

## Perturbed Two-Body motion

- “close” to 2-body motion
  - planetary systems
  - multi-star systems (triples, quadruples, ...)
  - close binaries (non-spherical stars)
  - general relativity corrections (Mercury, binary pulsars)
- Perturbed potential

$$\ddot{\mathbf{O}} = \ddot{\mathbf{O}}_0 + \mathbf{S}$$

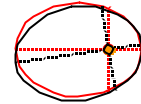
point mass potential
small corrections  
the disturbing function

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## Lagrange's Planetary Equations

- Developed for Solar System Planets
  - main potential due to the Sun
  - secondary potentials from other planets
  - slow secular and/or periodic changes to orbital elements due to orbit-averaged effect of S
- Milankovitch Cycles
  - effects on Earth's climate
- Apsidal Motion
  - precession of the orbit in its own plane
  - e.g. Mercury (43 arcsec / century)
  - tests General Relativity
  - observable in binaries (few deg / year)
  - tests stellar structure theory

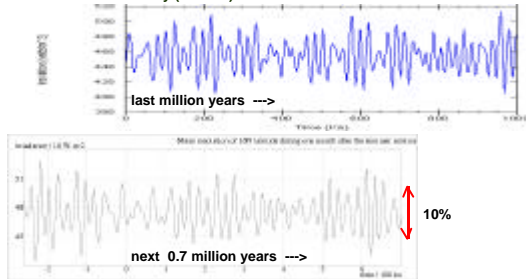


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## Milankovitch Cycles

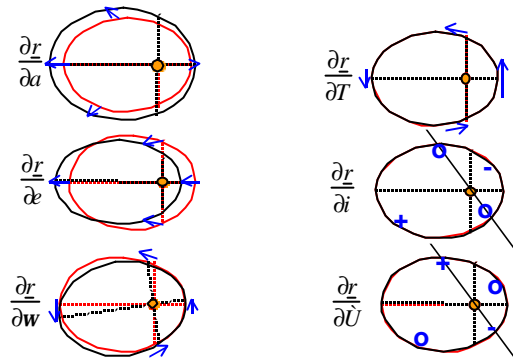
- Northern hemisphere insolation, including
  - rotation axis precession (23 ka), tilt (41 ka)
  - orbit eccentricity (100 ka)



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## changes in 6 orbital elements



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## Lagrange's Equations

Orbit depends on time and 6 orbital elements

$$\mathcal{L}(t, a, e, i, \dot{u}, \dot{w}, \dot{t}) \quad 4\pi^2 G M P^2 = a^3$$

$$\dot{\mathbf{r}} = \frac{\partial \mathbf{r}}{\partial t} + \sum_{i=1}^6 \frac{\partial \mathbf{r}}{\partial e_i} \dot{e}_i \quad \text{where } e_i = (a, e, i, \dot{u}, \dot{w}, \dot{t})$$

$$\frac{\partial^2 \mathbf{r}}{\partial t^2} + \frac{d}{dt} \left[ \sum_{i=1}^6 \frac{\partial \mathbf{r}}{\partial e_i} \dot{e}_i \right] = -\nabla \Phi_0 - \nabla S$$

2-body orbit      slow changes in orbital elements

$$\frac{d}{dt} \left[ \sum_{i=1}^6 \frac{\partial \mathbf{r}}{\partial e_i} \dot{e}_i \right] \cdot \frac{\partial \mathbf{r}}{\partial e_k} = -\nabla S \cdot \frac{\partial \mathbf{r}}{\partial e_k}$$

$$\sum_i M_{ik} \dot{e}_i = - \int \frac{\partial S}{\partial e_k} dt \quad \rightarrow \quad \dot{e}_i = - \sum_k M_{ik}^{-1} \int \frac{\partial S}{\partial e_k} dt$$

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## Lagrange's planetary equations

$$\dot{a} = \frac{2}{n a} \frac{\partial S}{\partial \mathbf{c}}; \quad \dot{e} = \frac{1}{n a^2 e} \left[ (1-e^2) \frac{\partial S}{\partial \mathbf{c}} - \sqrt{1-e^2} \frac{\partial S}{\partial \mathbf{w}} \right]$$

$$\dot{\mathbf{c}} = - \frac{(1-e^2)}{n a^2 e} \frac{\partial S}{\partial e} - \frac{2}{n a} \frac{\partial S}{\partial a}; \quad \dot{U} = \frac{1}{n a^2 \sqrt{1-e^2} \sin i} \frac{\partial S}{\partial i}$$

$$\dot{\mathbf{w}} = \frac{\sqrt{1-e^2}}{n a^2 e} \frac{\partial S}{\partial e} - \frac{\cot i}{n a^2 \sqrt{1-e^2}} \frac{\partial S}{\partial i}$$

$$\dot{i} = \frac{1}{n a^2 \sqrt{1-e^2}} \left[ \cot i \frac{\partial S}{\partial \mathbf{w}} - \text{cosec} i \frac{\partial S}{\partial U} \right]$$

where  $n^2 a^3 = G M$  and  $\mathbf{c} = -n T$ ,  
 $n = 2\pi / P$  is the mean daily motion

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### apsidal motion

- longitude of periastron,  $\varpi$ , changes as the orbit precesses

$$\dot{\varpi} = \frac{\sqrt{1-e^2}}{n a^2 e} \frac{\partial S}{\partial e} - \frac{\cot i}{n a^2 \sqrt{1-e^2}} \frac{\partial S}{\partial i}$$

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### relativistic apsidal motion

- Mercury perihelion precession
- binary neutron stars (pulsars)

gravitational potential (Newton + small GR correction)

$$\dot{O} = -\frac{GM}{r} \left[ 1 + \left( \frac{V}{c} \right)^2 + \dots \right] \quad V = L/r \quad L^2 = GM \ell = GM a (1-e^2)$$

disturbing function:

$$S = -\frac{G^2 M^2 a (1-e^2)}{c^2 r^3}$$

small inward acceleration

$$\frac{\partial S}{\partial r} = -\frac{G^2 M^2 a (1-e^2)}{c^2 r^4}$$

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### relativistic apsidal motion

$$S = -\frac{G^2 M^2 a (1-e^2)}{c^2 r^3}$$

$$\dot{\varpi} = \frac{\sqrt{1-e^2}}{n a^2 e} \frac{\partial S}{\partial e} - \frac{\cot i}{n a^2 \sqrt{1-e^2}} \frac{\partial S}{\partial i}$$

$$\frac{\partial S}{\partial e} \equiv \left\langle \nabla S \cdot \frac{\partial \underline{r}}{\partial e} \right\rangle$$

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orbit average of disturbing function:

$$\langle S \rangle = -\left\langle \frac{G^2 M^2 a (1-e^2)}{c^2 r^3} \right\rangle = -\left( \frac{GM}{ca} \right)^2 \frac{1}{(1-e^2)^{3/2}}$$

$\left\langle \frac{1}{r^3} \right\rangle = \frac{1}{P_0} \int_0^P \frac{dt}{r^3} = \frac{1}{P_0} \int_0^{2\pi} \frac{1}{r^2} \frac{dq}{\dot{q}}$ $= \frac{1}{P_0} \int_0^{2\pi} \frac{(1+e \cos q)}{\ell} \frac{dq}{L}$ $= \frac{2\pi}{P \ell L}$ $= \frac{1}{a^3 (1-e^2)^{3/2}}$	$r^2 \dot{q} = L = \sqrt{GM \ell}$ $r = \frac{\ell}{1+e \cos q}$ $\ell = a (1-e^2)$ $\frac{2\pi}{P L} = \sqrt{\frac{GM}{a^3}} \frac{1}{\sqrt{GM \ell}}$ $= \frac{1}{a^3 \sqrt{1-e^2}}$
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### An incorrect calculation

<p>apsidal motion :</p> $\dot{\varpi} = \frac{\sqrt{1-e^2}}{n a^2 e} \frac{\partial S}{\partial e}$ $= \frac{G^2 M^2}{n c^2 a^4 (1-e^2)}$ $= \frac{2\pi G M}{P c^2 a (1-e^2)}$ <p>Factor of 3 too small :(</p>	<p>disturbing function :</p> $\langle S \rangle = -\left\langle \frac{G^2 M^2 a (1-e^2)}{c^2 r^3} \right\rangle$ $= -\left( \frac{GM}{ca} \right)^2 \frac{1}{(1-e^2)^{3/2}}$ $\frac{\partial \langle S \rangle}{\partial e} = \left( \frac{GM}{ca} \right)^2 \frac{e}{(1-e^2)^{5/2}}$ <p>wrong because</p> $\frac{\partial \langle S \rangle}{\partial e} \neq \left\langle \nabla S \cdot \frac{\partial \underline{r}}{\partial e} \right\rangle$
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### Ron's derivation

Why this power ?

$$\frac{\partial r}{\partial e} = \frac{\partial}{\partial e} \langle r^{-3} \rangle = \frac{3 a^3 e (1-e^2)^{-5/2}}{a^3 (-3 r^{-4})}$$

$$\frac{\partial S}{\partial e} = \nabla S \cdot \frac{\partial \underline{r}}{\partial e} = \frac{\partial S}{\partial r} \frac{\partial r}{\partial e}$$

$$= \frac{-3 G^2 M^2 a (1-e^2)}{c^2 r^4} \times \frac{-e r^4}{a^3 (1-e^2)^{5/2}}$$

$$= \left( \frac{GM}{ca} \right)^2 \frac{3e}{(1-e^2)^{3/2}}$$

- put into equation for apsidal motion

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### Relativistic apsidal motion

$$\dot{w} = \frac{\sqrt{1-e^2}}{n a^2 e} \frac{\partial S}{\partial e}$$

with  $\frac{\partial S}{\partial e} = \left( \frac{GM}{ca} \right)^2 \frac{3e}{(1-e^2)^{3/2}}$

and  $n^2 a^3 = GM$

gives  $\dot{w} = \frac{6p GM}{P c^2 a (1-e^2)}$

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### Mass determinations

- combine apsidal motion and mass function to determine the total mass of the system
- measure  $\dot{w}/\dot{a}$  over many years
- measure  $f(m_2)$ , and  $a, \sin(i)$  from pulsar timing

$$\dot{w} = \frac{6p GM}{P c^2 a (1-e^2)}$$

- since  $a/M = a_1/m_2$

$$\dot{w} = \frac{6p G (m_2 \sin i)}{P c^2 (1-e^2) (a_1 \sin i)}$$

- solve for  $m_2 \sin(i)$ , then use  $f(m_2) = m_2^3 \sin^3(i) / M^2$

$$\therefore M = \left( \frac{m_2 \sin i}{f(m_2)} \right)^{1/2}$$

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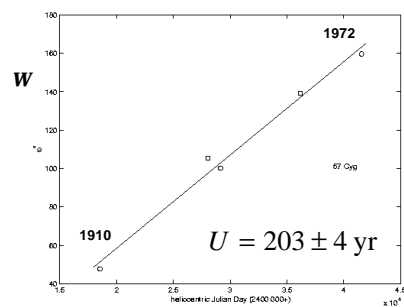
### Aspherical stars and apsidal motion

- For non-spherical stars
  - potential not that of 2 point masses
- observe  $w(t)$  from radial velocity curves
- or from eclipses
  - changing time of eclipse minima
- Relate apsidal period to orbital period,  $U/P$ 
  - see handout
  - measureable for  $U \sim 100$ s of years
- tell us about internal structure of stars
  - need gravitational potential of non-spherical bodies

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### 57 Cyg - radial velocity orbit



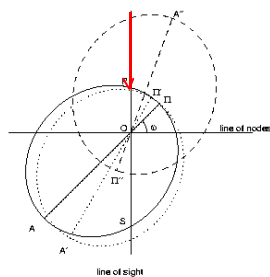
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### eclipse times for eccentric orbit

eclipses (for  $i = 90^\circ$ ) at

$$q + w = 90^\circ, 270^\circ$$



eclipses at minima of

$$d(q) = \frac{1-e^2}{1+e \cos q} \sqrt{1 - \sin^2 i \sin^2(q+w)}$$

minimum projected separation between stars

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### eclipse times with apsidal motion

advance of periastron :

$$w(t) = w_0 + \dot{w} (t - t_0)$$

eclipse times :

$$t_n - t_0 = n P + \frac{P}{2p} (w_0 - w_n) + \frac{eP}{p} (\cos w_0 - \cos w_n) + \dots$$

apparent period :

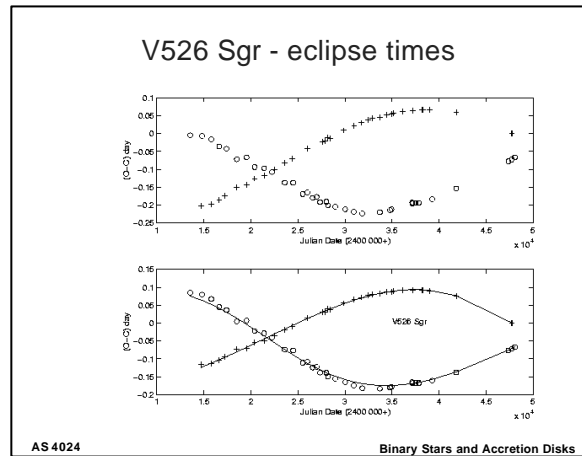
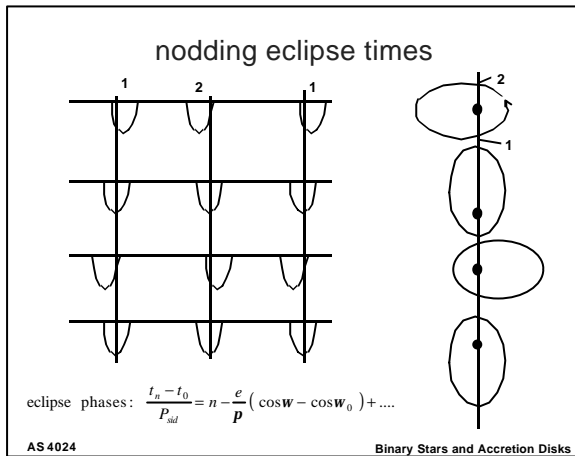
$$P_{sid} \rightarrow P \left( 1 - \frac{\dot{w}}{2p} \right)$$

eclipse phases :

$$\frac{t_n - t_0}{P_{sid}} = n - \frac{e}{p} (\cos w - \cos w_0) + \dots$$

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### Internal Structure of Stars

$U = \frac{2p}{w}$  = apsidal period

$$\frac{P}{U} = k_1 \left( \frac{R_1}{a} \right)^5 \left( \frac{m_2}{m_1} \{ 15f(e) + g(e) \} + g(e) \right) + k_2 \left( \frac{R_2}{a} \right)^5 \left( \frac{m_1}{m_2} \{ 15f(e) + g(e) \} + g(e) \right)$$

$$f(e) = 1 + \frac{13}{2}e^2 + \frac{181}{8}e^4 + \dots$$

$$g(e) = (1 - e^2)^{-2}$$

$R_1, R_2$  = radii of stars  
 $k_1, k_2$  = apsidal constants - internal structure - of stars

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