

13. Water waves

- Why do waves break when they reach the shore?



- How can you tell if a storm is coming?



13. 1 How do water waves propagate?

Assume:

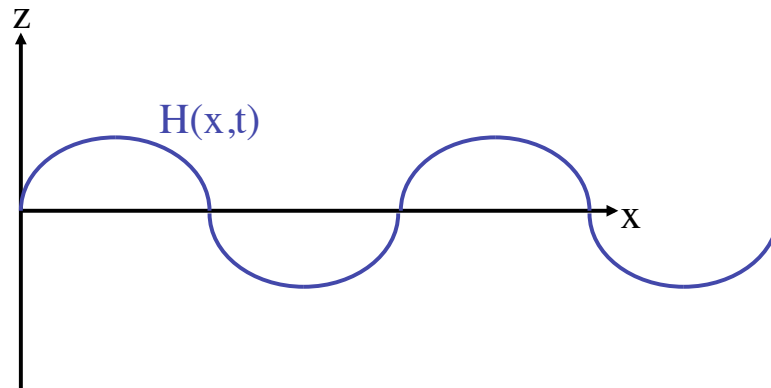
1. Incompressible, irrotational flow \Rightarrow
Laplace's equation

$$\nabla \times \underline{u} = 0, \quad \underline{u} = \nabla \Phi$$

$$\nabla \cdot \underline{u} = 0, \quad \Rightarrow \quad \nabla^2 \Phi = 0 \quad (13.1)$$

“potential” flow

2. Atmospheric (uniform) pressure at the water surface ($\Rightarrow p = p_0$ and ρ constant too)



Surface boundary condition:

- Match pressures at the surface using Bernoulli's equation (7.2) for time-dependent but *ideal* (zero vorticity) flow:

$$\frac{\partial u}{\partial t} = -\nabla \left(\frac{1}{2} u^2 + \frac{p}{\rho} + \psi \right)$$

$$\nabla \left(\frac{\partial \Phi}{\partial t} + \frac{1}{2} u^2 + \frac{p}{\rho} + \psi \right) = 0$$

Small \Rightarrow neglect

Uniform near surface

- Hence at the water surface ($z = H(x,t)$)

$$\frac{\partial \Phi}{\partial t} + \psi = F(t)$$

$= g H(x,t)$
 $\rightarrow 0$ as no spatial dependence

- So
$$H(x,t) = \frac{\partial \Phi / \partial t}{-g} \quad (13.2)$$

- Also, a parcel of fluid near the surface moves up and down *with the surface*

$$u_z = \frac{\partial \Phi}{\partial z} = \frac{\partial H}{\partial t}$$

- So differentiating (13.2) \Rightarrow

$$\frac{\partial^2 \Phi}{\partial t^2} + g \frac{\partial \Phi}{\partial z} = 0 \quad (13.3)$$


Method:

- Look for separable solutions of Laplace's equation (13.1):

$$\Phi = f(z)\phi(x)q(t)$$

- So

$$\nabla^2 \Phi = 0 \quad \Rightarrow \quad \phi_{xx} f + \phi f_{zz} = 0$$

$$\frac{f_{zz}}{f} = -\frac{\phi_{xx}}{\phi} = k^2$$


Positive to get $f \rightarrow 0$ below surface

- Choose

$$q(t) = e^{-i\omega t}$$

$$\phi(x) = e^{ikx}$$

- i.e.

$$\Phi = f(z)e^{i(kx - \omega t)}$$

$$f = Ae^{kz} + Be^{-kz}$$

- Choice for $f(z)$ depends on whether deep or shallow water

13.2 Deep water waves

- Choose $B = 0$ so that

$$f \rightarrow 0 \quad \text{as} \quad z \rightarrow -\infty$$

$$\Phi = C e^{kz} e^{i(kx - \omega t)} \quad (13.4)$$

- Note that wave penetrates only to about $z = -\lambda$ since here

$$e^{kz} = e^{z2\pi/\lambda} = e^{-2\pi} \approx 10^{-3}$$

- Boundary condition (13.3) now gives

$$-\omega^2 + gk = 0$$

- So $\omega = \sqrt{gk}$

- And the wave speed

$$c = \frac{\omega}{k} = \sqrt{\frac{g}{k}} = \sqrt{\frac{g\lambda}{2\pi}}$$

- Waves are *dispersive* (long λ i.e. small k travel faster)
- First sign of Atlantic storm is long waves of small amplitude (swell) period ~ 30 s, speed $\sim 47 \text{ ms}^{-1}$, travelling ~ 4000 km in a day.
- Hence arrive before storm (takes several days to cross Atlantic)
- Followed by shorter waves that travel more slowly

13.3 Shallow water waves

- Still have

$$\Phi = f(z)e^{i(kx - \omega t)}$$

$$f = Ae^{kz} + Be^{-kz}$$

- But choose

$$f = C \cosh[k(z + h)]$$

- So that $u_z \sim f' \rightarrow 0$ at $z = -h$

- So

$$\Phi = C \cosh[k(z + h)] e^{i(kx - \omega t)} \quad (13.5)$$

- (13.3) now gives (at $z = 0$)

$$(i\omega)^2 \cosh(kh) + gk \sinh(kh) = 0$$

- So

$$\omega = \sqrt{gk \tanh(kh)}$$

- And the wave speed

$$c = \frac{\omega}{k} = \sqrt{\frac{g}{k} \tanh(kh)}$$

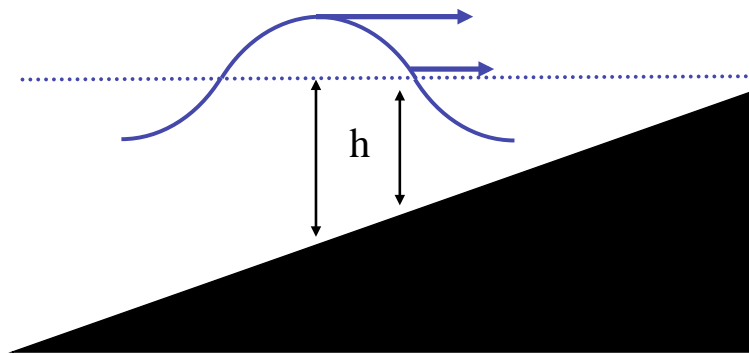
- Limits:

Deep water $h \rightarrow \infty, \quad c \rightarrow \sqrt{\frac{g}{k}} \quad (\text{as before})$

Shallow water $kh \rightarrow 0, \quad c \rightarrow \sqrt{gh} \quad (\text{indep. of } k)$

As h decreases, so does wave speed c

- Hence as waves approach shore:



- The back of the wave has a larger h \rightarrow larger speed so it catches up on the front of the wave.
- Wave steepens and breaks \rightarrow Surf!

13.4 Particle paths: deep water

- Write the particle position (x, z) as a displacement $(X(t), Z(t))$ from a mean position (x_0, z_0) .

$$x = x_0 + X(t)$$

$$z = z_0 + Z(t)$$

- Now from (13.4)

$$\Phi = Ce^{kz} e^{i(kx - \omega t)}$$

- So taking real parts

$$X'(t) = u_x(x_0, z_0) = \frac{\partial \Phi}{\partial x} = -kCe^{kz_0} \sin(kx_0 - \omega t)$$

$$Z'(t) = u_z(x_0, z_0) = \frac{\partial \Phi}{\partial z} = kCe^{kz_0} \cos(kx_0 - \omega t)$$

- Integrating gives

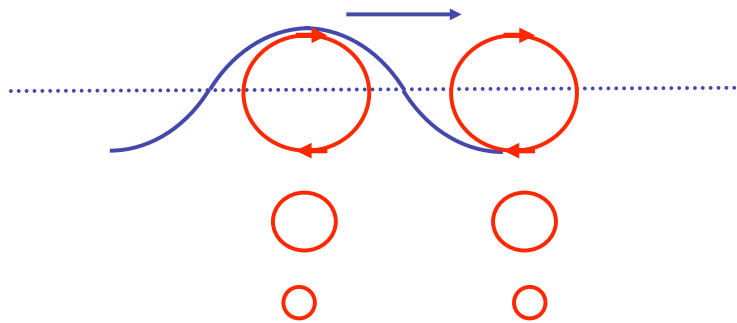
$$X(t) = -\frac{kC}{\omega} e^{kz_0} \cos(kx_0 - \omega t)$$

$$Z(t) = -\frac{kC}{\omega} e^{kz_0} \sin(kx_0 - \omega t)$$

- The particle paths are therefore circles,

$$R = \frac{kC}{\omega} e^{kz}$$

whose radius is the wave amplitude at $z = 0$ and decreases with z .



13.5 Particle paths: shallow water

- Once again

$$x = x_0 + X(t)$$

$$z = z_0 + Z(t)$$

- And with (13.5)

$$\Phi = C \cosh[k(z + h)] e^{i(kx - \omega t)}$$

- Obtain

$$X(t) = \frac{-kC}{\omega} \cosh[k(z_0 + h)] \cos(kx_0 - \omega t)$$

$$Z(t) = \frac{-kC}{\omega} \sinh[k(z_0 + h)] \sin(kx_0 - \omega t)$$

- Ellipses

$$\left\{ \frac{x - x_0}{\cosh[k(z_0 + h)]} \right\}^2 + \left\{ \frac{z - z_0}{\sinh[k(z_0 + h)]} \right\}^2 = \left(\frac{kC}{\omega} \right)^2$$

- h very large \Rightarrow circular motion
- h very small \Rightarrow almost horizontal motion

