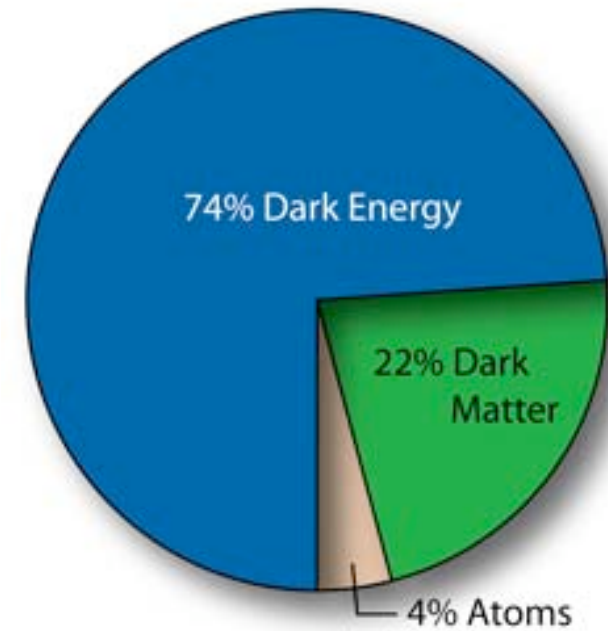


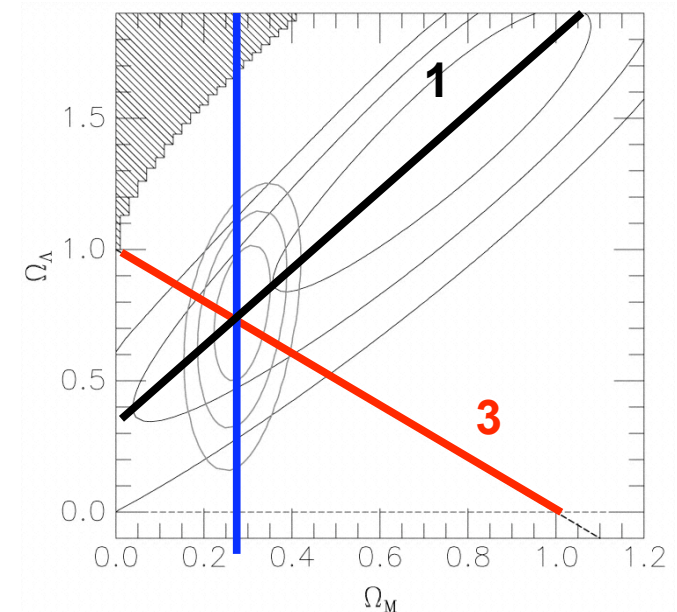
Cosmic Microwave Background

Flat Geometry

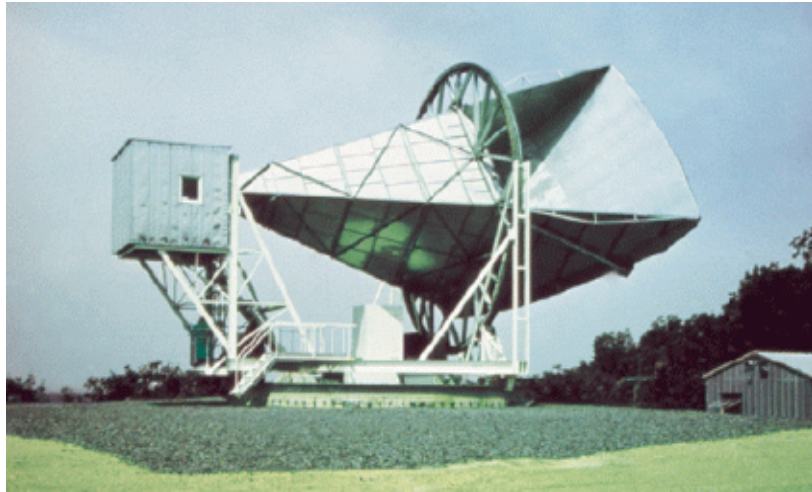
$$\Omega_0 = \Omega_M + \Omega_\Lambda$$
$$\approx 1.0$$



2



1965 -- Penzias + Wilson



Bell Labs telecommunications engineers find excess microwave noise from the sky.

~1% of thermal ($T \sim 300^\circ \text{K}$) noise $\rightarrow T \sim 3^\circ \text{K}$

Afterglow of the Big Bang

CMB = **C**osmic **M**icrowave **B**ackground

Confirms a forgotten 1948 prediction by Gamow.

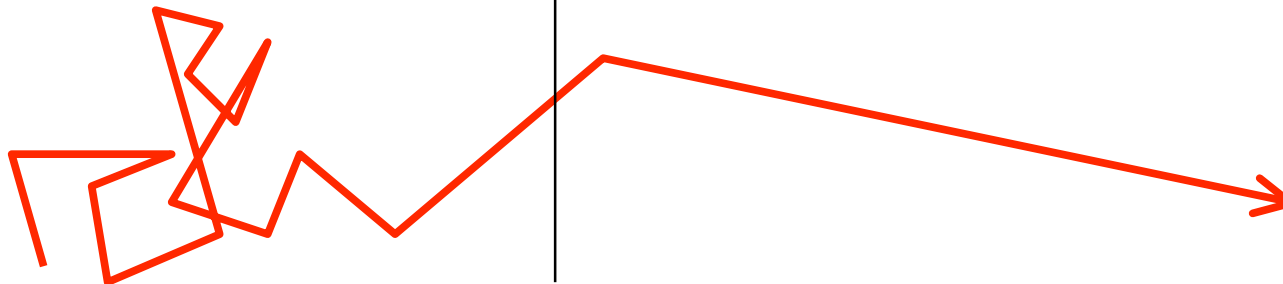
Nobel Prize \rightarrow P+W

Recombination Epoch ($z \sim 1100$)

ionised plasma \rightarrow *neutral gas*

- Redshift $z > 1100$
- Temp $T > 3000$ K
- H ionised
- electron -- photon
Thompson scattering

- $z < 1100$
- $T < 3000$ K
- H recombined
- almost no electrons
- neutral atoms
- photons set free

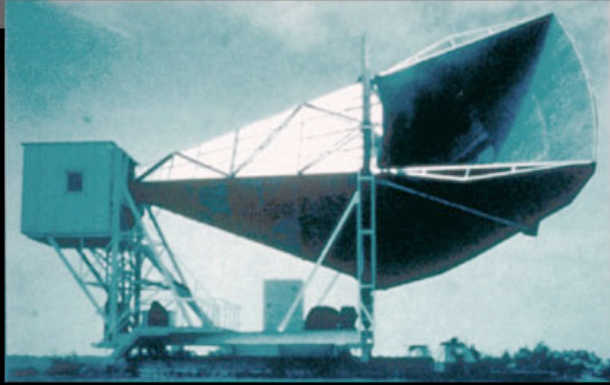


e - scattering optical depth

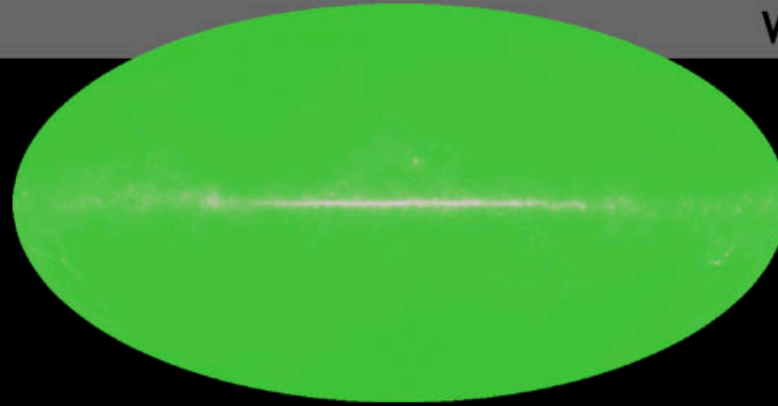
$$\tau(z) \approx \left(\frac{z}{1080} \right)^{13}$$

thin surface of last scattering

1965



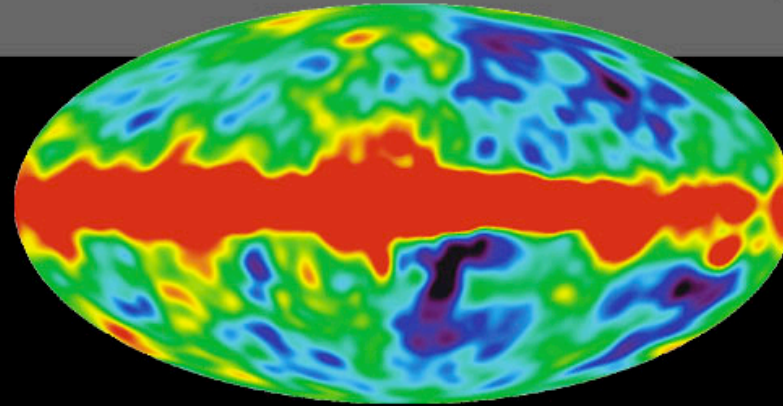
Penzias and
Wilson



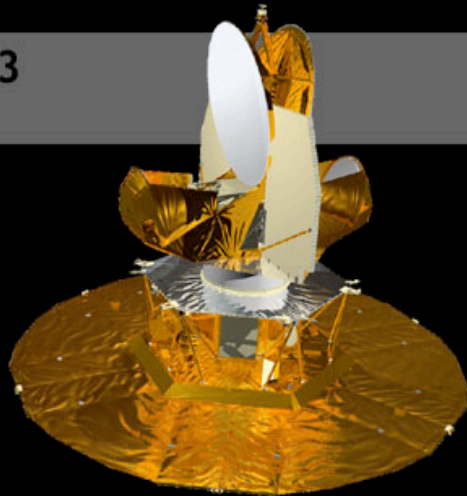
1992



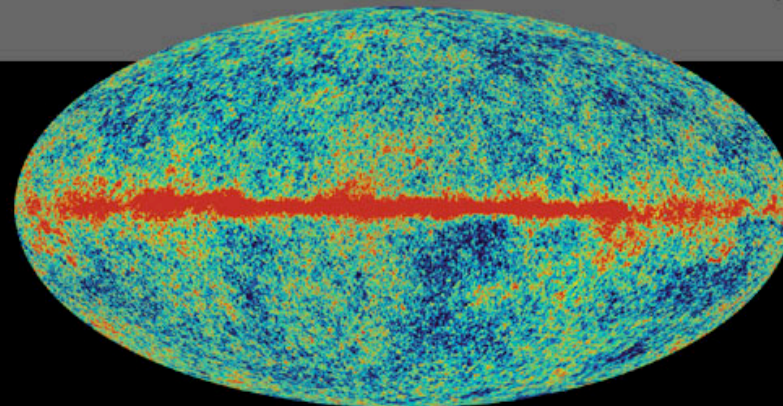
COBE



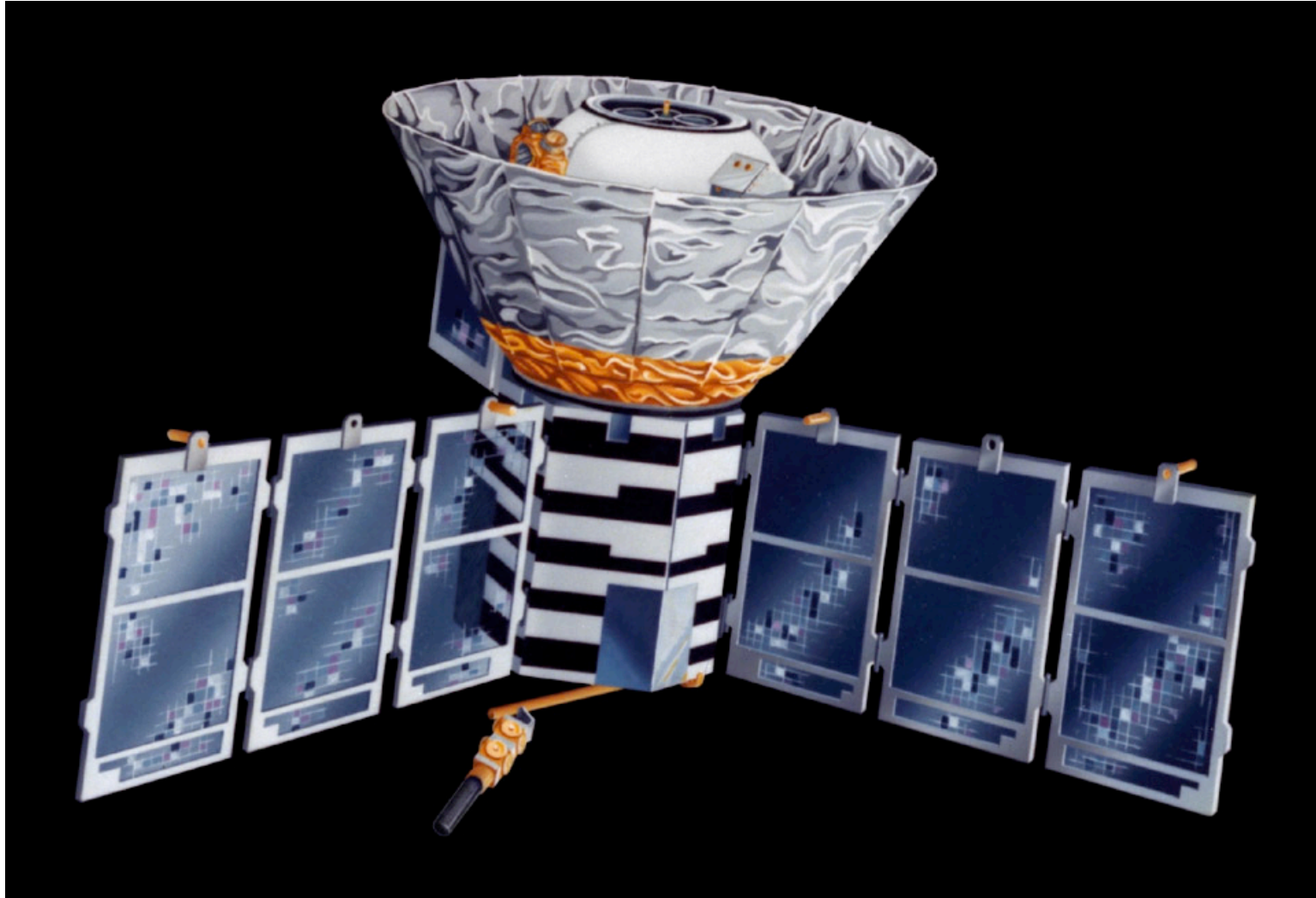
2003



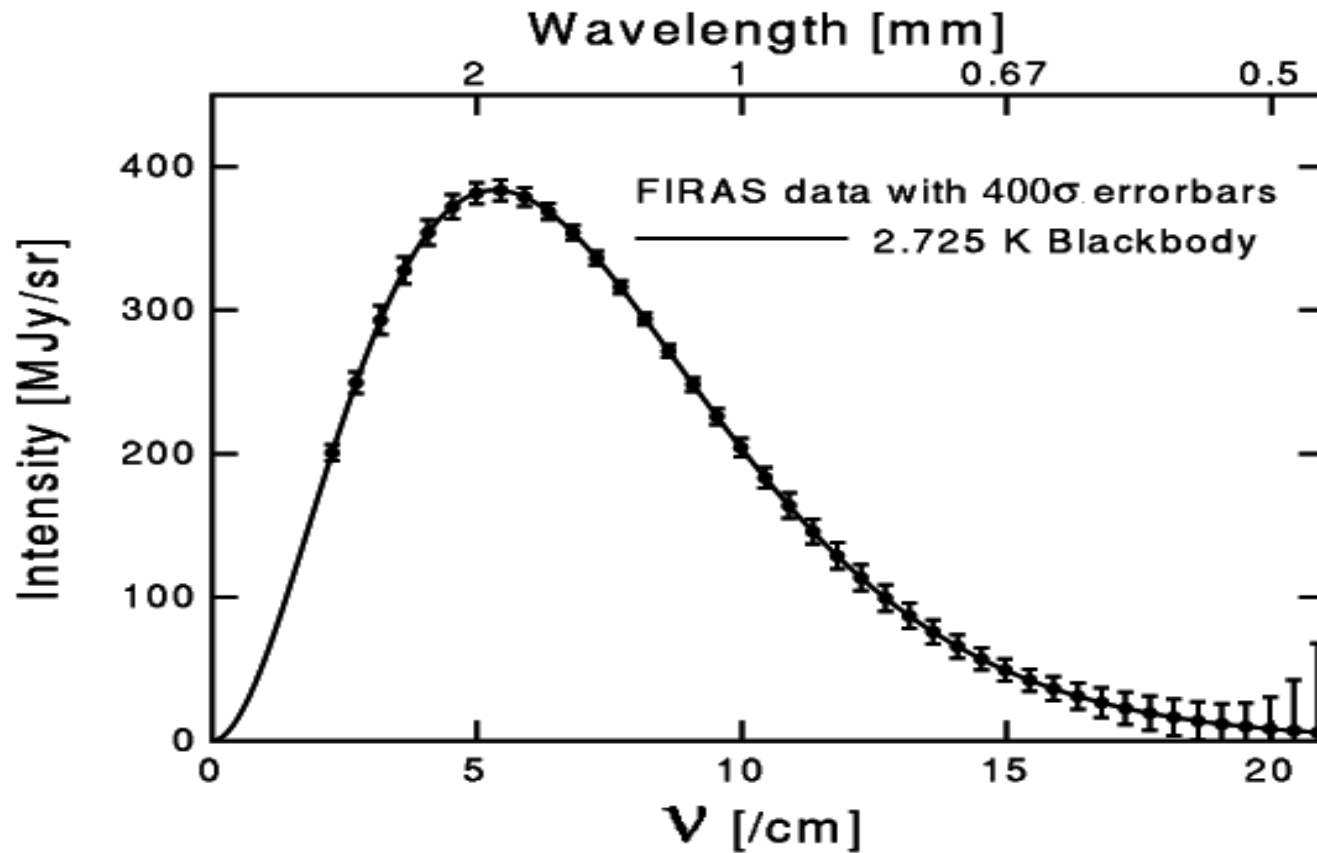
WMAP



NASA 1992 - **COBE** **CO**smic **B**ackground **E**xplorer



COBE spectrum of CMB

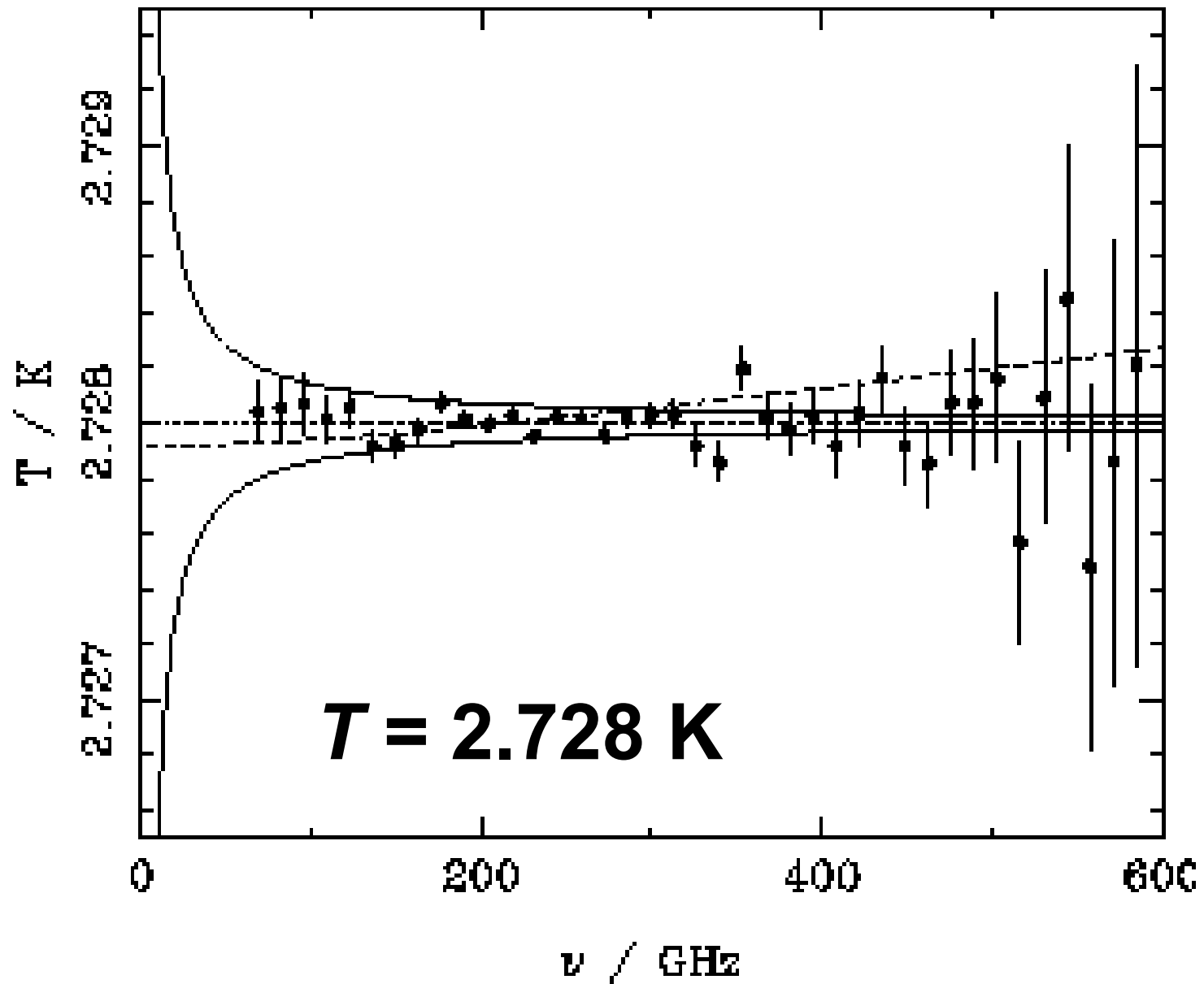


A perfect Blackbody !

**No spectral lines -- strong test of Big Bang.
Expansion preserves the blackbody spectrum.**

$$T(z) = T_0 (1+z) \quad T_0 \sim 3000 \text{ K} \quad z \sim 1100$$

COBE spectrum



Radiation -> Matter -> Vacuum

$$T = 2.728 \text{ K}$$

radiation energy density :

$$\rho_R = \frac{u(T)}{c^2} = \frac{4 \sigma}{c^3} T^4$$

$$\Omega_R = 8.6 \times 10^{-5} \left(\frac{0.7}{h} \right)^2 \sim 0.01\%$$

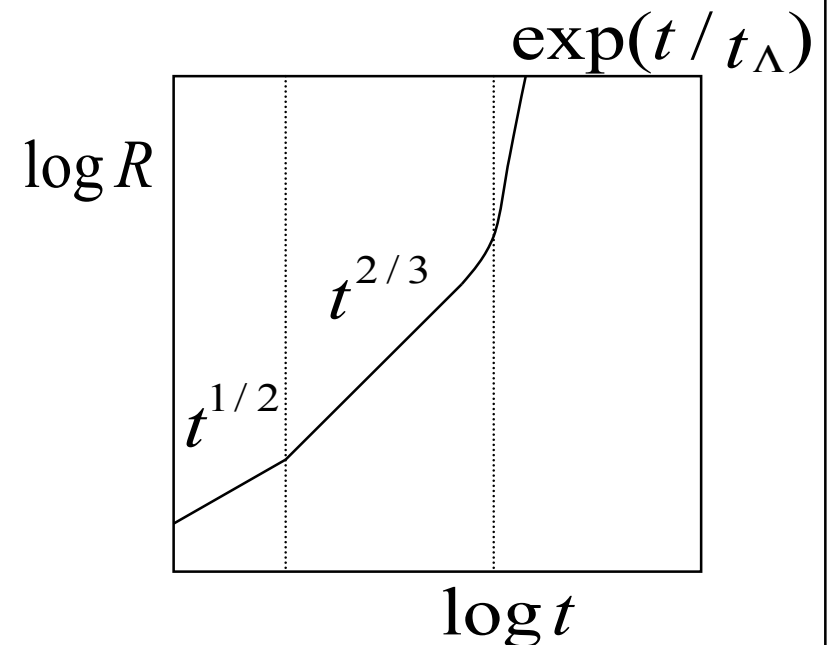
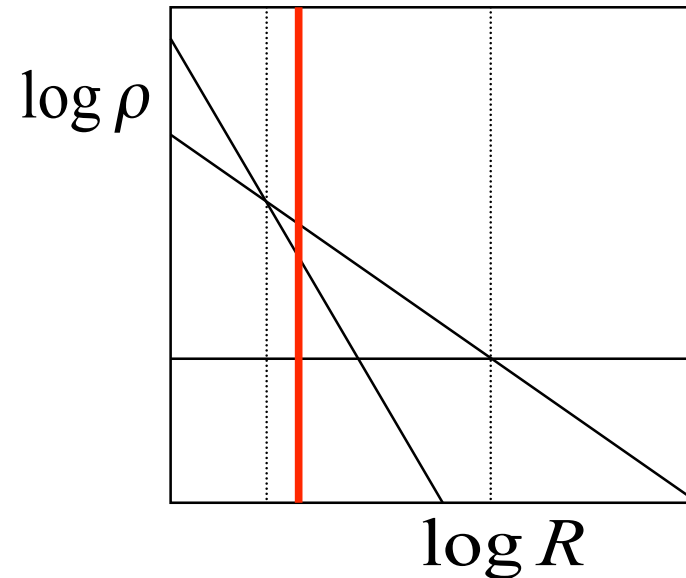
(including neutrinos)

matter - radiation equality :

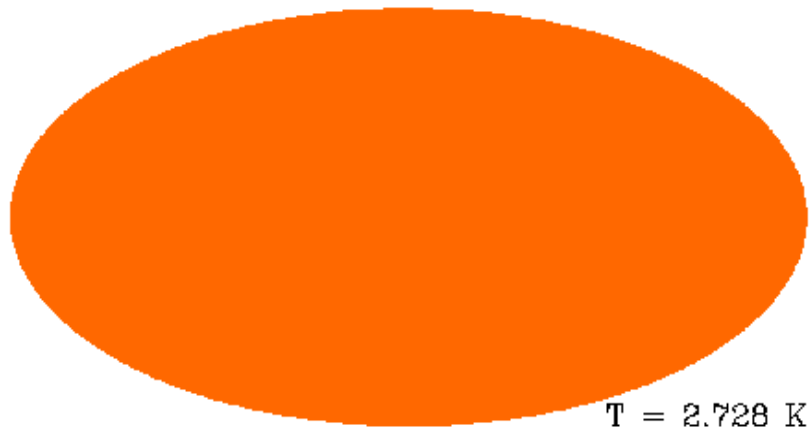
$$\frac{\Omega_R(z)}{\Omega_M(z)} = \frac{(1+z)^4 \Omega_R}{(1+z)^3 \Omega_M} = 1$$

$$(1+z_{eq}) = \frac{\Omega_M}{\Omega_R} = 3500 \left(\frac{\Omega_M}{0.3} \right) \left(\frac{2.73 \text{ K}}{T} \right)^4$$

Matter-dominated at z=1100

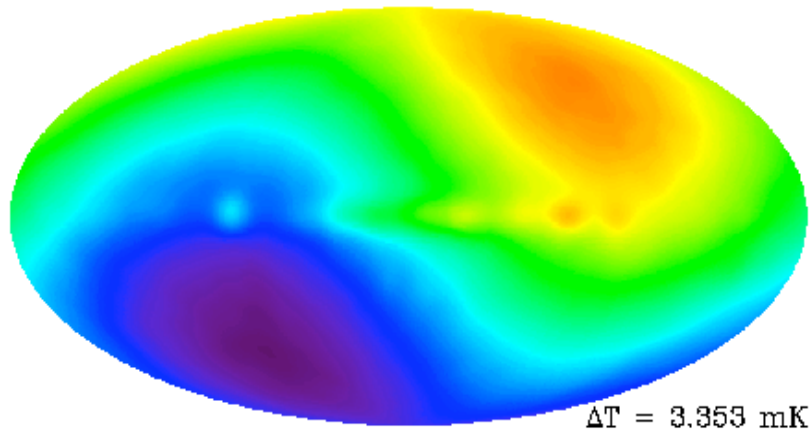


Cosmic Microwave Background



Almost isotropic

$$T = 2.728 \text{ K}$$

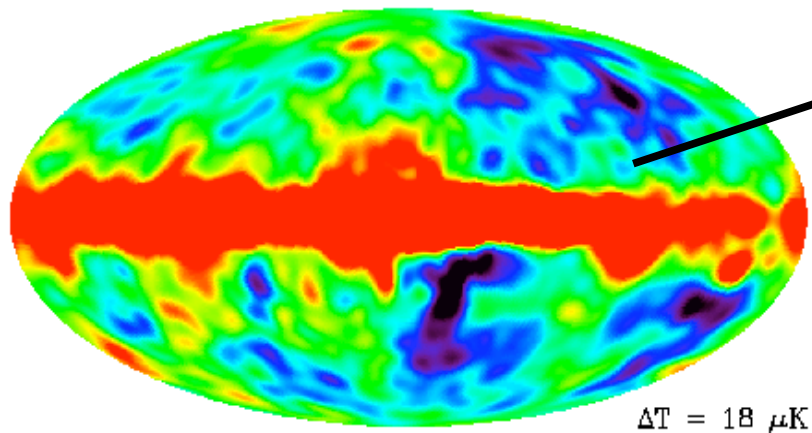


Dipole anisotropy

$$\frac{V}{c} = \frac{\Delta\lambda}{\lambda} = \frac{\Delta T}{T} \approx 10^{-3}$$

Our velocity:

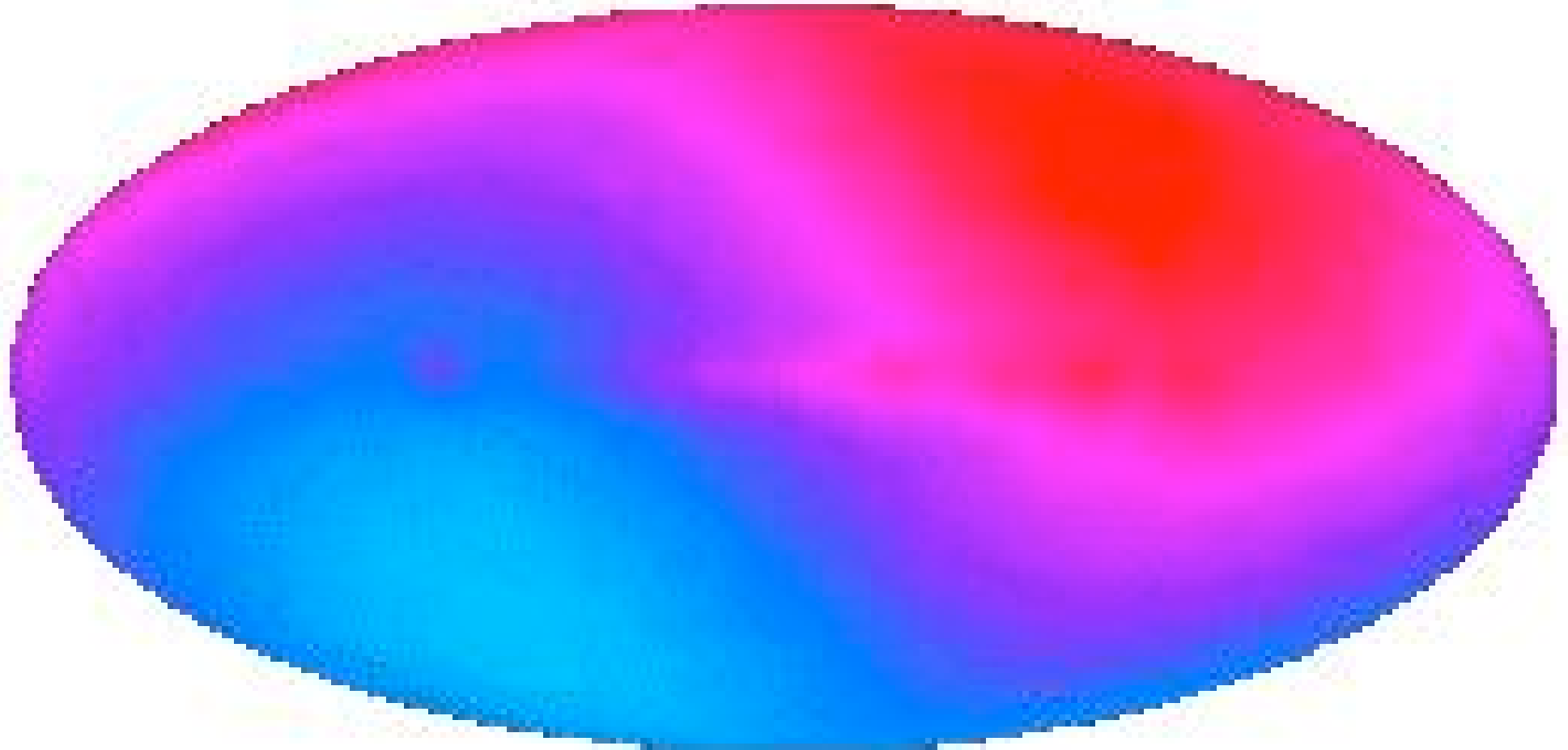
$$V \approx 600 \text{ km/s}$$



Milky Way sources

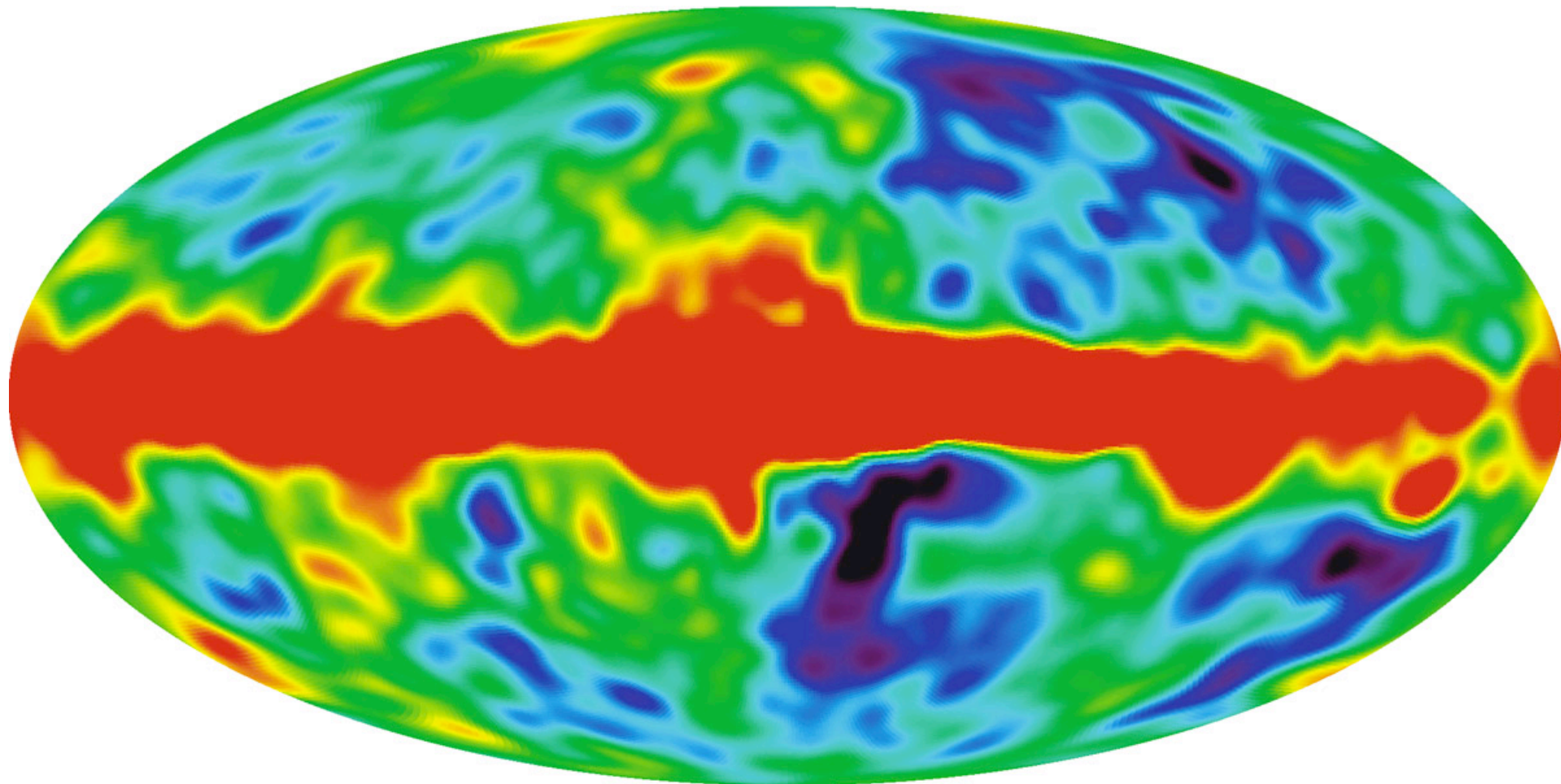
$$+ \text{anisotropies } \frac{\Delta T}{T} \sim 10^{-5}$$

0.1% CMB dipole anisotropy velocity relative to CMB photons



$$T(\theta) = T_0 \left(1 + \frac{V}{c} \cos \theta + \dots \right) \rightarrow \begin{array}{l} V_{SUN} = 371 \pm 1 \text{ km s}^{-1} \\ V_{MW} \approx 600 \text{ km s}^{-1} \end{array}$$

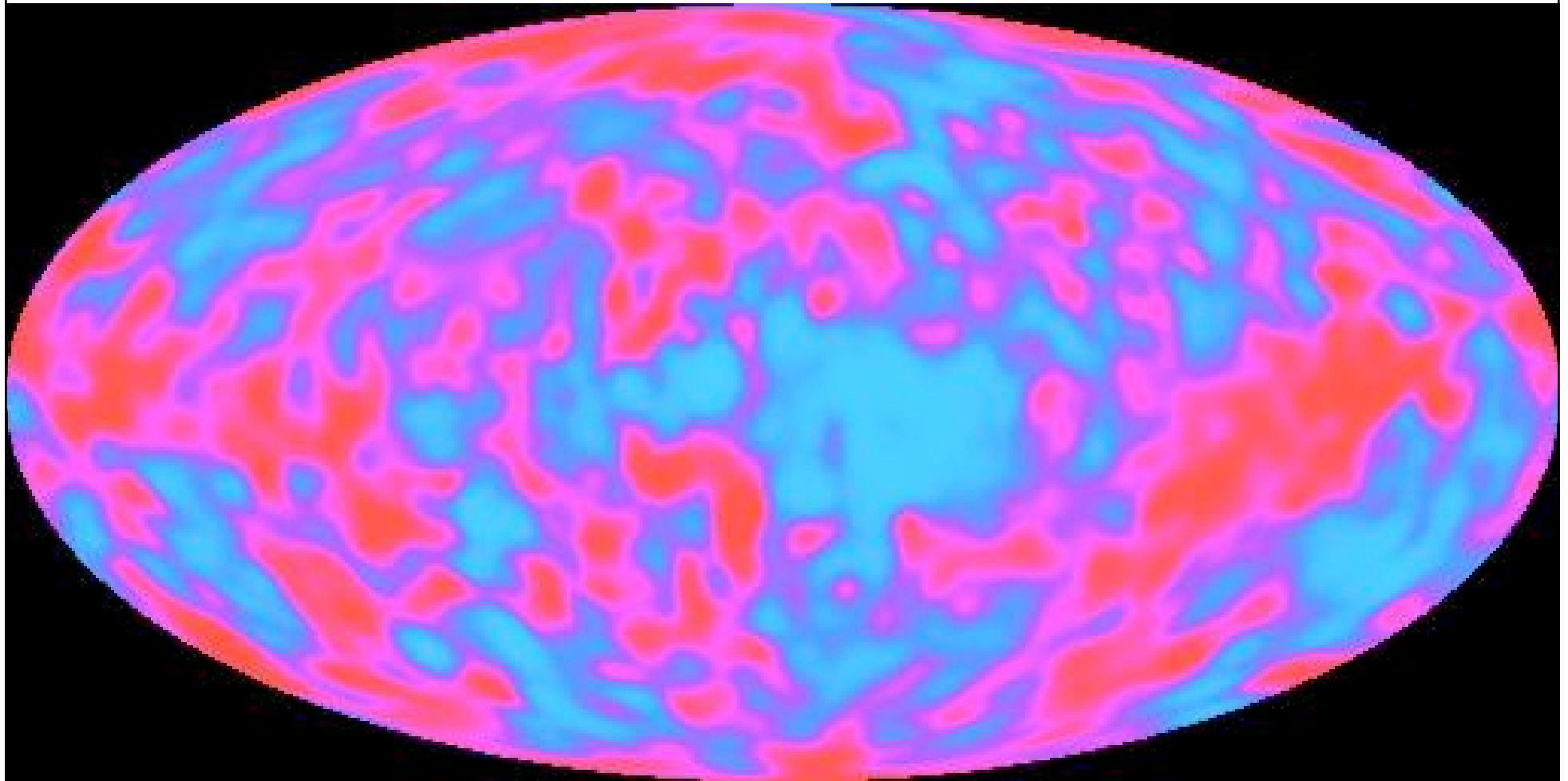
COBE 4-year map



Milky Way emission

subtract by using maps at several frequencies

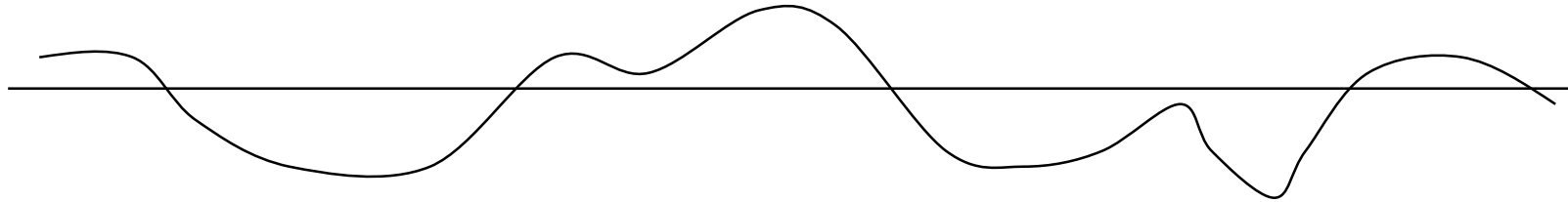
COBE - tiny ripples



$$\frac{\delta T}{T} \sim 10^{-5}$$

Resolution $\sim 7^\circ$

Tiny Ripples at Redshift 1100



$$\frac{\Delta T}{T} \approx \frac{\Delta \rho}{4 \rho} \sim 10^{-5} \text{ at } z = 1100$$

Ripples are :

relics of the Big Bang

initial quantum fluctuations expanded by early inflation

the seeds of later galaxy/cluster formation.

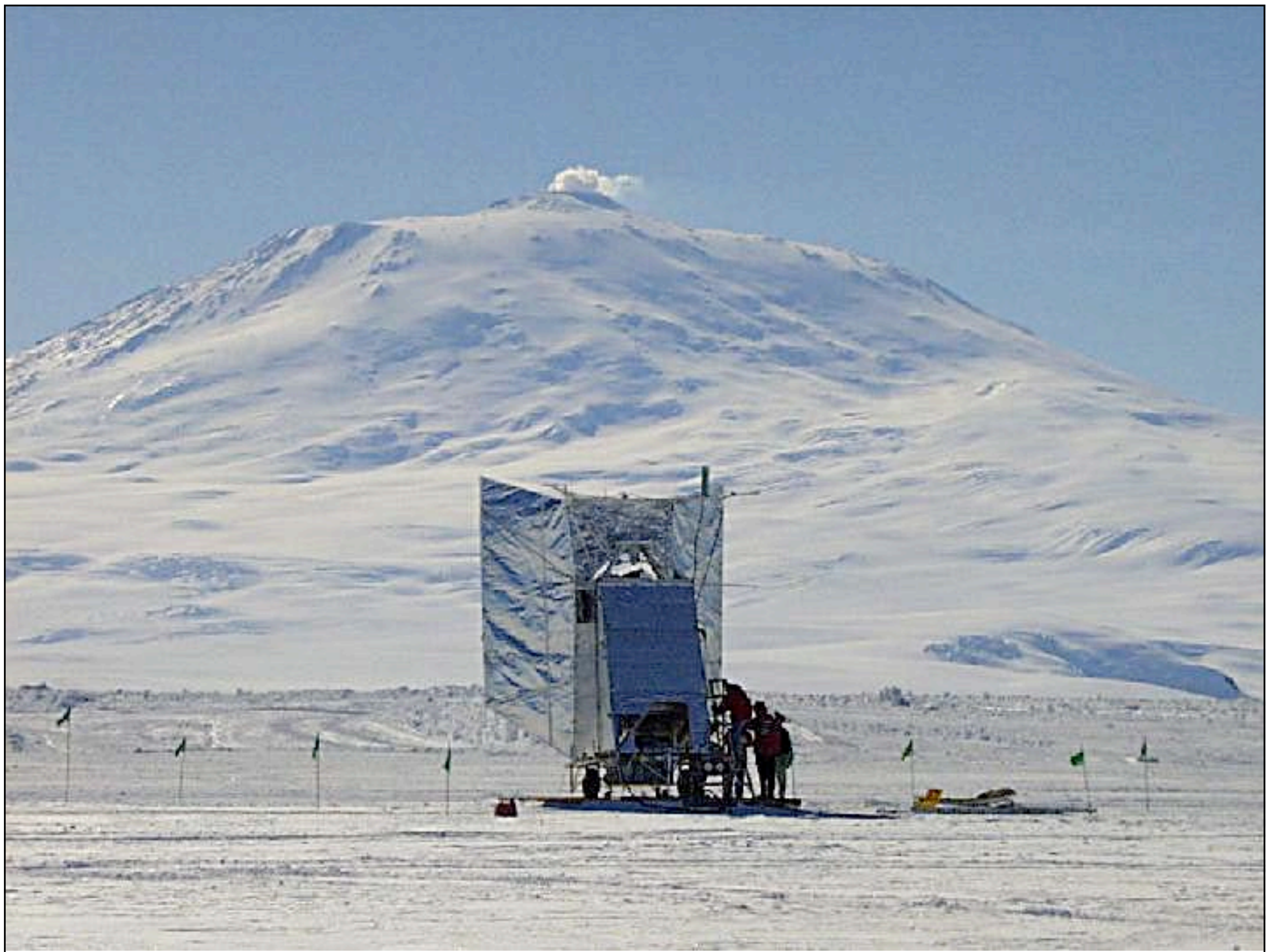
standard yardsticks for measuring curvature

(and other cosmology parameters)

1999 - Boomerang in Antarctica

Balloon Observations Of
Millimetric Extragalactic
Radiation Anisotropy
and Geophysics



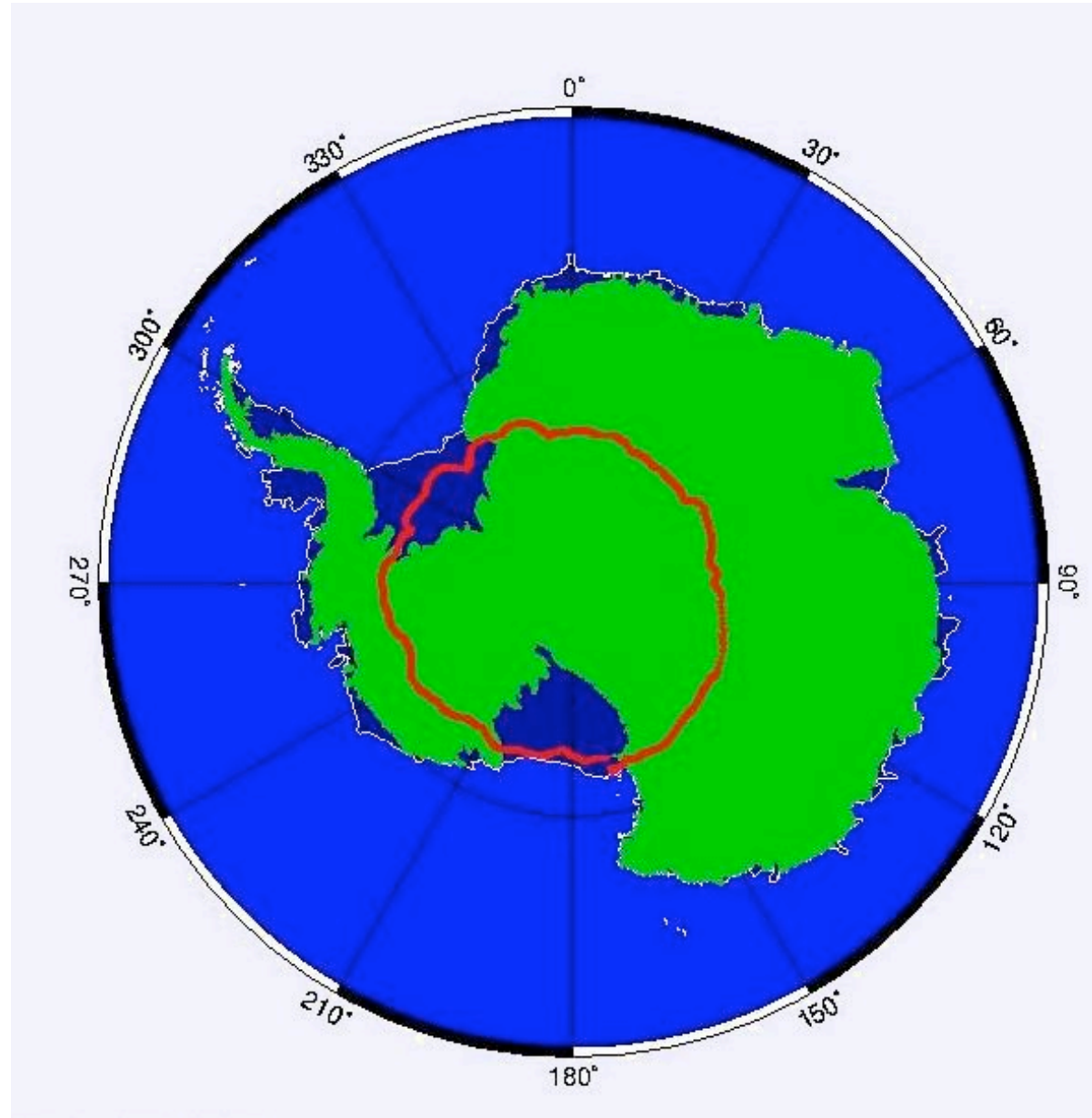


AS 4022 Cosmology

Boomerang's Balloon

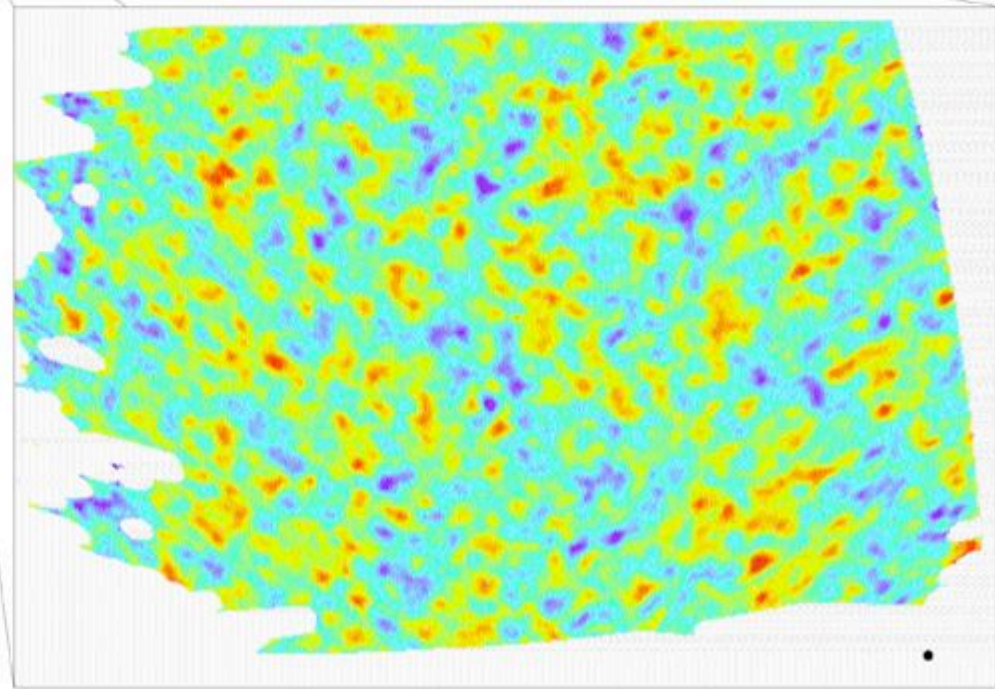
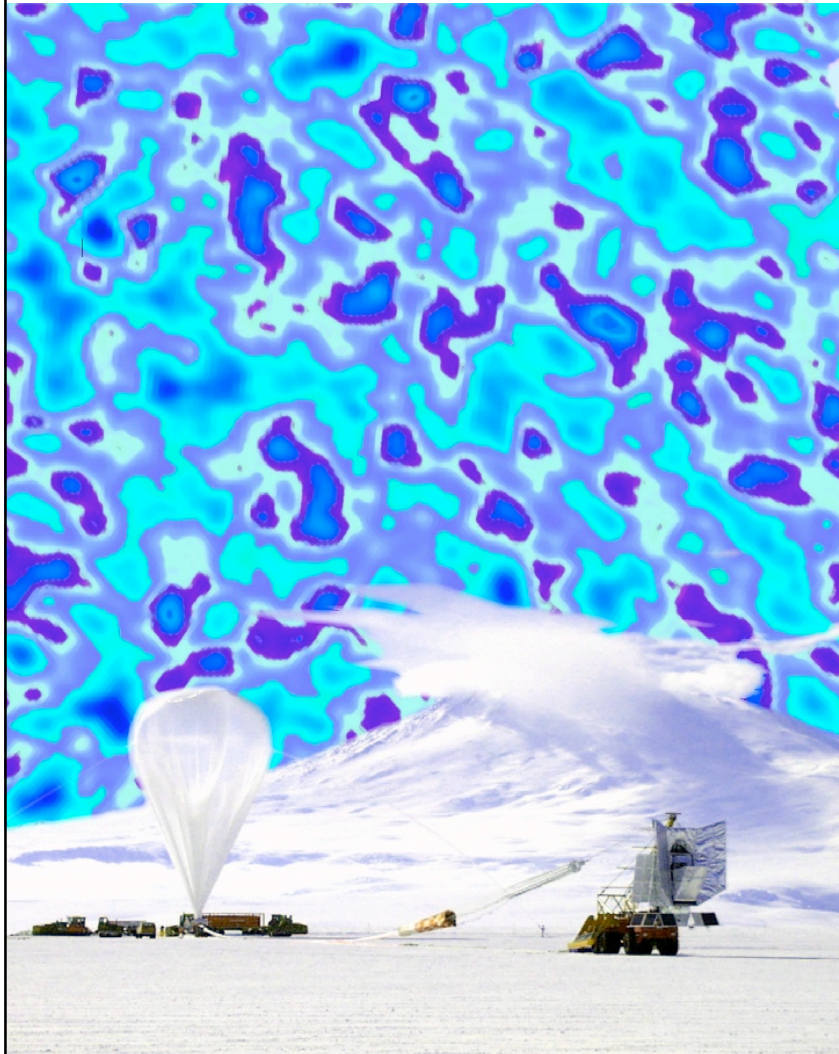
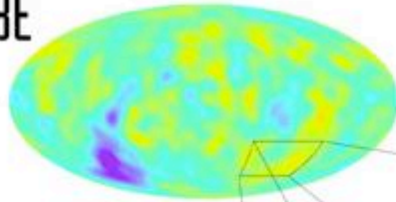


Boomerang's Stratospheric Flight Track



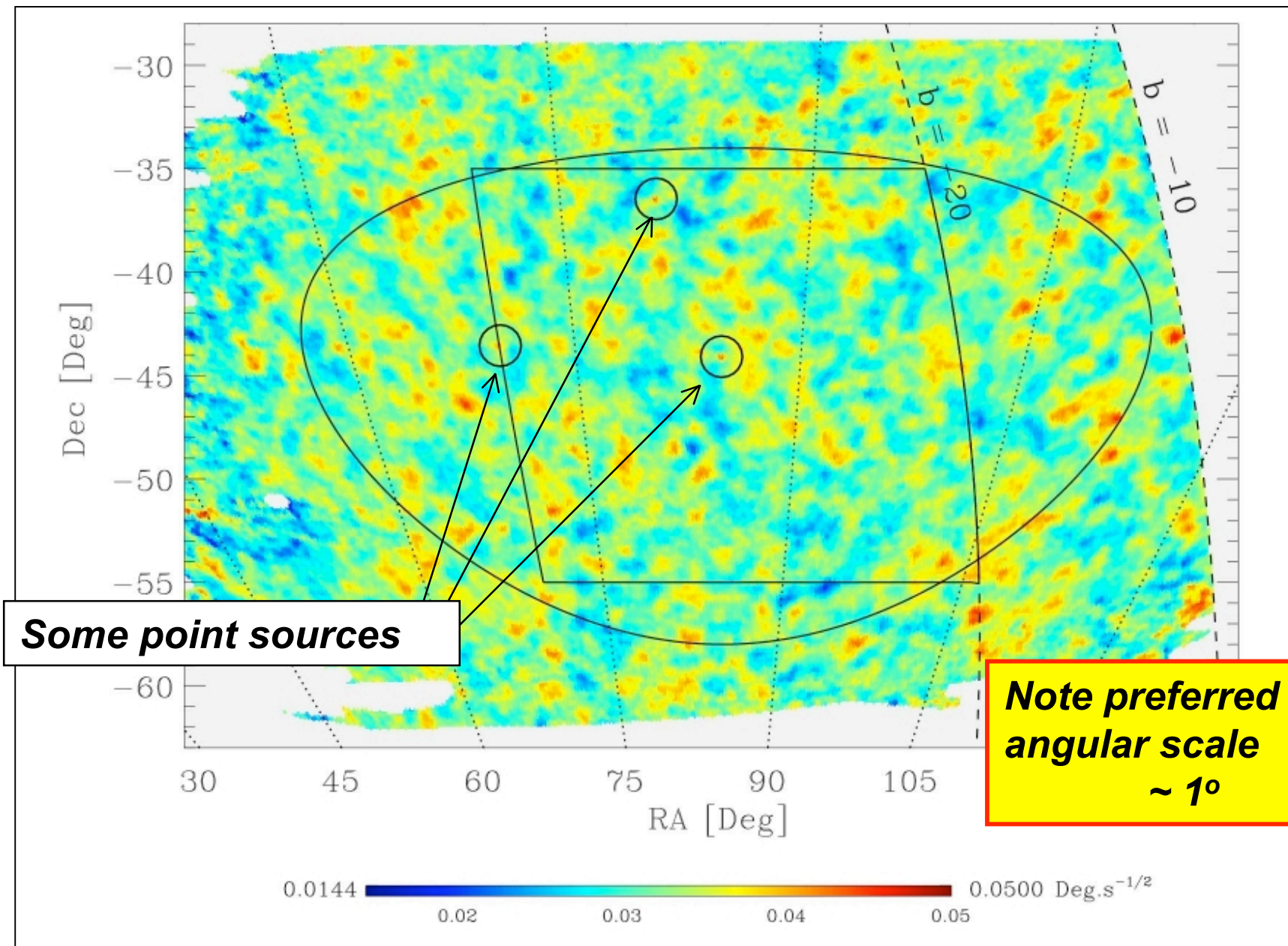
Altitude
37 km
10 days

COBE



Resolution $\sim 0.3^\circ$

Boomerang Map



Spherical Harmonics

Fit temperature map with a series of spherical harmonics

$$\frac{\Delta T}{T}(\theta, \phi) = \sum_{l,m} a_{lm} Y_{lm}(\theta, \phi)$$

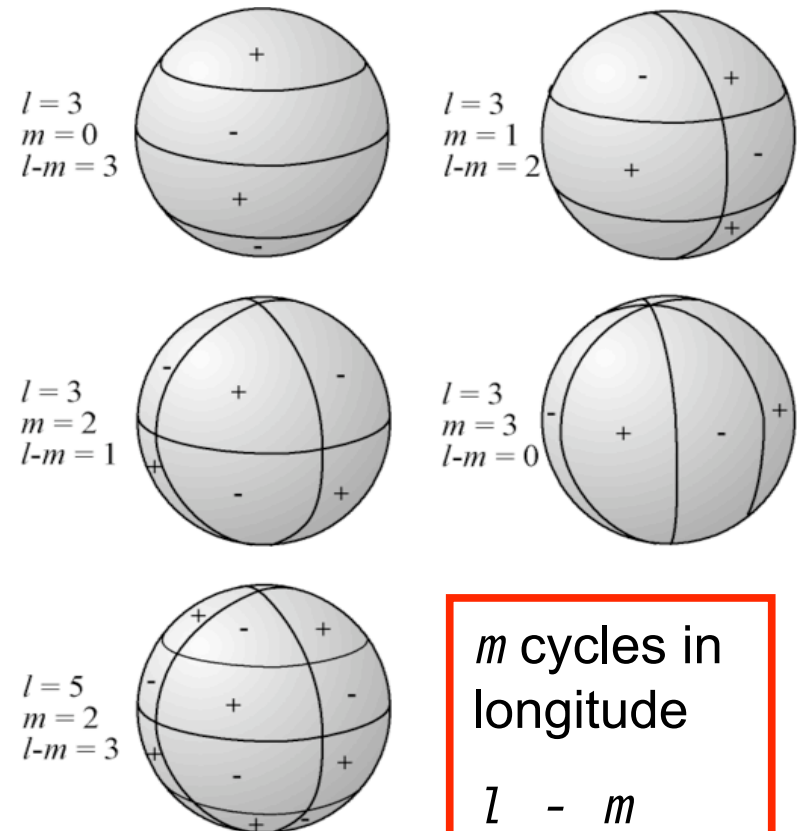
angular power spectrum

$$C_l = \left\langle |a_{lm}|^2 \right\rangle \text{ average } -l \leq m \leq l$$

dimensionless power spectrum

$$l(l+1) C_l \propto \frac{d \left\langle (\Delta T / T)^2 \right\rangle}{d \ln l}$$

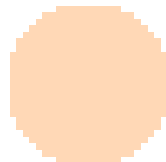
$$\text{angular scale: } \Delta\theta \approx \frac{\pi}{l} = \frac{180^\circ}{l}$$



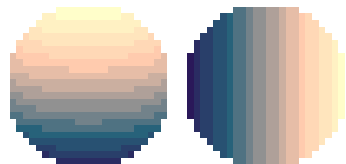
Spherical Harmonics

$$Y_0^0(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{1}{\pi}}$$

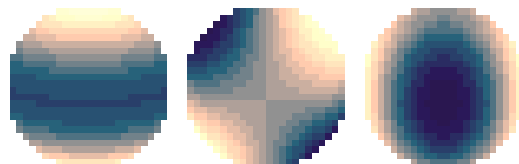
$l = 0$



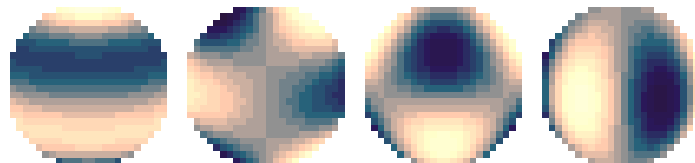
1



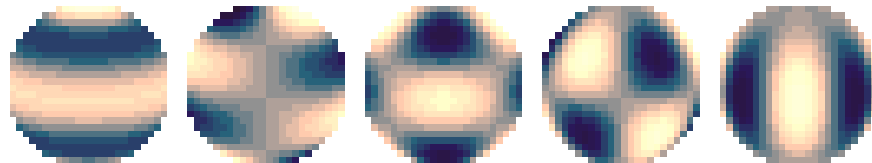
2



3



4



$m = 0$

1

2

3

4

$$Y_1^{-1}(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{3}{2\pi}} \sin \theta e^{-i\varphi}$$

$$Y_1^0(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{3}{\pi}} \cos \theta$$

$$Y_1^1(\theta, \varphi) = \frac{-1}{2} \sqrt{\frac{3}{2\pi}} \sin \theta e^{i\varphi}$$

$$Y_2^{-2}(\theta, \varphi) = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{-2i\varphi}$$

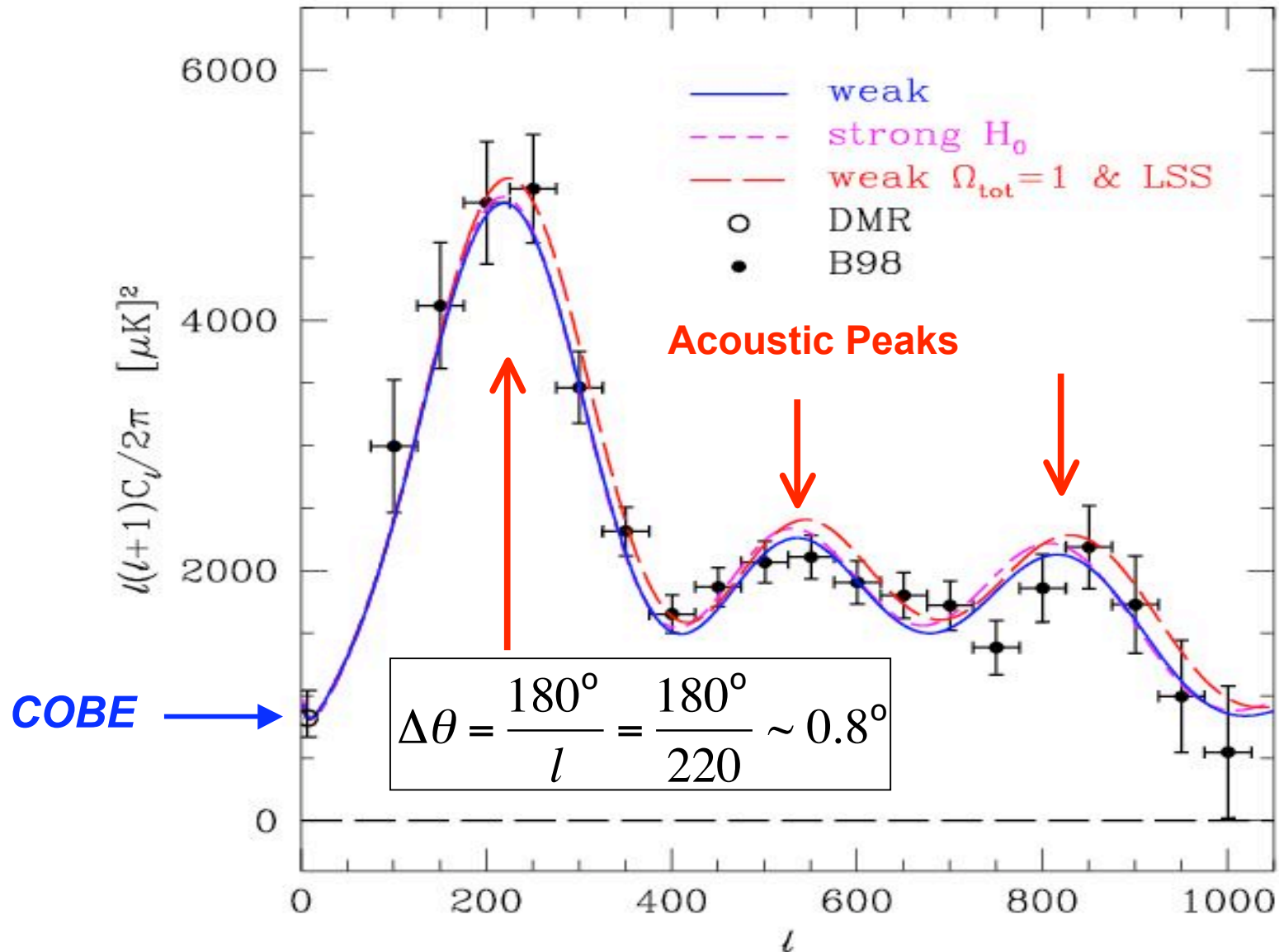
$$Y_2^{-1}(\theta, \varphi) = \frac{1}{2} \sqrt{\frac{15}{2\pi}} \sin \theta \cos \theta e^{-i\varphi}$$

$$Y_2^0(\theta, \varphi) = \frac{1}{4} \sqrt{\frac{5}{\pi}} (3 \cos^2 \theta - 1)$$

$$Y_2^1(\theta, \varphi) = \frac{-1}{2} \sqrt{\frac{15}{2\pi}} \sin \theta \cos \theta e^{i\varphi}$$

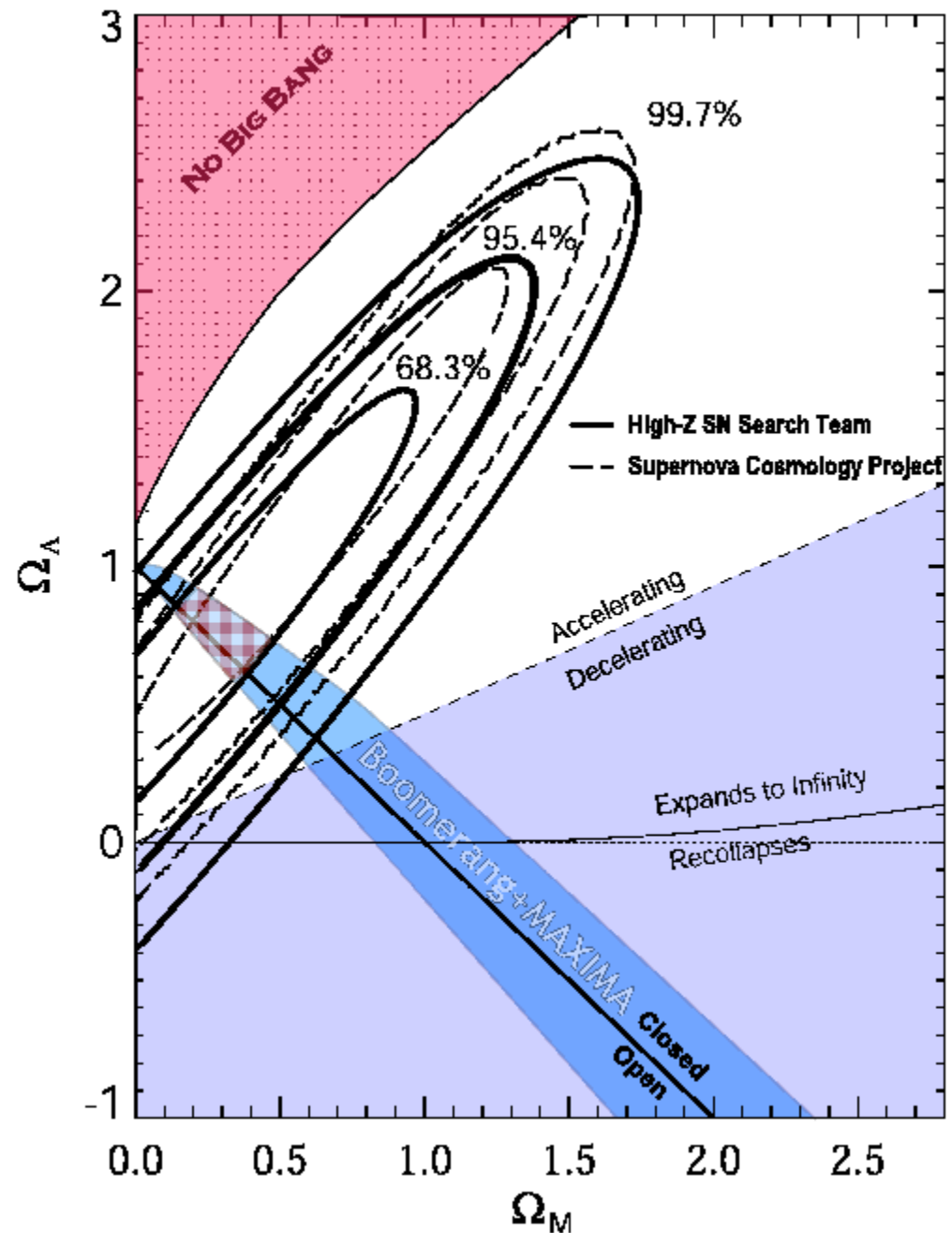
$$Y_2^2(\theta, \varphi) = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\varphi}$$

Boomerang Power Spectrum



Supernovae + CMB ripples

Pre-WMAP
constraints
From BOOMERANG
and MAXIMA
circa 2002



WMAP

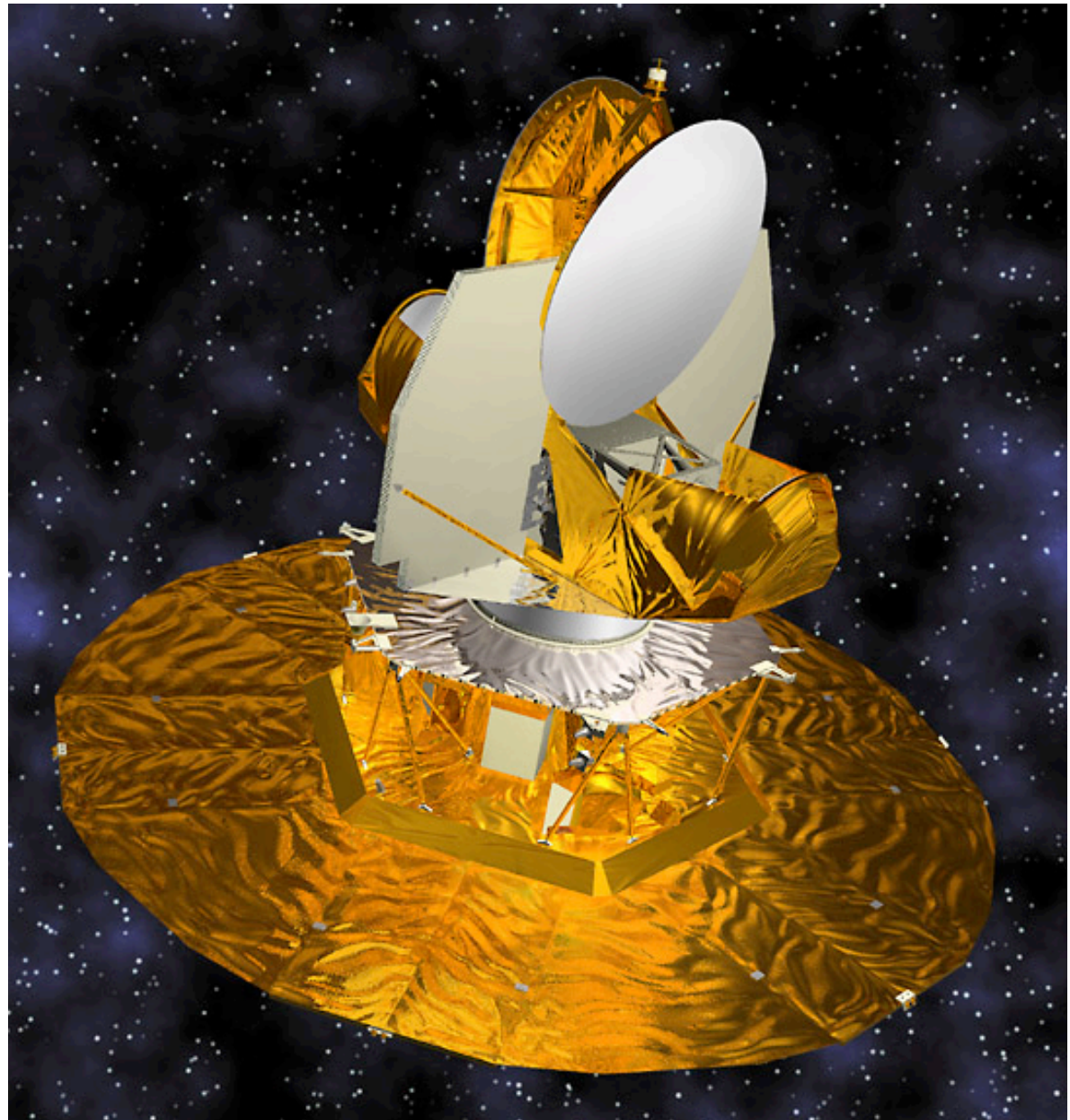
NASA 2001...

Wilkinson

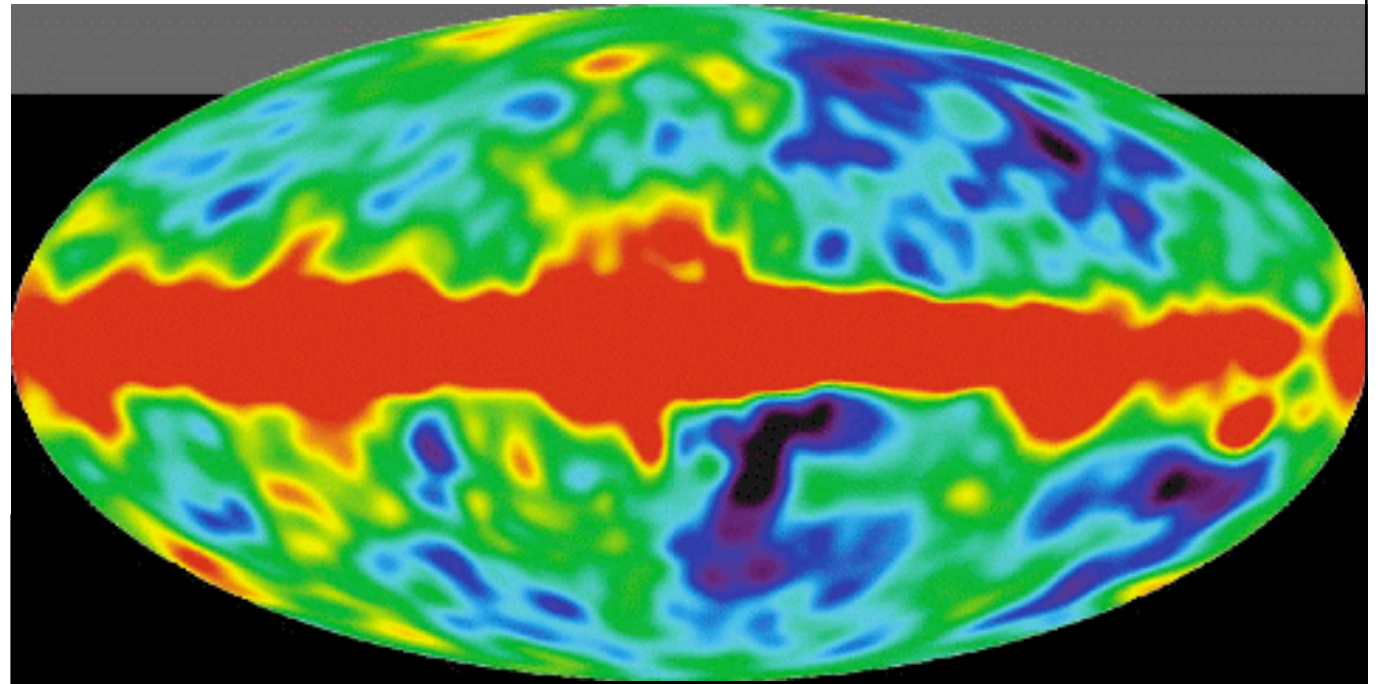
Microwave

Anisotropy

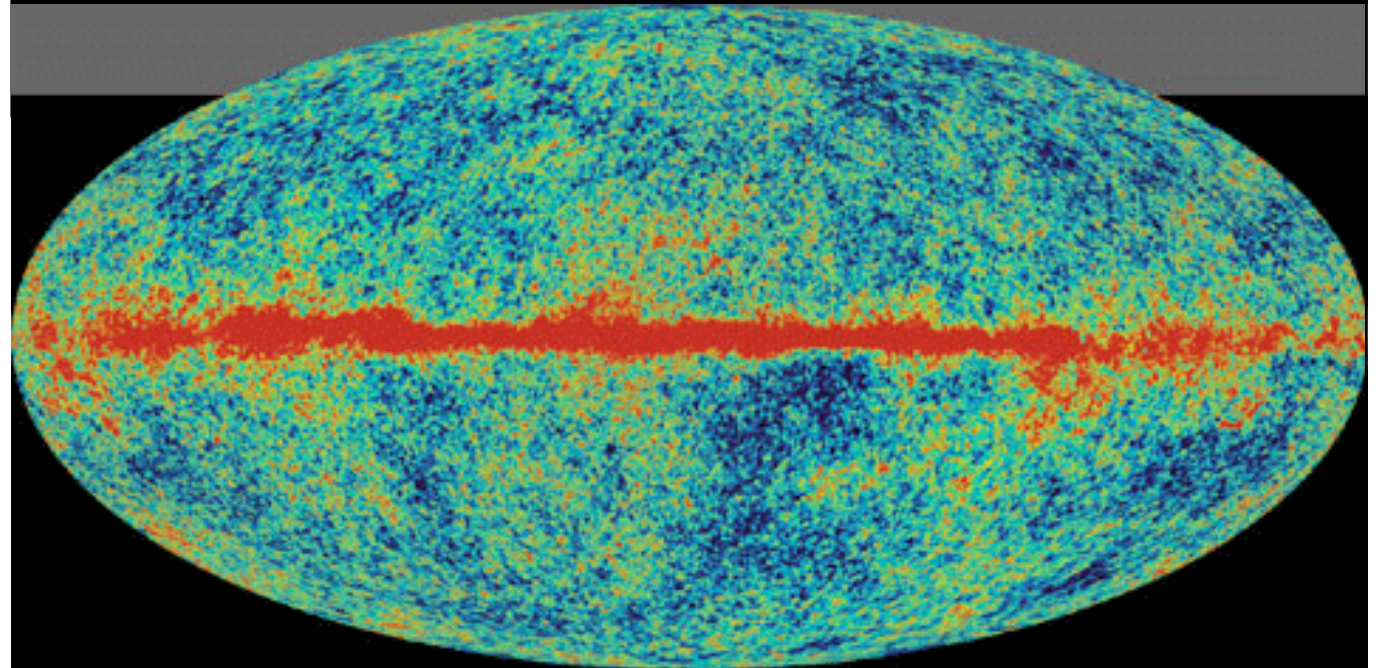
Probe



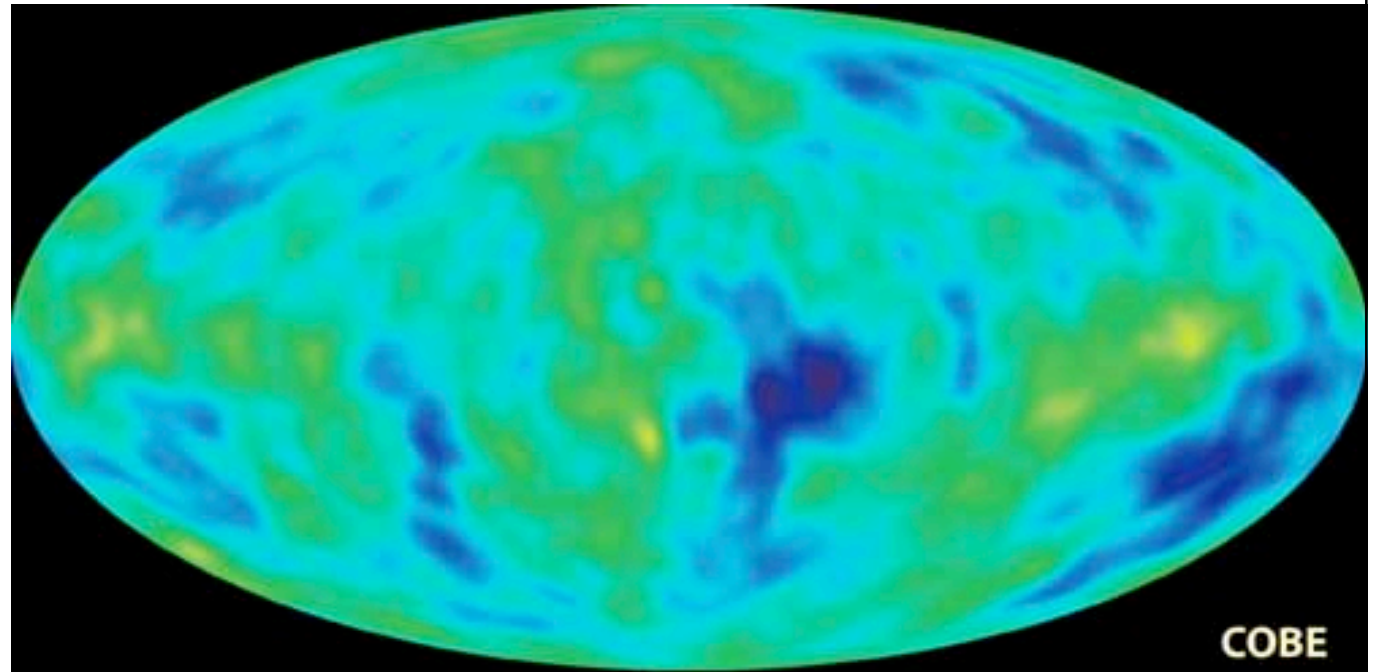
COBE
1992
7 degree
resolution



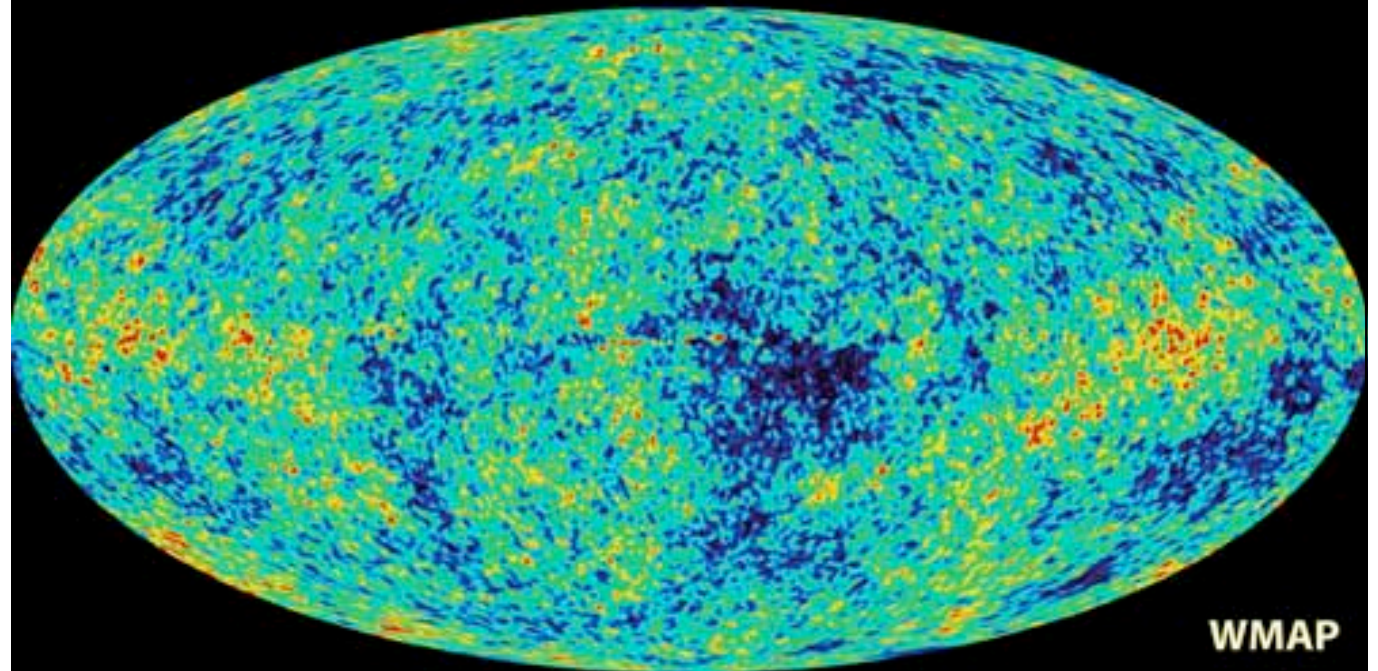
WMAP
2003
20 arcmin
resolution



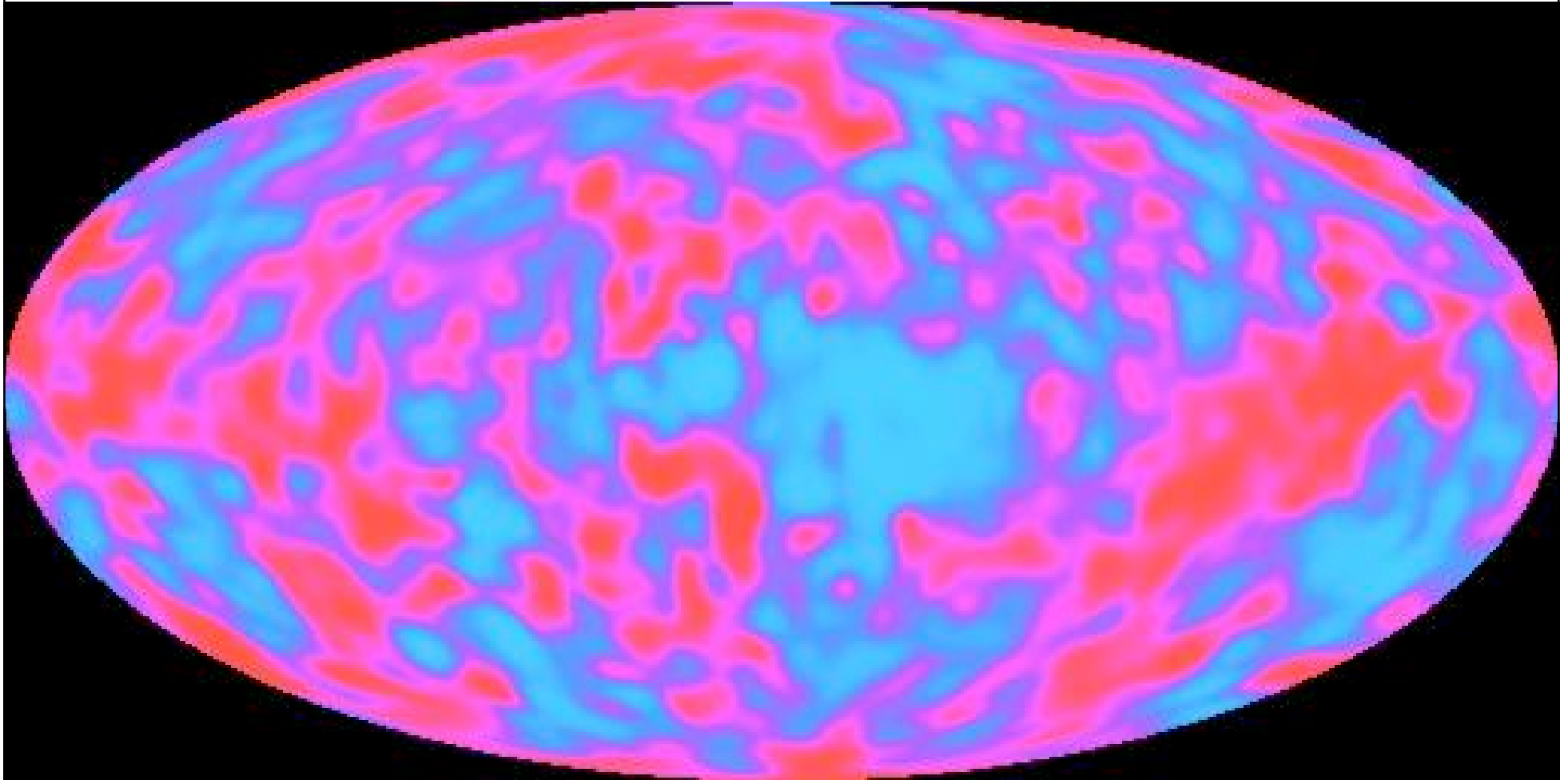
COBE
1992
7 degree
resolution



WMAP
2003
20 arcmin
resolution



COBE - temperature ripples

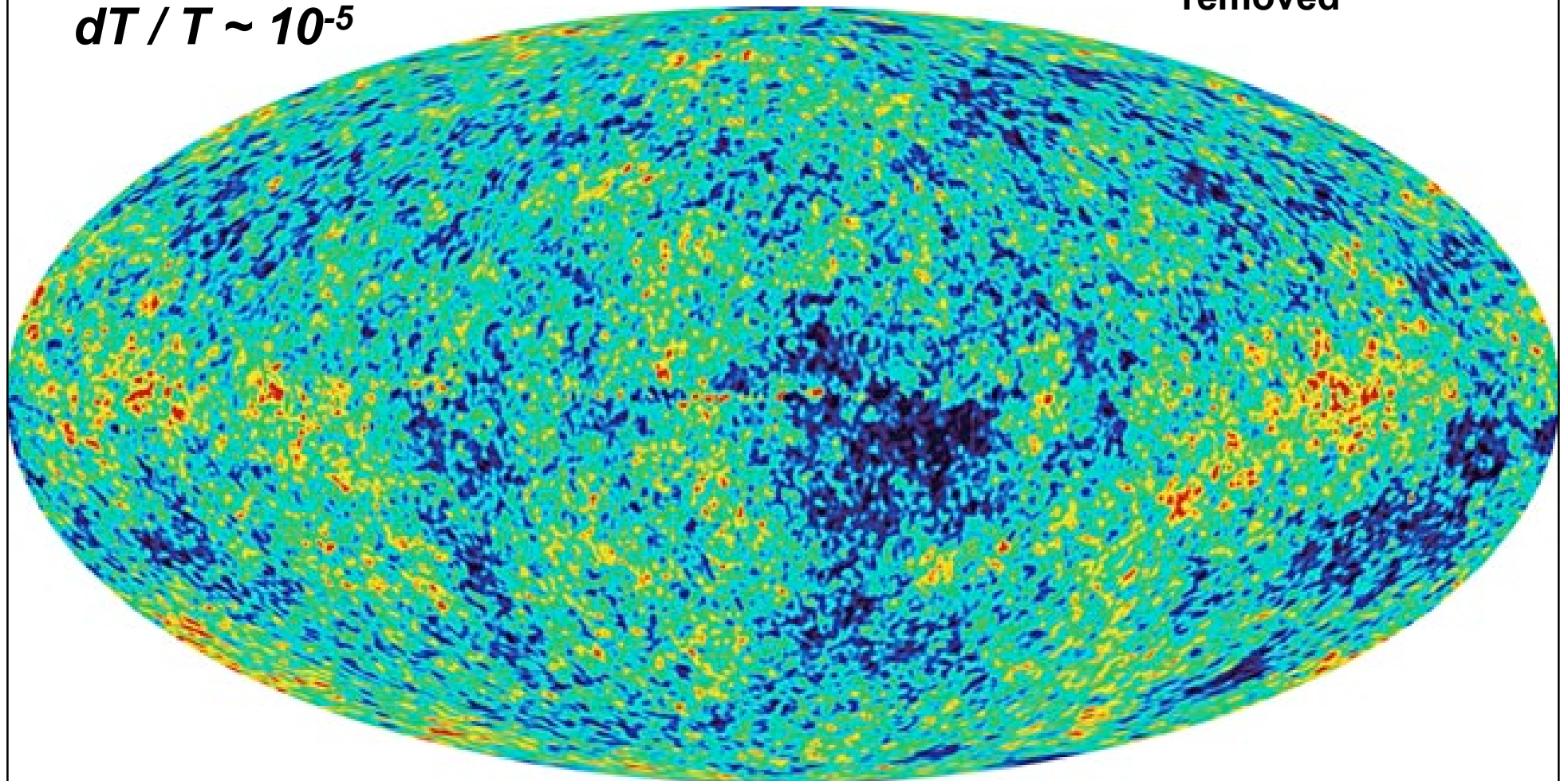


$T = 2.73 \text{ K}$

2003 WMAP all-sky

Dipole and
foreground galaxy
removed

$dT / T \sim 10^{-5}$

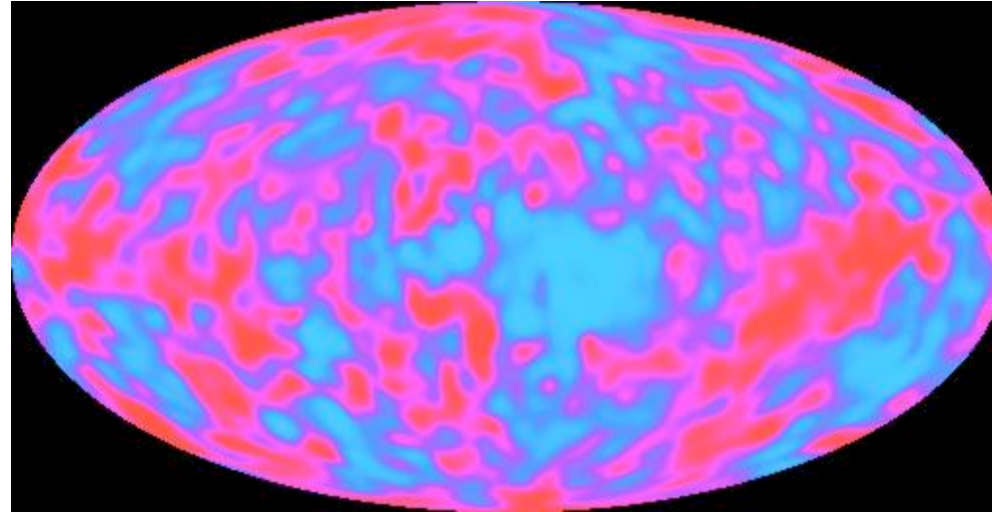


Snapshot at $z=1100$ of quantum fluctuations stretched by inflation.

Dark matter potential wells that seed later galaxy formation.

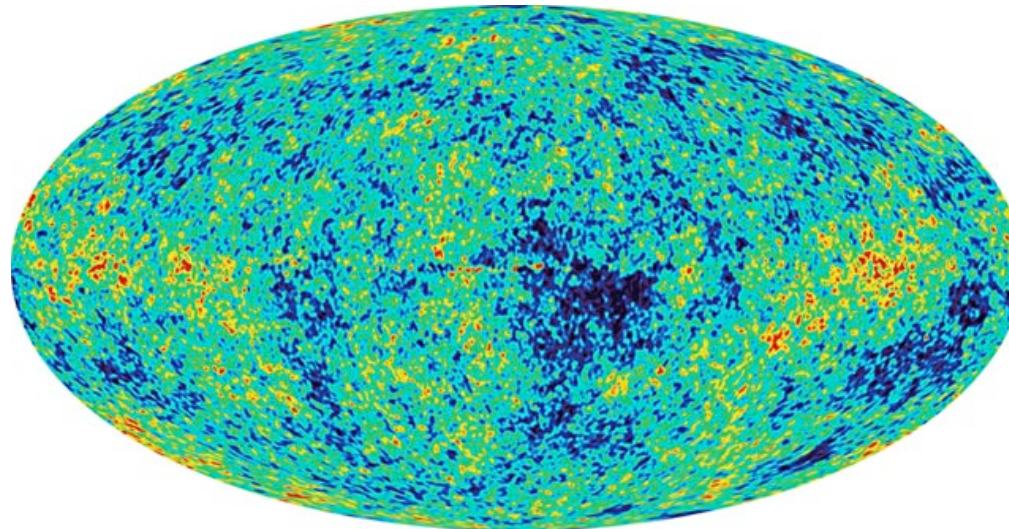
CMB Anisotropies

COBE
1994



$$\frac{\Delta T}{T} \sim 10^{-5}$$

WMAP
2004



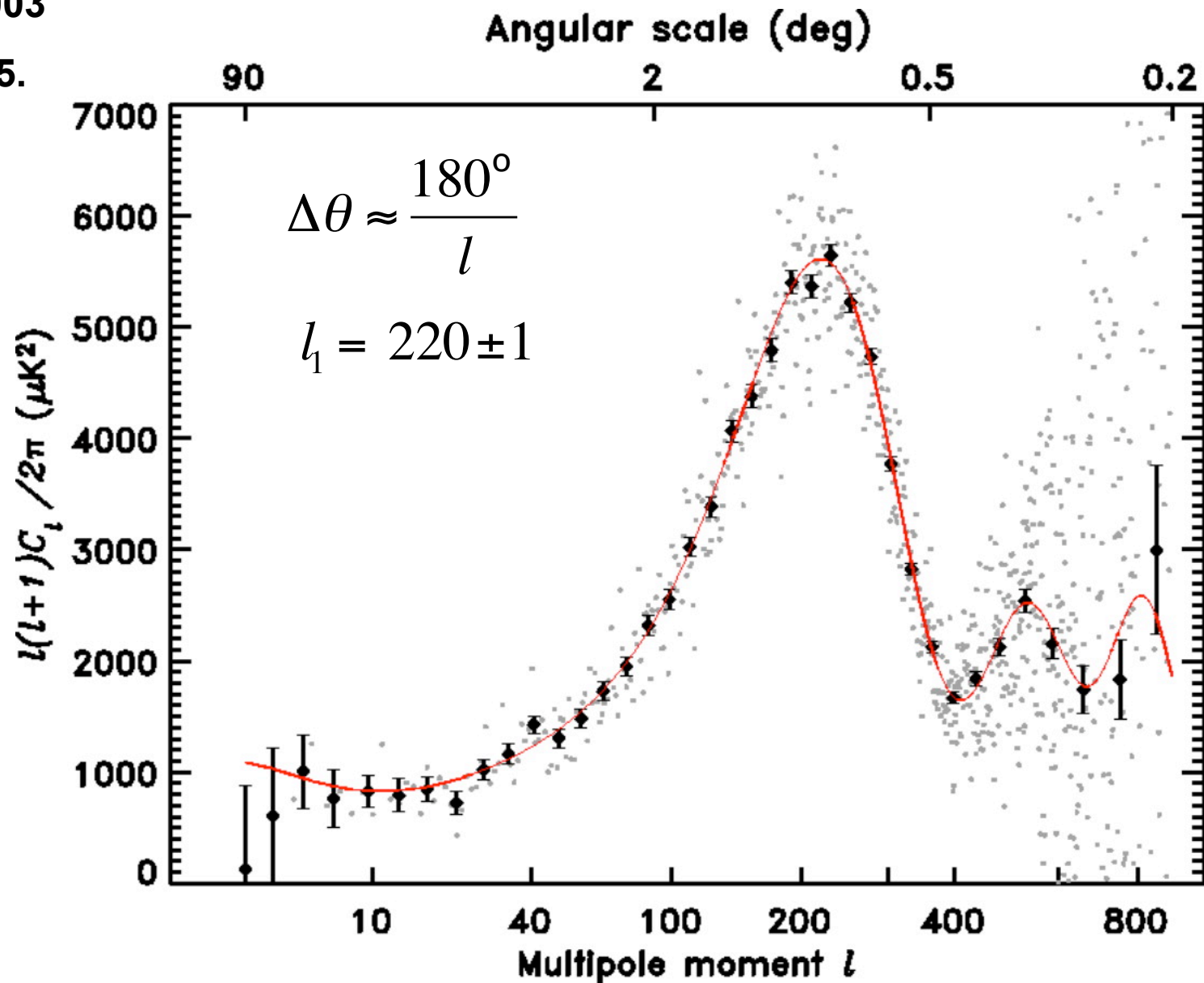
$$\Delta\theta \sim 1^\circ$$

Snapshot of Universe at $z = 1100$
Seeds that later form galaxies.

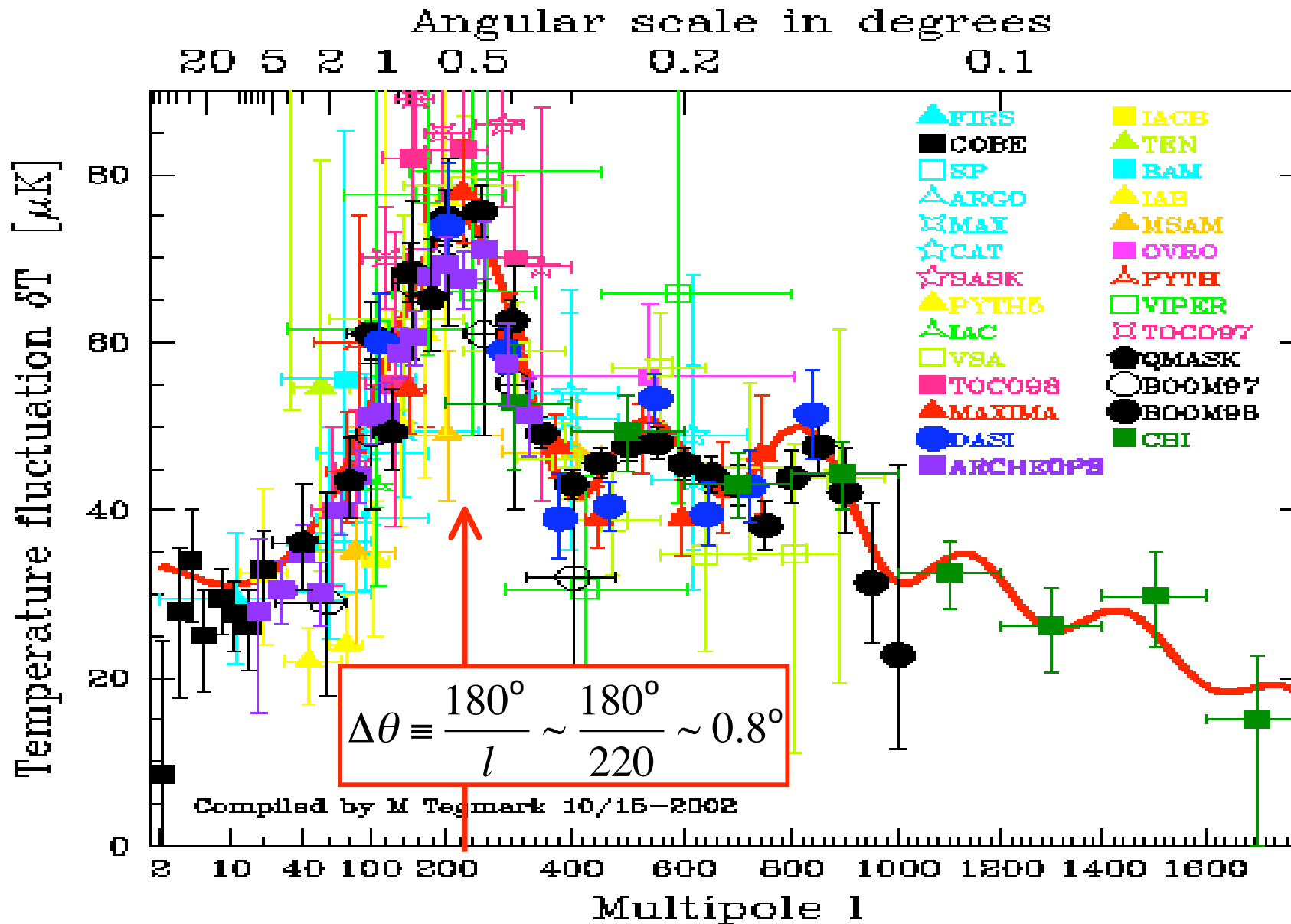
2003 -- WMAP Power Spectrum

Spergel et al. 2003

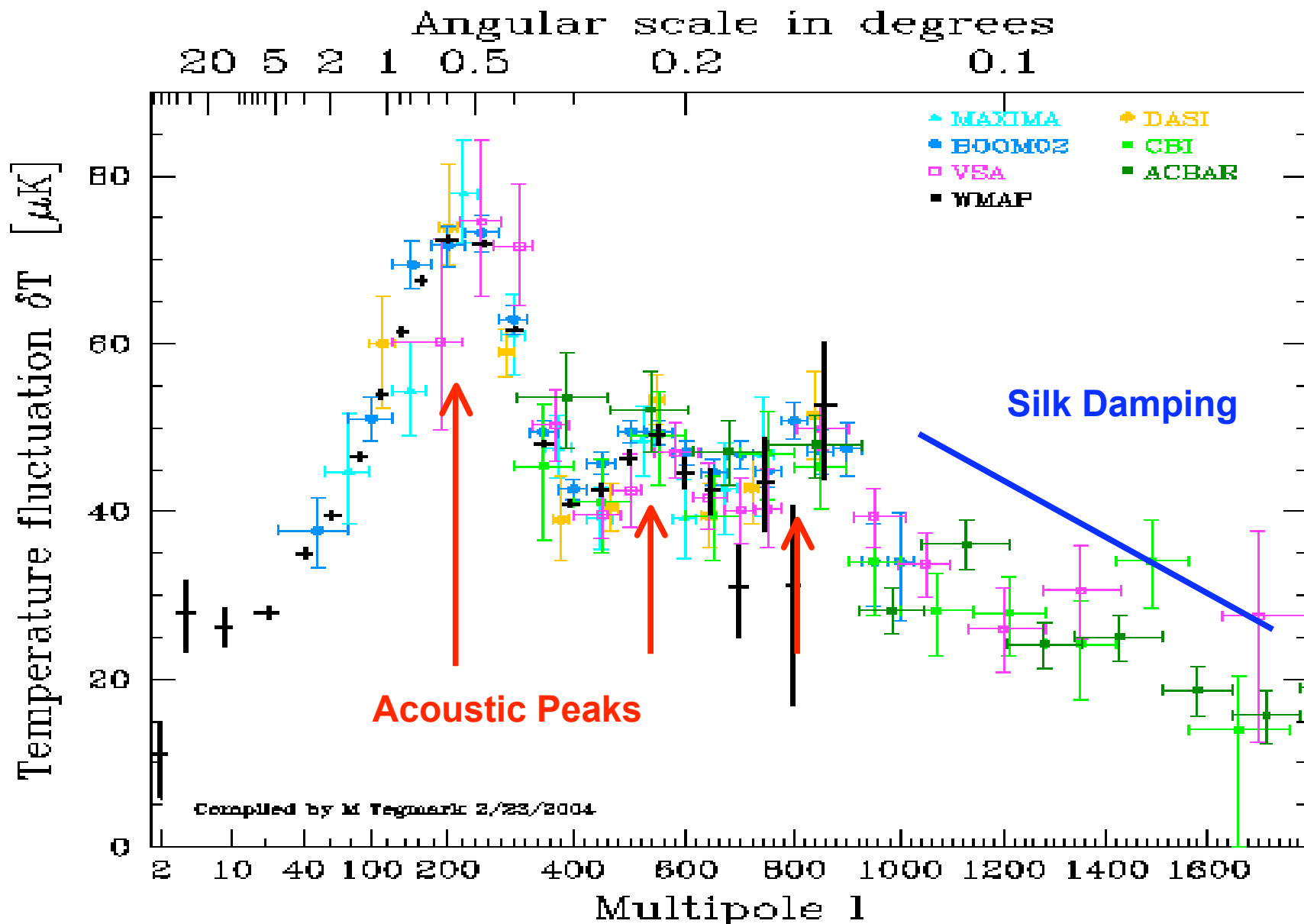
ApJSup 148,175.



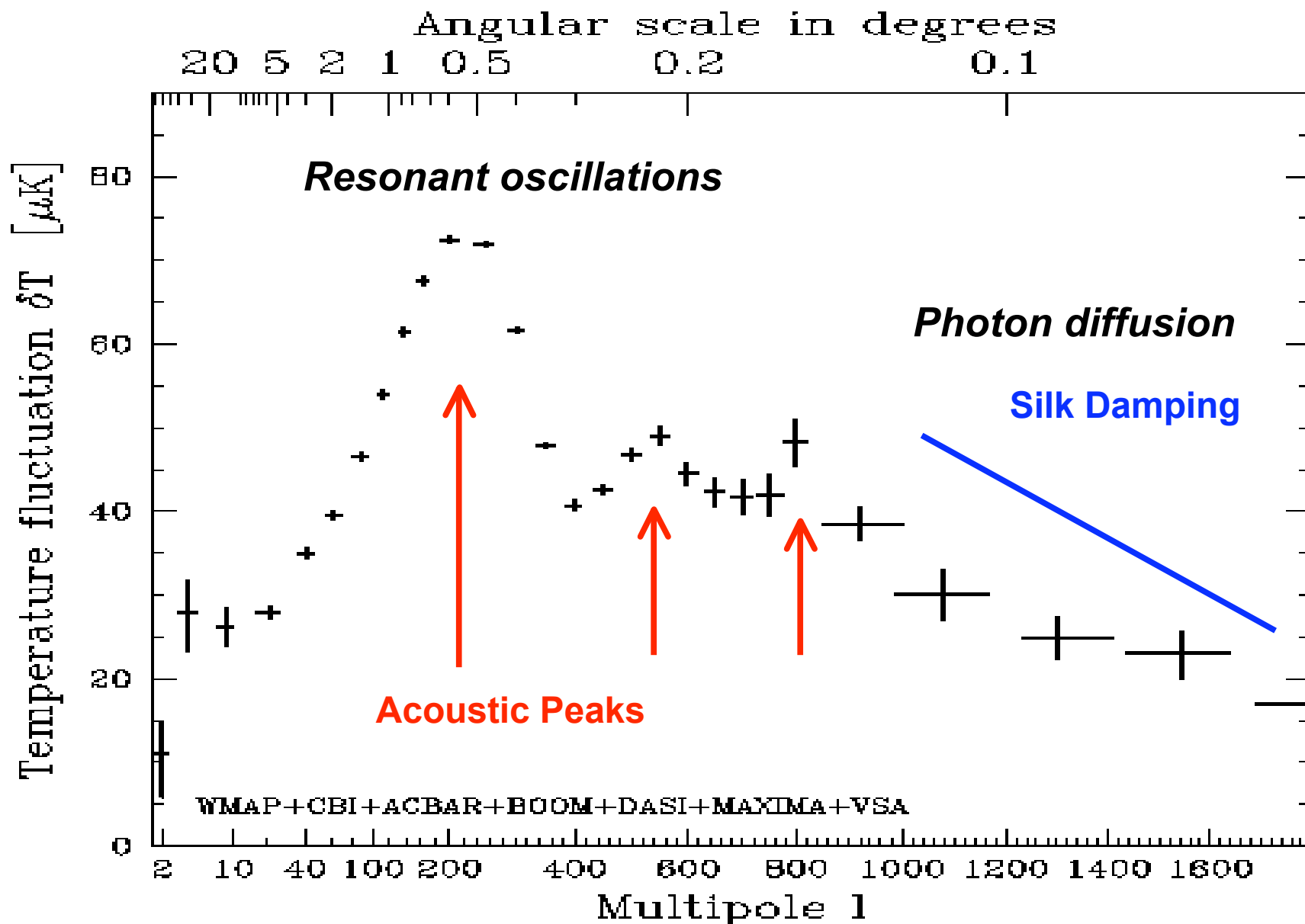
2002 - CMB Power Spectrum



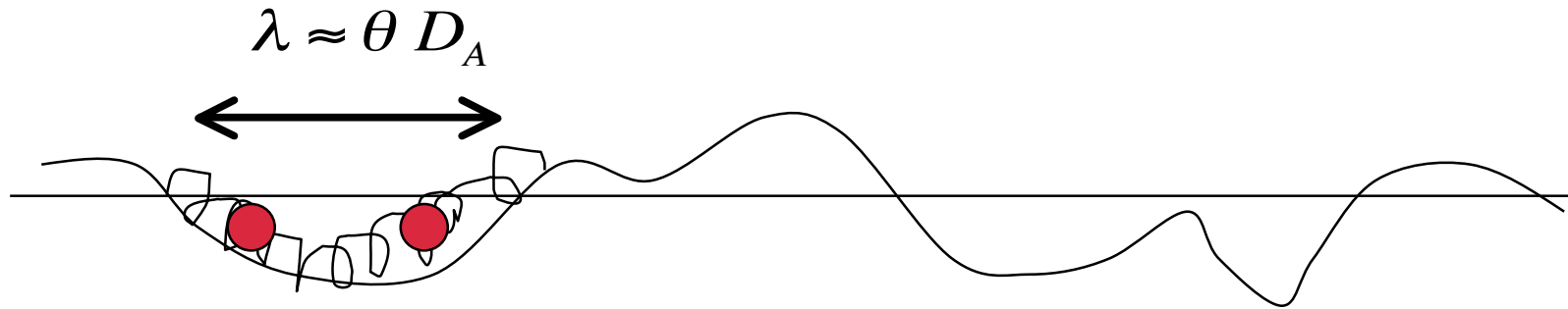
2004 - CMB Power Spectrum



2004 - CMB Binned



Acoustic Oscillations



Dark Matter potential wells - many sizes.

photon-electron-baryon fluid

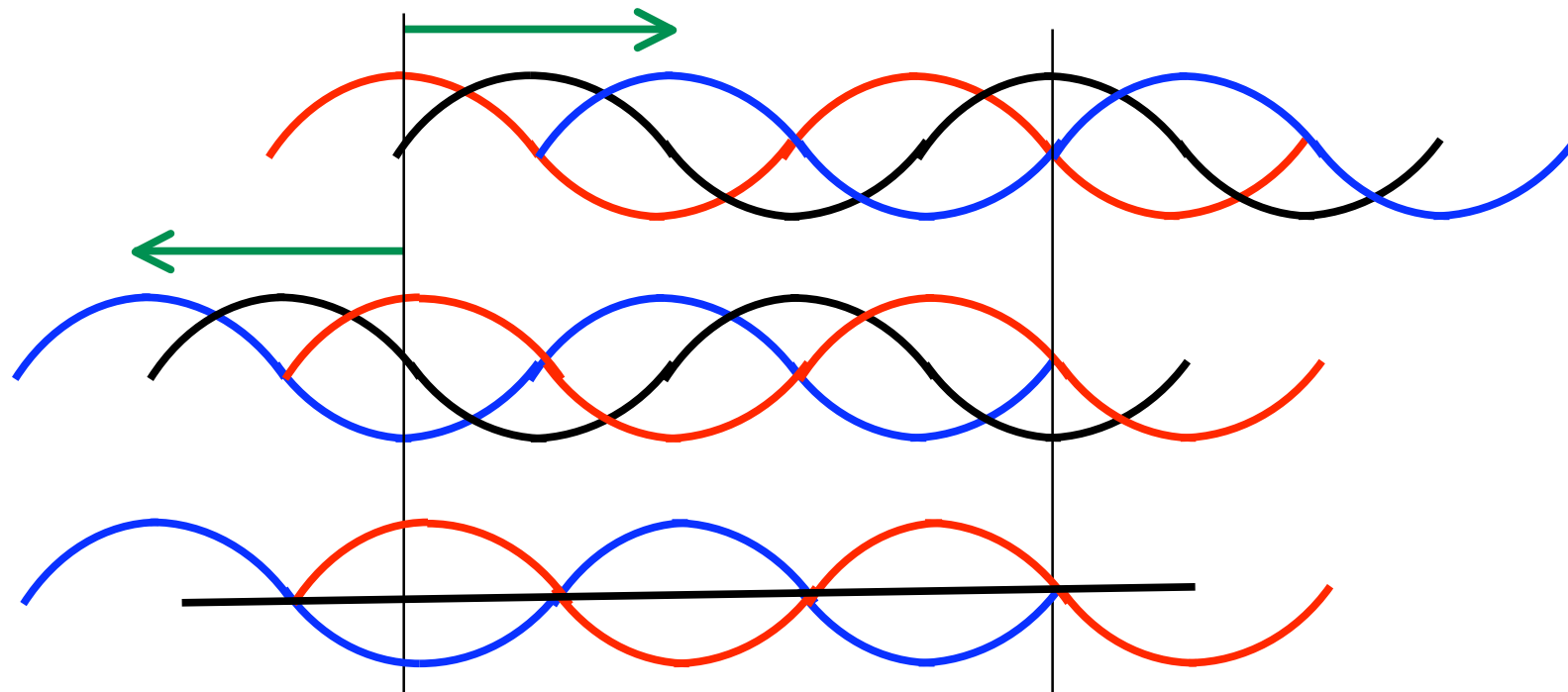
fluid falls into DM wells

photon pressure pushes it out again

oscillations starting at $t = 0$ (post-inflation)

stopping at $z = 1100$ (recombination)

Standing Sound Waves



temperature oscillations:

$$c_s = \frac{\lambda}{P} = \frac{\omega}{k} = \frac{c}{\sqrt{3}}$$

$$\Delta T(x, t) = a \cos(\omega t) \cos(k x) \quad \omega = \frac{2\pi}{P} \quad k = |\mathbf{k}| = \frac{2\pi}{\lambda}$$

$$\Delta T(\mathbf{x}, t) = \sum_{\mathbf{k}} a(\mathbf{k}) \cos(\omega t) \cos(\mathbf{k} \cdot \mathbf{x}) \quad \langle a(k) \rangle = a_0 \left(\frac{k}{k_0} \right)^{n_s}$$

Resonant Oscillations

size of potential well λ

oscillation period $P \approx \frac{\lambda}{c_s}$

sound speed $c_s = \frac{c}{\sqrt{3}}$

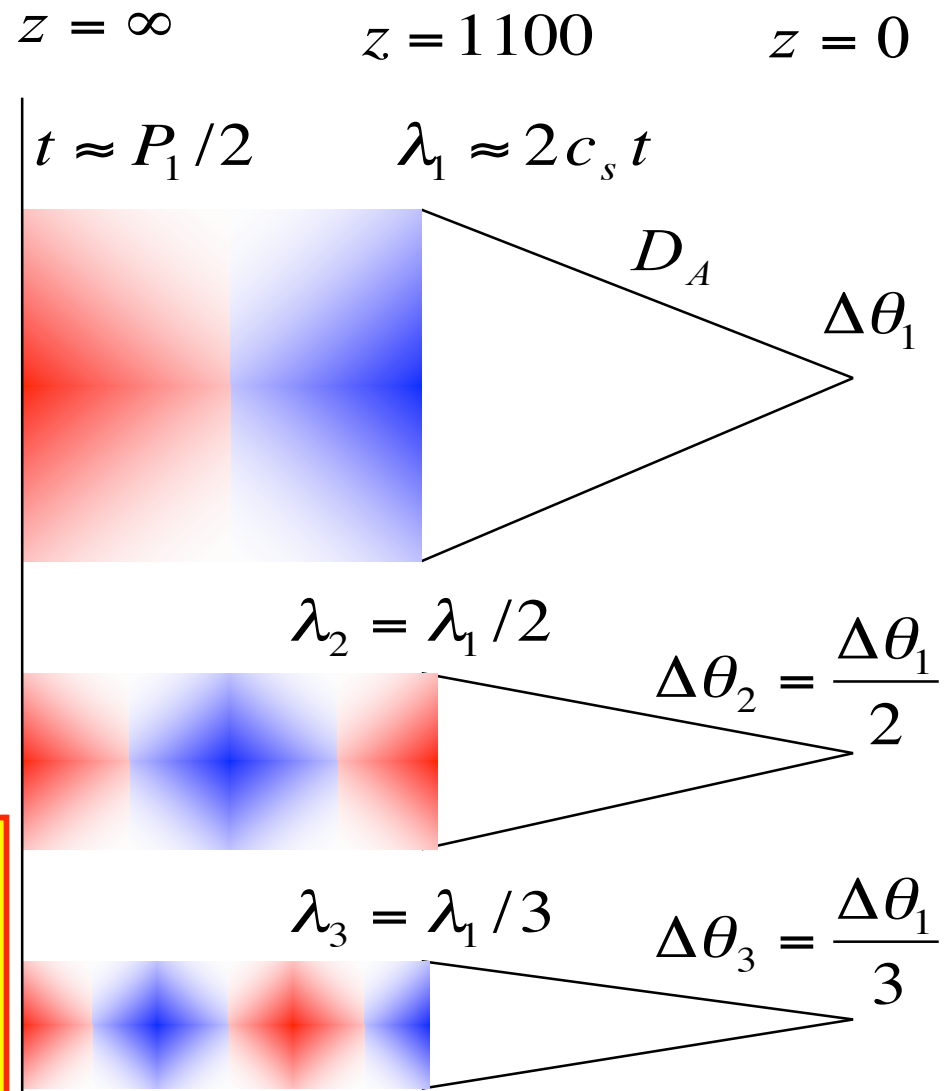
temperature oscillations

$$\Delta T(t) = \Delta T(0) \cos(2\pi t/P)$$

$\max|\Delta T|$ at $t = \frac{nP}{2} \sim \frac{n\lambda}{2c_s}$

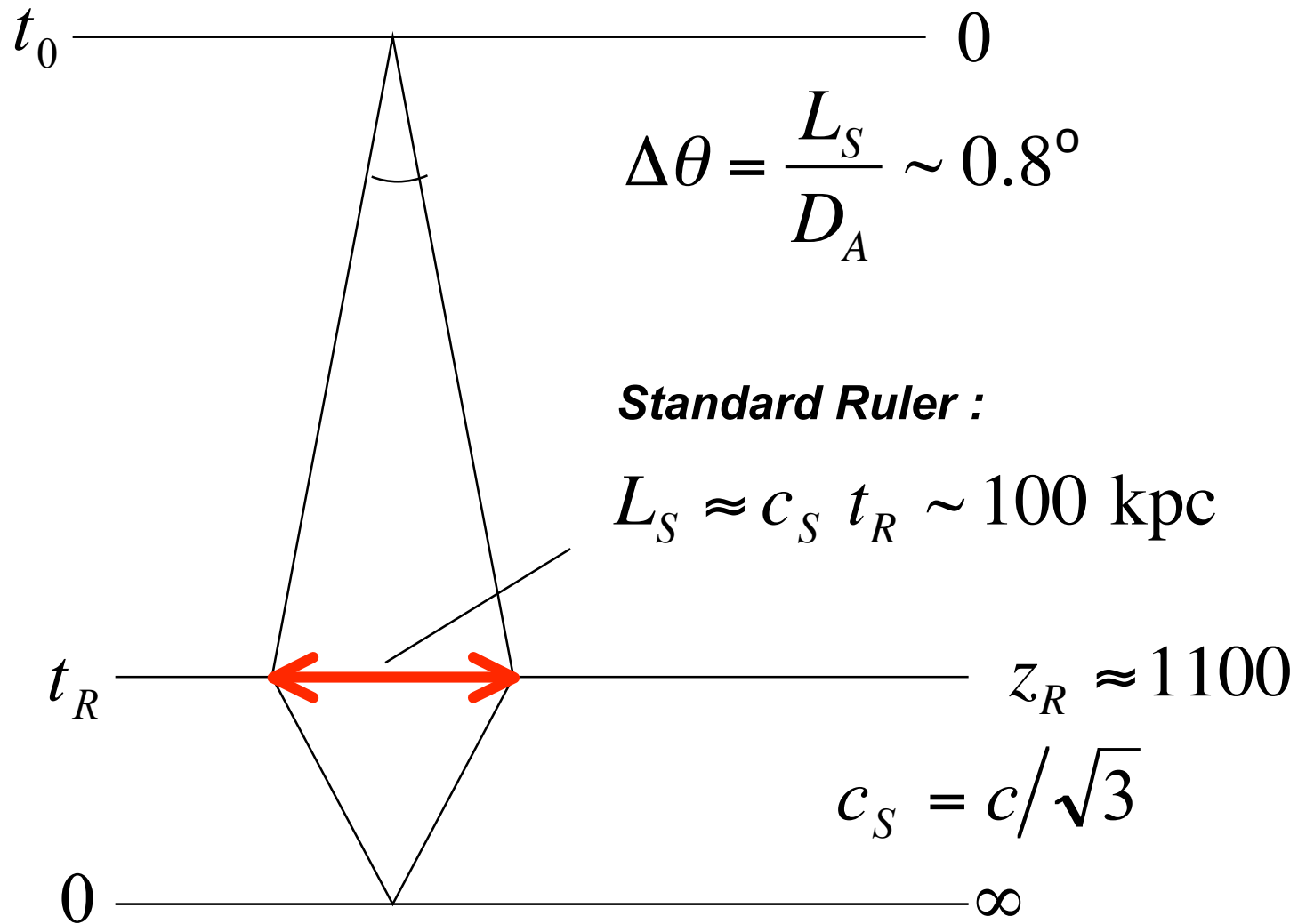
angular size

$$\Delta\theta_n = \frac{\lambda_n}{D_A} = \frac{\Delta\theta_1}{n} \quad \Delta\theta_1 \approx \frac{2c_s t}{D_A} \sim 0.8^\circ$$

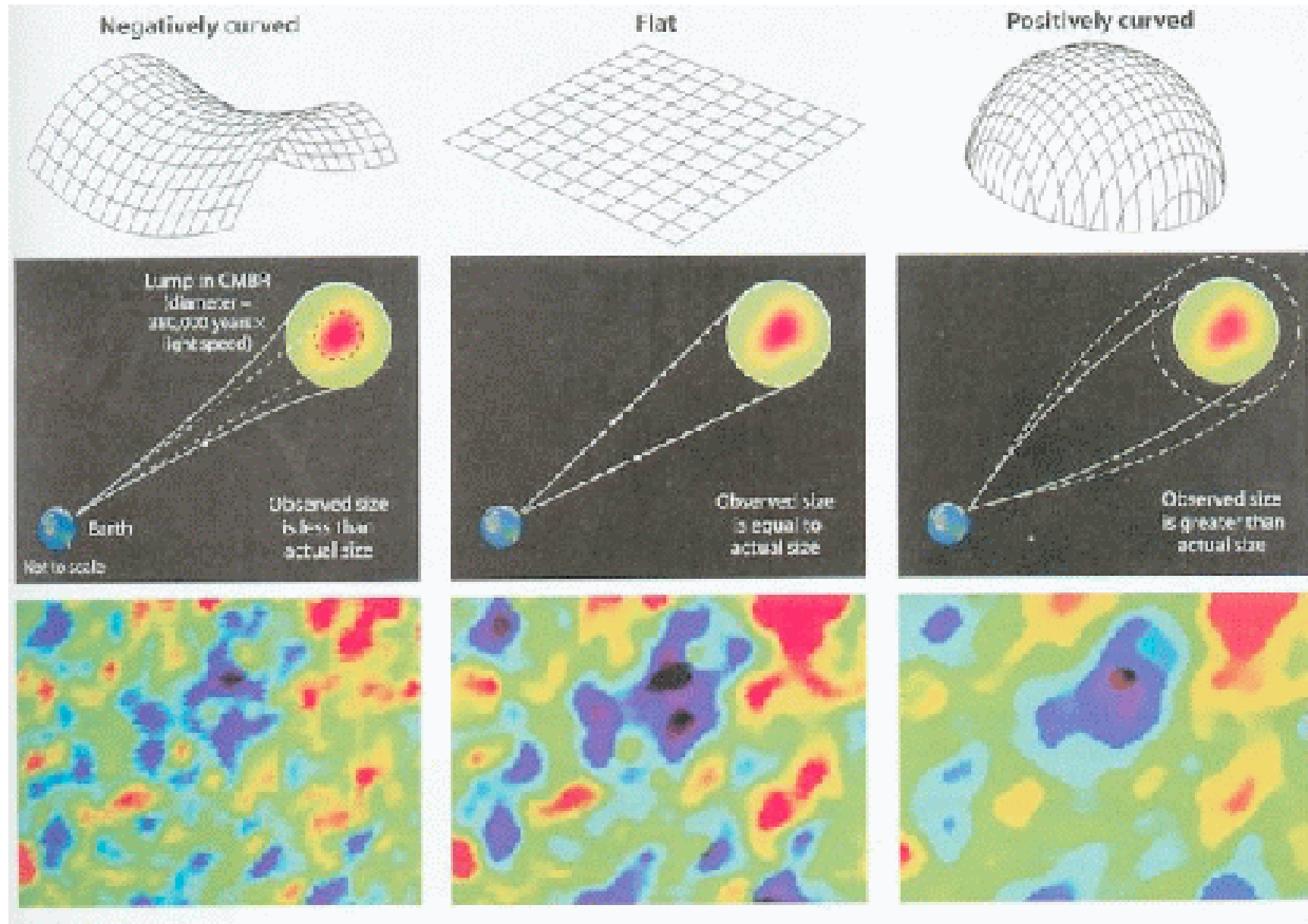


Smaller wells oscillate faster.

Sound Horizon at $z = 1100$



Angular scale --> Geometry



Sound Horizon at $z = 1100$

distance travelled by a sound wave

$$c_S dt$$

recombination
at $z = 1100$

$$x \equiv 1 + z = \frac{R_0}{R(t)}$$

expand each step by factor $R(t_R)/R(t)$:

$$dt = \frac{-dx}{x H(x)}$$

$$L_S(t_R) = R(t_R) \int_0^{t_R} \frac{c_S dt}{R(t)}$$

sound speed

$$= \frac{R_0}{1+z} \int_{1+z}^{\infty} \frac{x}{R_0} \frac{c_S dx}{x H(x)}$$

$$dt = -dx / x H(x)$$

$$R(t) = R_0 / x$$

$$c_S \approx \frac{c}{\sqrt{3}}$$

$$= \frac{c_S}{(1+z)} \int_{1+z}^{\infty} \frac{dx}{H(x)}$$

$H(x)$ from Friedmann Eqn.

$$= \frac{c_S}{(1+z) H_0} \int_{1+z}^{\infty} \frac{dx}{\sqrt{x^4 \Omega_R + x^3 \Omega_M + \Omega_\Lambda + (1 - \Omega_0) x^2}}$$

$$\approx \frac{c_S}{(1+z) H_0} \int_{1+z}^{\infty} \frac{dx}{\sqrt{x^4 \Omega_R + x^3 \Omega_M}}$$

keep 2 largest terms.

Sound Horizon at $z = 1100$

$$\begin{aligned}
 L_S(t_R) &= \frac{c_s}{(1+z)} \int_{1+z}^{\infty} \frac{dx}{H(x)} \approx \frac{c_s}{(1+z) H_0} \int_{1+z}^{\infty} \frac{dx}{\sqrt{x^4 \Omega_R + x^3 \Omega_M}} \\
 &= \frac{c_s}{(1+z) H_0 \sqrt{\Omega_R}} \int_{1+z}^{\infty} \frac{dx}{\sqrt{x^3(x+x_0)}} \quad x_0 \equiv \frac{\Omega_M}{\Omega_R} \approx 3500 \left(\frac{\Omega_M}{0.3} \right) \\
 &= \frac{c_s}{(1+z) H_0 \sqrt{\Omega_R}} \left(-\frac{2}{x_0} \sqrt{1 + \frac{x_0}{x}} \right)_{1+z}^{\infty} \\
 &= \frac{2c_s}{(1+z) H_0 \sqrt{\Omega_M x_0}} \left(\sqrt{1 + \frac{x_0}{1+z}} - 1 \right) \quad c_s = \frac{c}{\sqrt{3}} \\
 &= \frac{c}{H_0} \frac{2(\sqrt{4.6} - 1)}{1100 \sqrt{3 \times 0.3 \times 3500}} \\
 &= 3.4 \times 10^{-5} \frac{c}{H_0} \approx 110 \left(\frac{0.7}{h} \right) \left(\frac{0.3}{\Omega_M} \right)^{1/2} \text{ kpc}
 \end{aligned}$$

**Expands by factor
 $1 + z = 1100$
to ~ 120 Mpc today.**

Angular Scale measures Ω_0

sound horizon :

$$L_S(z) = \frac{1}{1+z} \int_{1+z}^{\infty} \frac{c_S dx}{H(x)}$$

angular diameter distance :

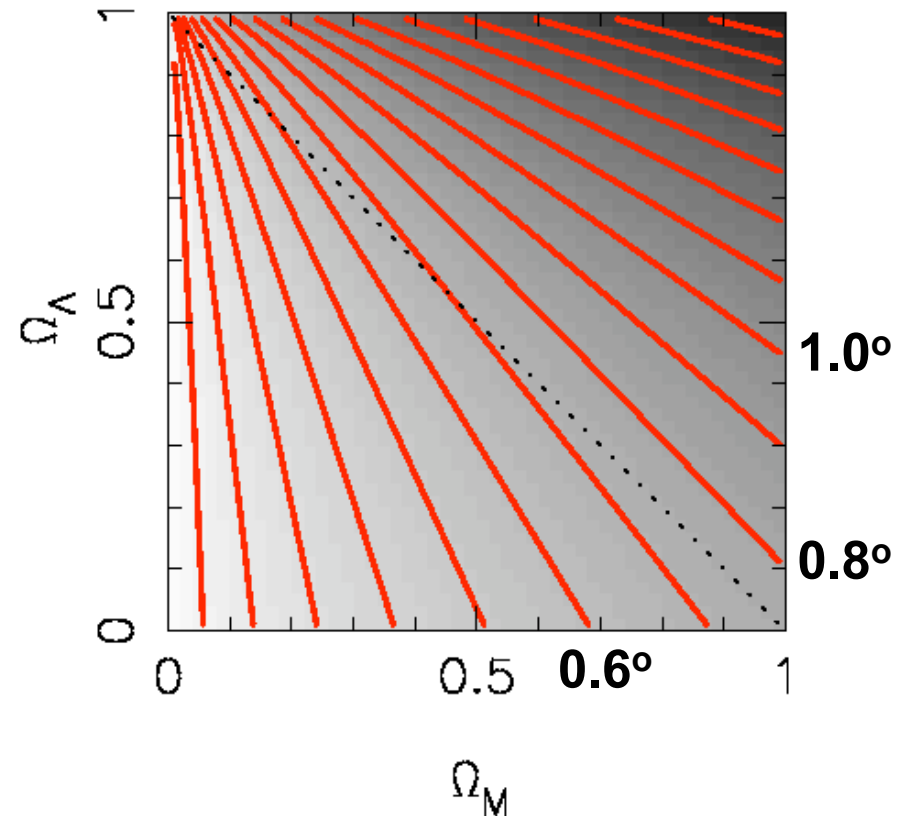
$$D_A(z) = \frac{R_0 S_K(\chi)}{1+z}$$

$$\chi = \int_t^{t_0} \frac{c dt}{R(t)} = \frac{c}{R_0} \int_1^{1+z} \frac{dx}{H(x)}$$

angular scale

$$\theta = \frac{L_S(z)}{D_A(z)} = \frac{\int_{1+z}^{\infty} \frac{c_S dx}{H(x)}}{R_0 S_k \left(\frac{c}{R_0} \int_1^{1+z} \frac{dx}{H(x)} \right)}$$

$$\Omega_R = 0.000086$$



Angular scale depends mainly on the curvature.

Gives $\theta \sim 0.8^\circ$ for flat geometry,

$$\Omega_0 = \Omega_M + \Omega_\Lambda = 1$$

Finer Details: measure Ω_b and Ω_M

Sound speed not constant :

$$c_s(z) = \frac{c}{\sqrt{3(1+R(z))}}$$

$$R(z) \equiv \frac{3\rho_b(z)}{4\rho_R(z)} = \frac{3\Omega_b(1+z)}{4\Omega_R}$$

Acoustic peaks not quite equally spaced:

$$l_n = l_A (n + \delta_n)$$

phase shifts

$$\delta_n \approx a_n \left(\frac{r}{0.3} \right)^{0.1} \quad a_{1,2,3} \approx 0.267, 0.24, 0.35$$

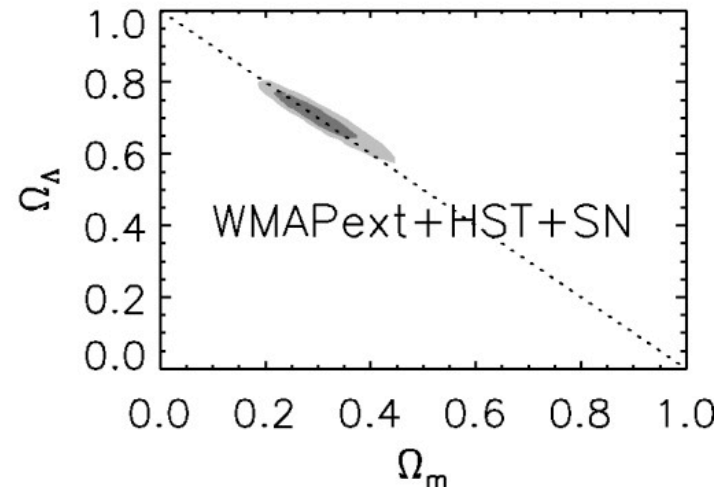
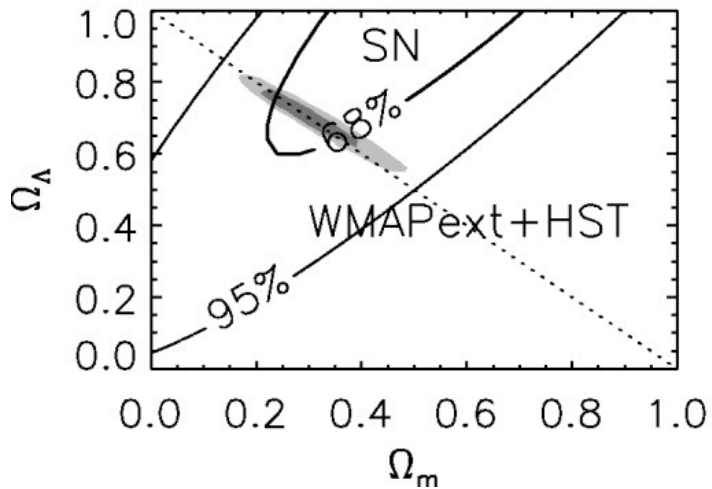
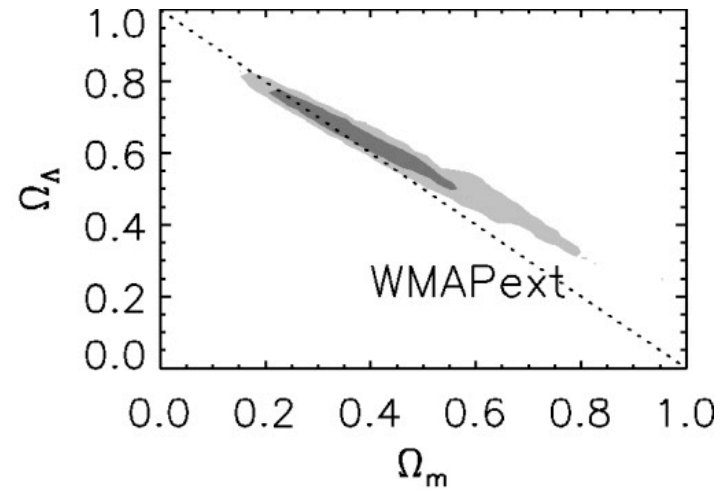
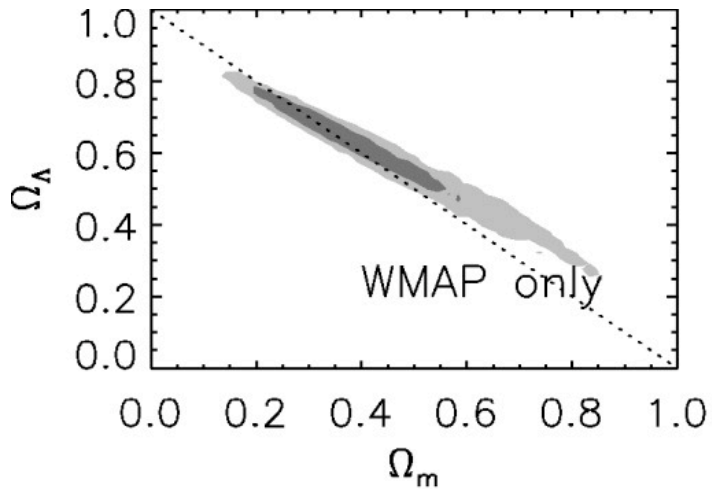
$$r \equiv \frac{\rho_M(z)}{\rho_R(z)} = \frac{\Omega_M(1+z)}{\Omega_R}$$

Max Tegmark's CMB Movies

Shows how the **CMB power spectrum**
(and the **baryon power spectrum**)
depend on the cosmological parameters.

A link to these is available on the course web page.

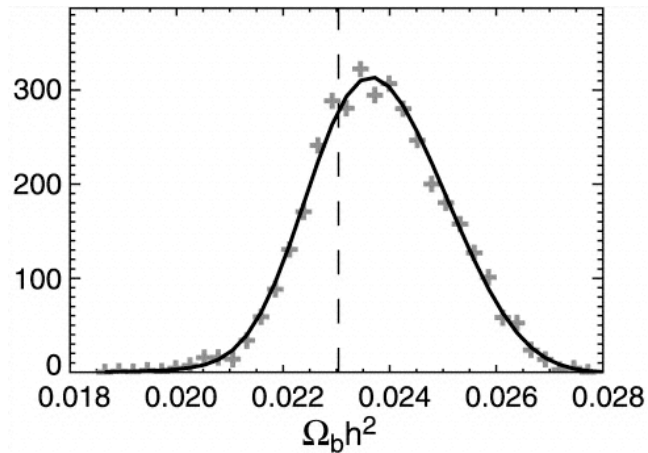
WMAP parameter constraints



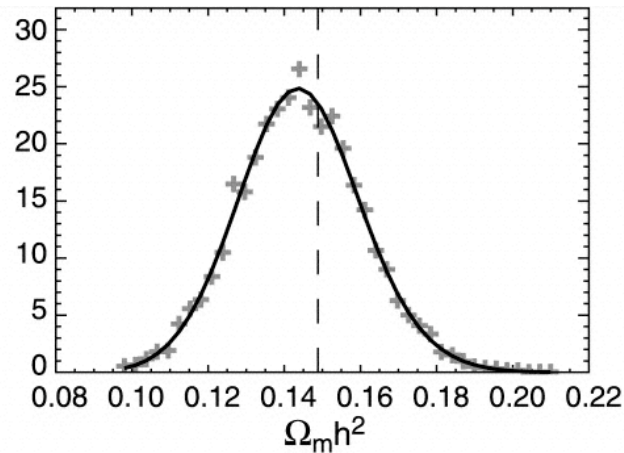
WMAP parameter constraints

Spergel et al. 2003 ApJSup 148,175.

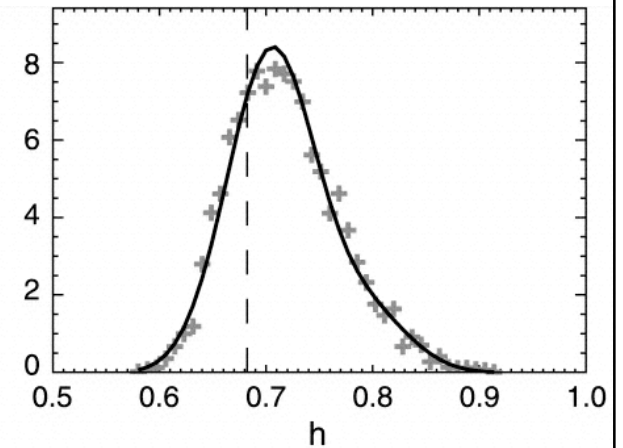
baryons



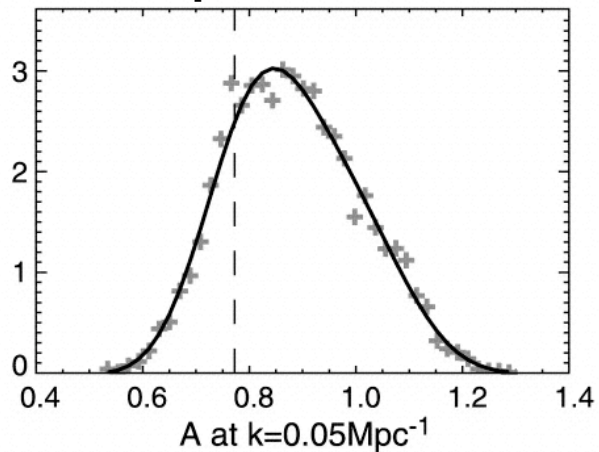
Dark Matter



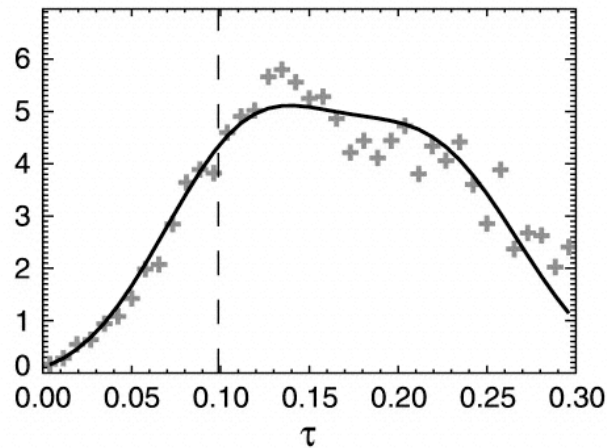
expansion



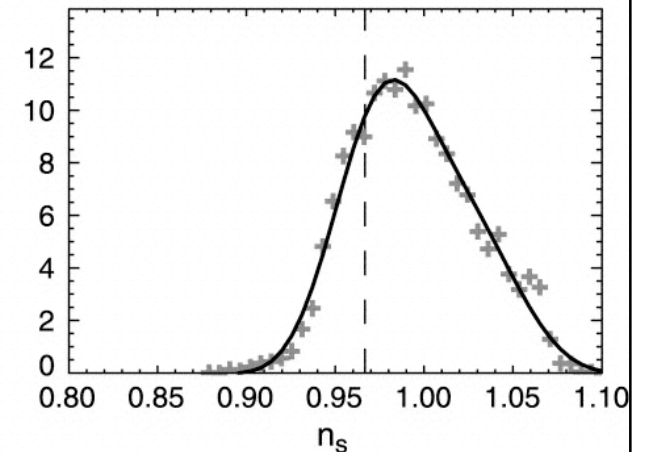
amplitude



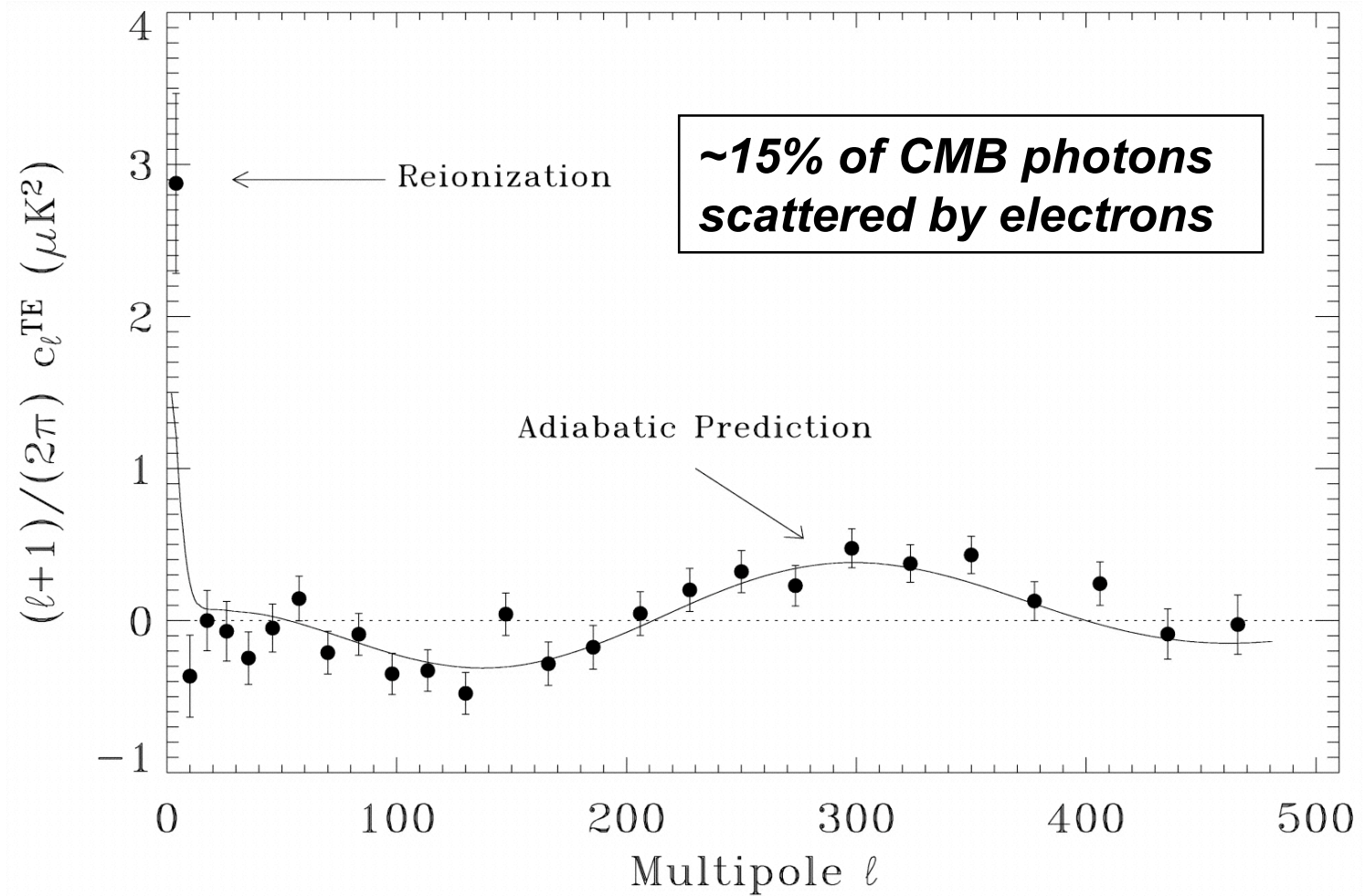
reionisation



tilt



WMAP Polarisation Power Spectrum



Epoch of Re-Ionisation

UV from first stars re-ionises gas.

Scatters ~15% of CMB photons yielding ~15% polarisation.

WMAP measured this !

Electron scattering optical depth:

$$d\tau = n \sigma_T dr$$

$$= n_0 (1+z)^3 \sigma_T c dt$$

$$\tau = n_0 \sigma_T c \int_1^{1+z} \frac{x^3 dx}{x H(x)}$$

$$= \frac{n_0 \sigma_T c}{H_0} \int_1^{1+z} \frac{x^2 dx}{\sqrt{\Omega_M x^3 + \Omega_\Lambda + (1 - \Omega_0) x^2}}$$

Gives ~15% optical depth at $z \sim 20$

$$dt = \frac{-dx}{x H(x)}$$

$$x \equiv 1 + z$$

Thompson cross - section

σ_T

electron density today

$$n_0 = \frac{\Omega_b}{m_H} \frac{3 H_0^2}{8 \pi G} \left(X + \frac{Y}{2} \right)$$

Precision Cosmology

$h = 71 \pm 3$ expanding

$\Omega = 1.02 \pm 0.02$ flat

$\Omega_b = 0.044 \pm 0.004$ baryons

$\Omega_M = 0.27 \pm 0.04$ Dark Matter

$\Omega_\Lambda = 0.73 \pm 0.04$ Dark Energy

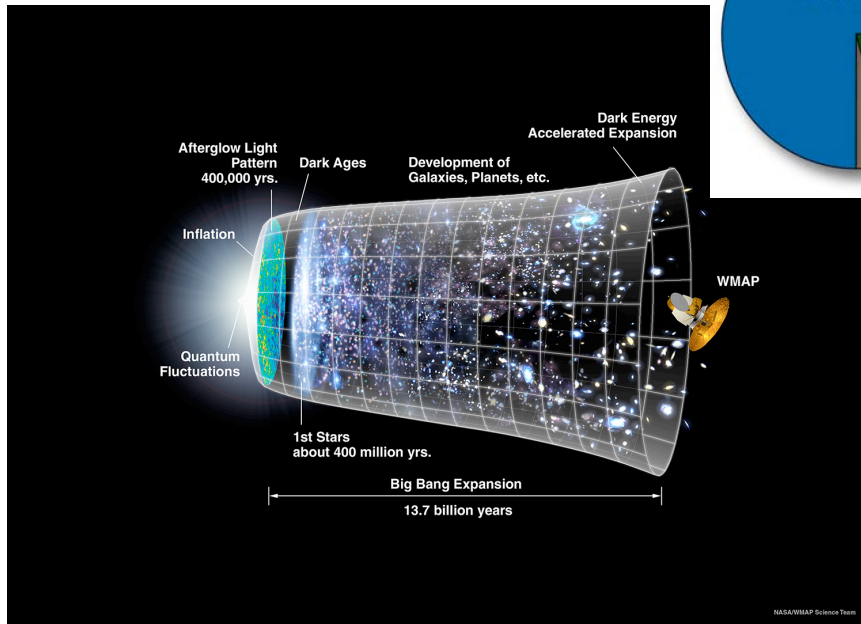
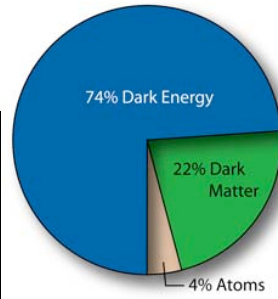
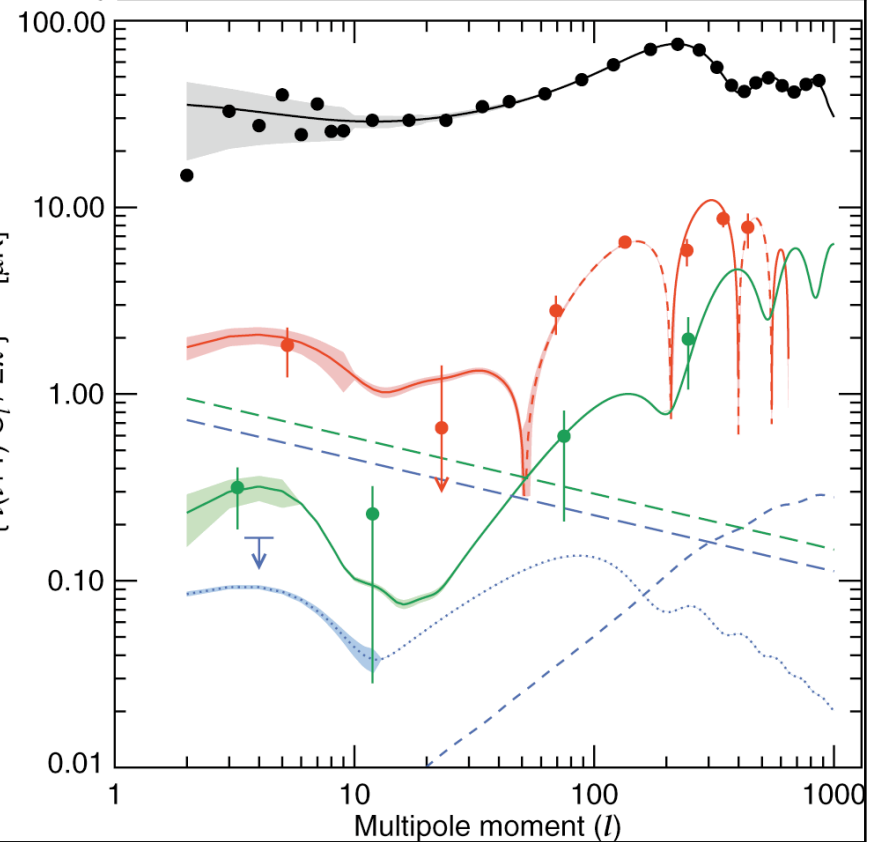
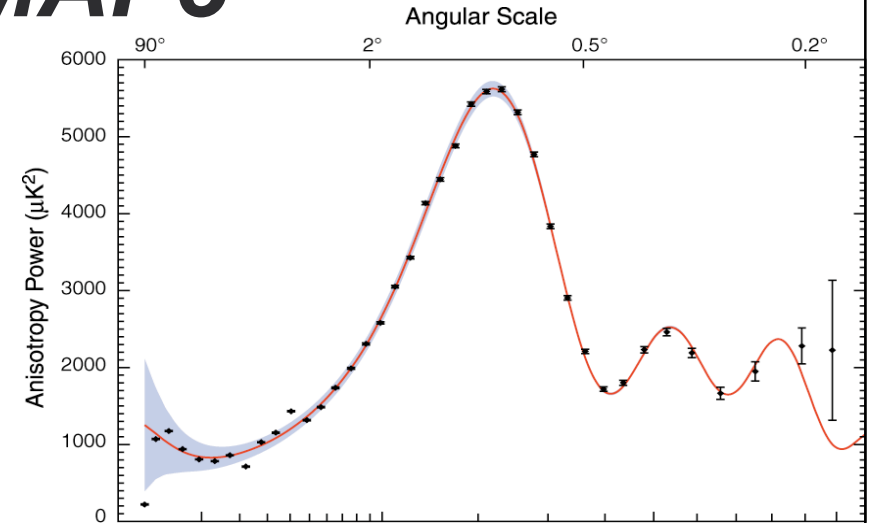
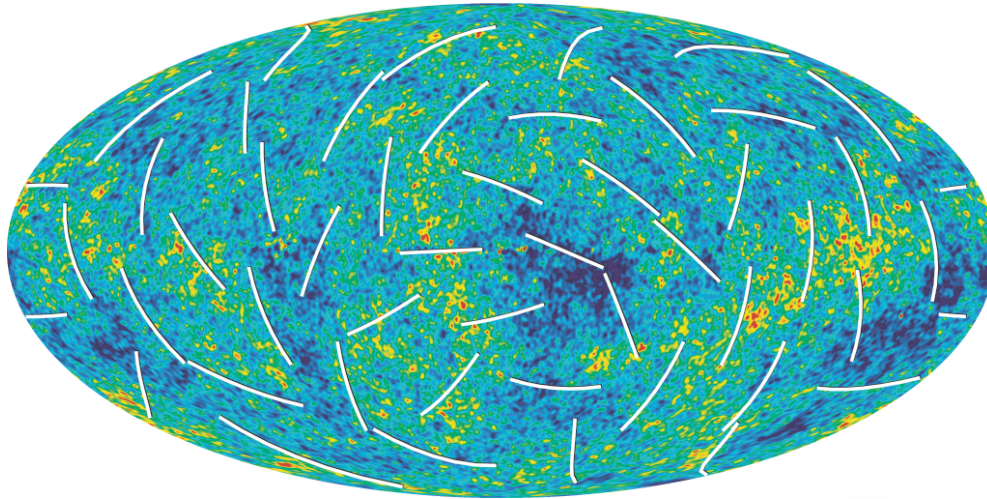
$t_0 = 13.7 \pm 0.2 \times 10^9$ yr now

$t_* = 180^{+220}_{-80} \times 10^6$ yr $z_* = 20^{+10}_{-5}$ reionisation

$t_R = 379 \pm 1 \times 10^3$ yr $z_R = 1090 \pm 1$ recombination

(From the WMAP 1-year data analysis)

2006 - WMAP3



Planck -- 2009 launch

