^2$$

orbital speed: $V^2 = G M \left(\frac{2}{r} - \frac{1}{a} \right)$

Potential energy:

$$PE = W = -\int_r^{\infty} \frac{G m_1 m_2}{r^2} dr = -\frac{G m_1 m_2}{r}$$

Total energy: $E = T + W = -\frac{G m_1 m_2}{2 a} < 0$

Binding energy: $-E > 0$

Angular momentum of the orbit

- **Angular momentum vector \underline{J} , defines orbital plane**

- $\underline{J} = m_1 \underline{L}_1 + m_2 \underline{L}_2$ and $L^2 = G M a (1-e^2)$
and $L_1^2 = G (m_2^3 / M^2) a_1 (1-e^2)$
and $a_1/a = m_2 / M$

- same for \underline{L}_2

- hence
$$L_1 = \frac{m_2^2}{M^2} L; \quad L_2 = \frac{m_1^2}{M^2} L$$

therefore

$$J^2 = \frac{G m_1^2 m_2^2}{M} a (1-e^2)$$

and the final expression for J is

$$J = \frac{2\mathbf{p} a^2 m_1 m_2}{P M} \sqrt{1-e^2}$$

Orbital Angular momentum

- **Given masses m_1, m_2 and Energy E ,**
 - the angular momentum J determines the shape of the orbit
 - ie the eccentricity (or the conic section parameter l)
- **For given E ,**
 - circular orbits have maximum J
 - J decreases as $e \rightarrow 1$
 - orbit becomes rectilinear ellipse
- **relation between E , and J very important in**
determining when systems interact
mass exchange and orbital evolution