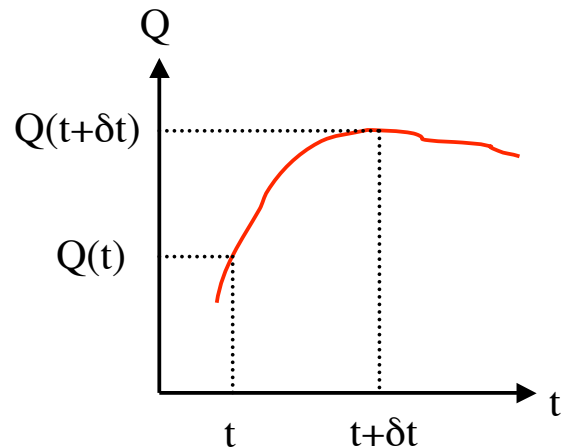


## 2.3 Relating Eulerian / Lagrangian descriptions

- Reminder: consider a function  $Q(t)$

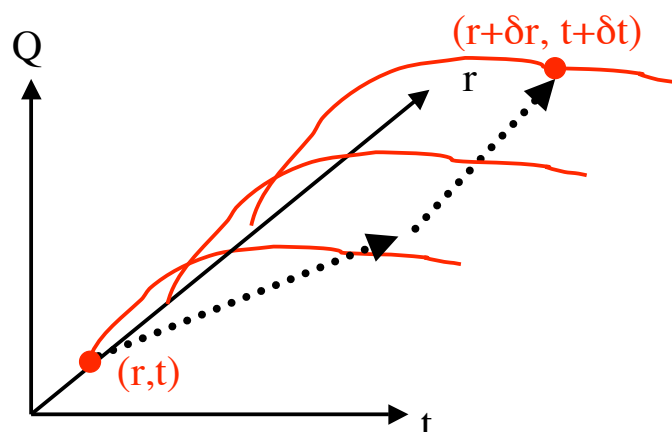
$$\frac{dQ}{dt} = \frac{Q(t + \delta t) - Q(t)}{\delta t}$$



- Now make  $Q$  a function of two variables:

$$\frac{dQ}{dt} = \frac{Q(\underline{r} + \delta \underline{r}, t + \delta t) - Q(\underline{r}, t)}{\delta t}$$

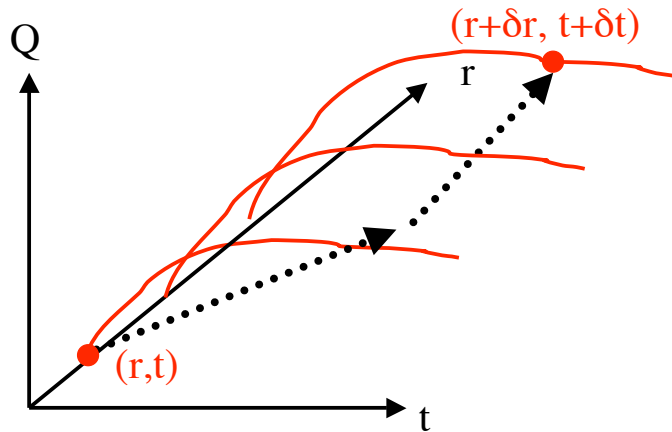
- $\underline{r}$  = position of fluid element at time  $t$
- $\underline{r} + \delta \underline{r}$  = position of same fluid element at time  $t + \delta t$



- First of all, separate out the variation in  $\underline{r}$  and  $t$

$$\begin{aligned}
 & Q(\underline{r} + \delta \underline{r}, t + \delta t) - Q(\underline{r}, t) \\
 &= \underbrace{Q(\underline{r}, t + \delta t) - Q(\underline{r}, t)}_{\text{Variation in time } t \text{ at fixed } \underline{r}} + \underbrace{Q(\underline{r} + \delta \underline{r}, t + \delta t) - Q(\underline{r}, t + \delta t)}_{\text{Variation in space } \underline{r} \text{ at fixed } t + \delta t}
 \end{aligned}$$

Variation in time  $t$  at fixed  $\underline{r}$       Variation in space  $\underline{r}$  at fixed  $t + \delta t$



- Then, write the numerator as an expansion in  $\delta \underline{r}$  and  $\delta t$

Remember that

$$\begin{cases}
 Q(t + \delta t) = Q(t) + \delta t \frac{\partial Q}{\partial t} + \dots \\
 Q(\underline{r} + \delta \underline{r}) = Q(\underline{r}) + \delta \underline{r} \cdot \underline{\nabla} Q + \dots
 \end{cases}$$

So

$$\begin{aligned}
 & Q(\underline{r} + \delta \underline{r}, t + \delta t) - Q(\underline{r}, t) \\
 &= \delta t \frac{\partial Q}{\partial t} + \dots + \delta \underline{r} \cdot \underline{\nabla} Q + \dots
 \end{aligned}$$

- So we have:

$$Q(\underline{r} + \delta \underline{r}, t + \delta t) - Q(\underline{r}, t)$$

$$\approx \frac{\partial Q}{\partial t} \delta t + \delta \underline{r} \cdot \underline{\nabla} Q$$

← Evaluated at  $t + \delta t$

$$\approx \frac{\partial Q}{\partial t} \delta t + \delta \underline{r} \cdot \left[ \underline{\nabla} Q + \delta t \frac{\partial}{\partial t} \underline{\nabla} Q \right]$$

← Evaluated at  $t$

- finally:

$$\frac{dQ}{dt} = \frac{Q(\underline{r} + \delta \underline{r}, t + \delta t) - Q(\underline{r}, t)}{\delta t}$$

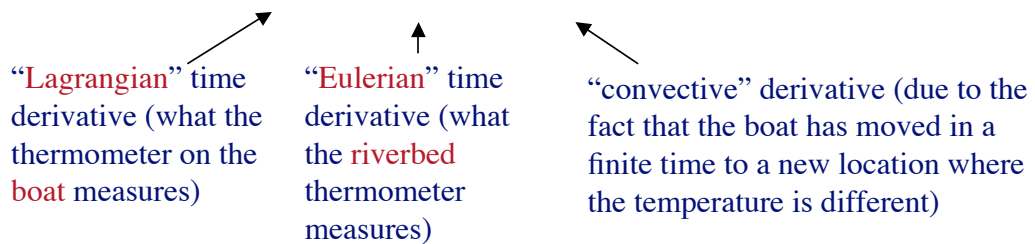
$$\approx \frac{\frac{\partial Q}{\partial t} \delta t + \delta \underline{r} \cdot \left[ \underline{\nabla} Q + \frac{\partial}{\partial t} \underline{\nabla} Q \delta t \right]}{\delta t}$$

$$\approx \frac{\partial Q}{\partial t} + \frac{\delta \underline{r}}{\delta t} \cdot \underline{\nabla} Q$$

← But this is just the velocity!

- Hence if the flow velocity is  $\underline{u}$

$$\frac{dQ}{dt} = \frac{\partial Q}{\partial t} + \underline{u} \cdot \nabla Q \quad (2.1)$$



**NB:** In a *steady state* the Eulerian time derivative would be zero, but the Lagrangian time derivative would only be zero if  $Q$  was also *uniform*.

## 2.4 Streamlines and streamfunctions

- We can write a 2D flow  $\underline{u}(x,y)$  in terms of a scalar  $\psi$  (known as a *stream function*) such that

$$u_x = -\frac{\partial \psi}{\partial y}, u_y = \frac{\partial \psi}{\partial x}$$

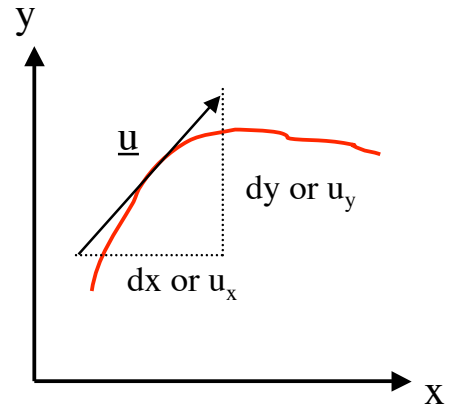
- Hence  $\underline{u}$  is divergence-free

A *streamline* is defined as a curve that has  $\underline{u}$  in the tangential direction,

$$\frac{dx}{u_x} = \frac{dy}{u_y}$$

$$\Rightarrow u_y dx - u_x dy = 0$$

$$\Rightarrow \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy = d\psi = 0$$



Hence  $\psi$  is constant on a streamline.

## Reminder

- If  $Q$  is a scalar:

$$\underline{u} \cdot \nabla Q = u_x \frac{\partial Q}{\partial x} + u_y \frac{\partial Q}{\partial y} + u_z \frac{\partial Q}{\partial z} \quad \text{cartesians}$$

$$= u_r \frac{\partial Q}{\partial r} + \frac{u_\theta}{r} \frac{\partial Q}{\partial \theta} + \frac{u_\phi}{r \sin \theta} \frac{\partial Q}{\partial \phi} \quad \text{sphericals}$$

$$= u_R \frac{\partial Q}{\partial R} + u_z \frac{\partial Q}{\partial z} + \frac{u_\phi}{R} \frac{\partial Q}{\partial \phi} \quad \text{cylindricals}$$

- If  $\underline{Q}$  is a vector, then  $(\underline{u} \cdot \underline{\nabla})\underline{Q}$  is also a vector, each component of which is  $(\underline{u} \cdot \underline{\nabla})$  acting on each component of  $\underline{Q}$ .
- Hence in cartesian:

$$\begin{aligned}
 (\underline{u} \cdot \underline{\nabla})\underline{Q} &= \left( u_x \frac{\partial}{\partial x} + u_y \frac{\partial}{\partial y} + u_z \frac{\partial}{\partial z} \right) (Q_x, Q_y, Q_z) \\
 &= \left[ \begin{aligned}
 &u_x \frac{\partial Q_x}{\partial x} + u_y \frac{\partial Q_x}{\partial y} + u_z \frac{\partial Q_x}{\partial z}, \\
 &u_x \frac{\partial Q_y}{\partial x} + u_y \frac{\partial Q_y}{\partial y} + u_z \frac{\partial Q_y}{\partial z}, \\
 &u_x \frac{\partial Q_z}{\partial x} + u_y \frac{\partial Q_z}{\partial y} + u_z \frac{\partial Q_z}{\partial z}
 \end{aligned} \right]
 \end{aligned}$$

## Question 1

- The temperature variation in a river is

$$T(x, t) = e^x \sin t$$

- And the river flows with velocity

$$\underline{u}(x, t) = (xt^2, 0)$$

- Write down both the Lagrangian and Eulerian temperature derivatives.

# Answer 1

- Did you get:

$$\frac{dT}{dt} = \frac{\partial T}{\partial t} + u_x \frac{\partial T}{\partial x} = e^x \cos t + xt^2 e^x \sin t$$