

Inclinations

- Radial velocities give minimum masses

SB2: $m \sin^3 i$ SB1: $\frac{m^3 \sin^3 i}{M^2}$

- need inclinations
 - astrometric orbits
 - eclipsing binaries
 - polarimetry
 - apsidal motion

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Astrometric Orbits

Interferometry measures the relative orbit (milli-arcsec accuracy).
Fit elliptical orbit model to $x(t)$ and $y(t)$.

Get inclination, hence masses

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Polarimetry

Scattered light is polarised (e.g. our sky)

Get inclination, hence masses

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Eclipsing Binaries

- Binary properties from eclipses
 - sizes, shapes of stars, inclination,
 - temperatures, limb darkening, apsidal motion
 - 4 contact phases:

primary eclipse secondary eclipse

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Partial Eclipse

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Eclipses

$R_1 > R_2 \quad q \equiv 2p f$

edge-on ($i = 90^\circ$):
4 contact phases:
 $R_1 \pm R_2 = a \sin q$

$$\frac{R_2}{a} = \frac{\sin q_1 - \sin q_2}{2}$$

$$\frac{R_1}{a} = \frac{\sin q_1 + \sin q_2}{2}$$

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Eclipses

4 contact phases

$$R_1 \pm R_2 = a \sqrt{\sin^2 q + \cos^2 q \cos^2 i}$$

$$= a \sqrt{1 + \cos^2 q \sin^2 i}$$

4 measurements: f_1, f_2, f_3, f_4

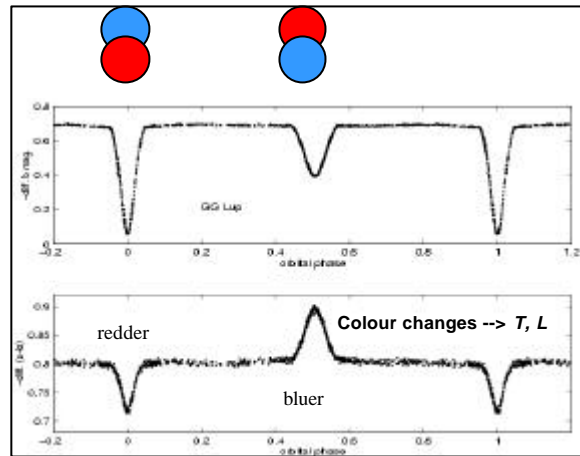
4 parameters: $\frac{R_1}{a}, \frac{R_2}{a}, i, f_0$

mid-eclipse $q = 0$

total eclipse $a \cos i < R_1 - R_2$

partial eclipse $R_1 - R_2 < a \cos i < R_1 + R_2$

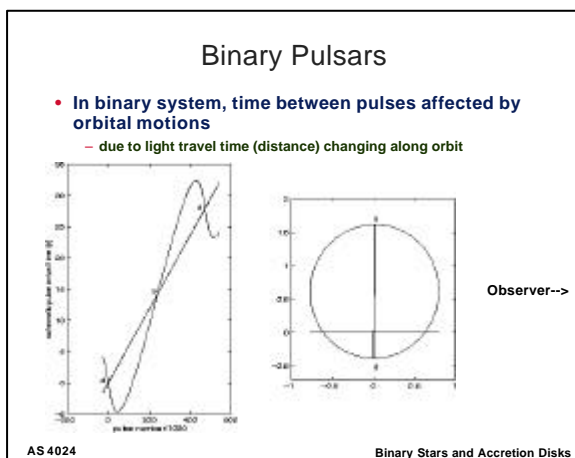
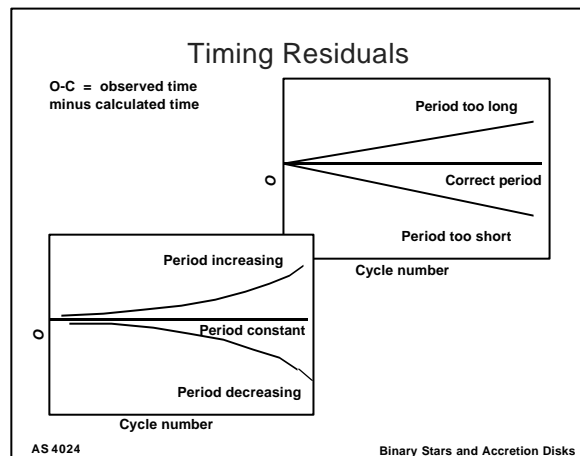
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Application to Binary Pulsars

- binary system where one star is a pulsar**
 - emits 'pulses' of radiation
 - accurate timing possible (accurate clocks)
 - need narrow pulses
 - radio signals from neutron stars
- solitary pulsar**
 - if at 0 velocity relative to us
 - time between pulses, $dt = \text{constant}$ (unless being spun up/down)
 - if at V_{rel} relative velocity
 - $dt = \text{constant} \times \text{pulse number}$
 - if pulsar spins up, dt decreases with pulse number
 - concave curve
 - if pulsar spins down, dt increases with pulse number
 - convex curve

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Light travel time

pulsar orbit: $r_p = \frac{a_p (1 + e^2)}{1 + e \cos q}$

distance along line of sight:

$$z_p = r_p \sin i \sin(q + w)$$

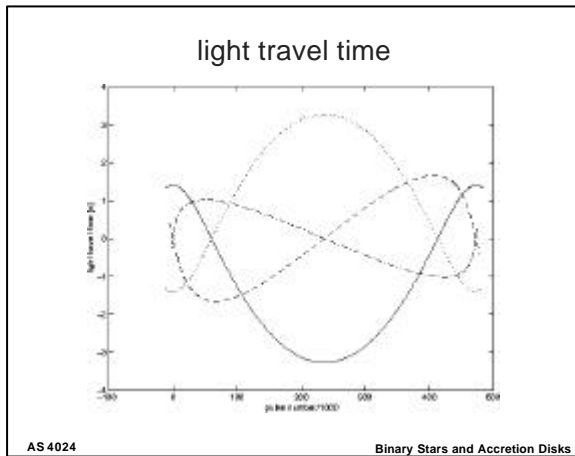
light travel time:

$$\frac{z_p}{c} = \frac{a_p \sin i}{c} (1 + e^2) \frac{\sin(q + w)}{1 + e \cos q}$$

circular orbit

$$\frac{z_p}{c} = \frac{a_p \sin i}{c} \sin\left(\frac{2p}{P}(t - T_0)\right)$$

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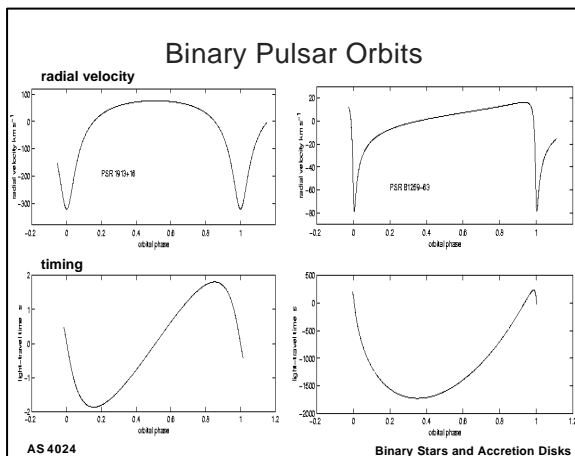


Binary Pulsar timing residuals

- Time difference between predicted, $j_n t$, and actual (binary) pulse arrival times, t_n is

$$\Delta t = t_n - j_n t = a t + b \sin \left[\frac{2\pi (t - T_0)}{P} \right]$$
 - P is the orbital period, T_0 is a reference time
 - a, b are determined by the velocity of the pulsar
 - a : from systematic velocity
 - b : from Keplerian velocity
 - for circular orbits: $b = (a_p/c) \sin i$

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Mass determinations

- visible companion star
 - O-B star in High-Mass X-ray Binaries (HMXB)
 - A-K star in Low-Mass X-ray Binaries (LMXB)

$$a_c \sin i = \frac{(1 - e^2)^{1/2} K_c P}{2\pi}$$

mass function

$$f(m_p) = \frac{m_p^3 \sin^3 i}{M^2} = \frac{(1 - e^2)^{3/2} K_c^3 P}{2\pi G}$$

mass ratio, q ,

$$q = \frac{m_p}{m_c} = \frac{a_c \sin i}{a_p \sin i}$$

- If inclination, i , can be found, then masses follow

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Frequency shifts

- Binary orbit also affects pulsar frequency
 - radio pulsars, very narrow pulse widths
 - pulse frequency affected by orbital velocity
 - Doppler shift:

$$\Delta f = f \frac{V_{rad}}{c} = f \frac{\dot{z}}{c}$$

- gives a phase lag of:

$$\Delta \mathbf{f} = \int_{T_0}^t \Delta f dt \approx f_0 \frac{\dot{z}}{c} (t - T_0)$$

$$= f_0 \frac{\dot{z}}{c} \left[\frac{z}{c} \right]$$

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Pulsar Phase lag

- Combined phase lag is
 - from light travel time due to orbit

$$\Delta \mathbf{f}_L = -f \frac{z}{c}$$

- and from Doppler shift

$$\Delta \mathbf{f}_D = f_0 \frac{\dot{z}}{c} \left[\frac{z}{c} \right]$$

- hence

$$\Delta \mathbf{f} = \Delta \mathbf{f}_D + \Delta \mathbf{f}_L \approx \left[\frac{z}{c} \right] \left[f_0 \frac{\dot{z}}{c} - f_0 - \dot{f}_0 (t - T_0) \right]$$

- generally $\Delta \mathbf{f}_D \approx 0.001 \Delta \mathbf{f}_L$
- but measurable in radio pulsars

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